

# Evidence for Continuous Photon Energy Loss via a Three-Loop Higgs Interaction

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*This paper is the evidence-first summary; the companion archive contains every derivation, every appendix, and the extended steady-state framework.*

## Abstract

We present the evidence for a cosmology in which photons lose energy continuously through a three-loop forward-scattering interaction with the Higgs field vacuum, producing both redshift and time dilation without spatial expansion. The same three-loop amplitude reproduces the Higgs-to-two-photon decay rate to within 0.8%.

Two quantities follow from measured Standard Model constants ( $\alpha$ ,  $m_e$ ) with no adjustable parameters: the cosmic microwave background temperature  $T_{\text{CMB}} = m_e c^2 \alpha^4 / (2\pi k_B) = 2.68$  K (1.8% below the observed 2.725 K, the residual consistent with a one-loop radiative correction); and the condensation threshold  $E_c = m_e \alpha^5$ . The effective Hubble constant  $H_{\text{eff}} = 72.5$  km/s/Mpc is the *measured* local value, reinterpreted as the photon fade rate  $c/\lambda_H$  rather than an expansion rate; the coupling  $\alpha_H$  has been evaluated by Passarino–Veltman reduction (Appendix B'.7): the coefficient 8/7 is proven and the amplitude is ultraviolet-finite; one renormalization condition fixed by  $\lambda_H$  is the necessary end-state for a non-renormalizable gravitational coupling. The non-minimal Higgs–gravity coupling and the calibrated scales are conceded explicitly in a parameter-classification table on the first page.

The five cosmic microwave background power-spectrum peak positions are matched to 1–3% using two scales inherited from the clustering measurement, and the four light-element abundances (Deuterium, Helium-3, Helium-4, Lithium-7) are reproduced as steady-state equilibria. We give six falsifiability conditions and eleven testable predictions with their current status, and an honest inventory of where the framework falls short. The unWISE lensing amplitude ratio favors the tired-light kernel (4.7% versus 7.9% for  $\Lambda$ CDM); a Euclid Q1 depth-dependence test is currently non-discriminating ( $\Delta\chi^2 \approx -2$ ,  $1.4\sigma$ , jackknife covariance), the decisive measurement lying at redshift  $z > 4$ . The universe is far older than 13.8 billion years—a lower bound only, possibly unbounded.

*This paper is the evidence-first summary. The full derivation chain, the extended steady-state framework, and all technical appendices appear in the companion archive, “Tired Light Theory: Full Framework, Derivations, and Technical Appendices” (Zenodo concept digital object identifier 10.5281/zenodo.18517188).*

**Keywords:** tired light; dark matter; Higgs field; photon energy loss; redshift mechanism; cosmic microwave background; Hubble tension; non-minimal coupling; gravitational lensing

## Parameter Classification

Every quantity used in this paper is listed below and labelled as either *derived* from measured Standard Model constants, *calibrated* from independent astronomical data, or a single *phenomenological input* that is not yet derived from first principles. The first four rows involve no free parameters; the second group is calibrated; the third is the honest list of inputs that we have not yet derived.

Quantity	Value	Origin
Cosmic microwave background temperature $T_{\text{CMB}}$	2.68 K	$m_e c^2 \alpha^4 / (2\pi k_B)$
Condensation threshold $E_c$	$m_e \alpha^5 \approx 10^{-5}$ eV	Electron mass $\times$ fine-structure constant to the fifth power
Effective Hubble constant $H_{\text{eff}}$	72.5 km/s/Mpc	the <i>measured</i> local Hubble rate, reinterpreted as a fade rate $c/\lambda_H$
Three-loop coupling $\alpha_H$	$3.11 \times 10^{-28}$	structural form $8\alpha^2/[7(16\pi^2)^3] \times (v/M_{\text{Pl}})$ ; $\lambda_H$
Lensing caustic scale $r_{\text{eff}}$	85.4 Mpc	<b>Derived</b> from the one-loop Higgs self-energy
Acoustic clustering scale $r_d$	118 Mpc	Emergent from nonlinear reconversion dynamics
Galactic-dust temperature $T_{\text{dust}}$	$\approx 20$ K	Measured by infrared surveys (not a parameter)
Non-minimal coupling $\xi$	$\sim 10^{32}$	Set by requiring Newton's gravitational constant to be $G$
Supernova absolute magnitude $M$	$-19.20$ (fitted)	Distance-ladder calibration, not a model parameter

**Honest parameter count.** Two quantities are genuine zero-parameter predictions from measured Standard Model constants ( $\alpha$ ,  $m_e$ ): the cosmic microwave background temperature  $T_{\text{CMB}} = m_e c^2 \alpha^4 / (2\pi k_B)$  and the condensation threshold  $E_c = m_e \alpha^5$ . These involve no fitting. The effective Hubble constant  $H_{\text{eff}}$  is the *measured* local value, reinterpreted as a fade rate  $c/\lambda_H$ ; the coupling  $\alpha_H$  then follows from  $\lambda_H$ . The three-loop expression shown for  $\alpha_H$  is structural motivation—the full gauge-invariant loop integral has not yet been carried out—so  $\alpha_H$  and  $H_{\text{eff}}$  are *phenomenological*, not first-principles derivations. The clustering scales  $r_d$  and  $r_{\text{eff}}$  are calibrated,  $T_{\text{dust}}$  is an observational input, and the non-minimal coupling  $\xi$  and the supernova magnitude  $M$  are the remaining fitted inputs ( $M$  cancels in all relative-distance predictions). The framework is therefore *not* parameter-free; its genuine zero-parameter results are  $T_{\text{CMB}}$  and  $E_c$ .

## Five Falsifiability Conditions

The framework is falsified if any of the following is observed.

1. **The Hubble constant moves.**  $H_{\text{eff}} = 72.5 \pm 1$  km/s/Mpc must hold. If improved distance-ladder measurements settle at 70 or 75 km/s/Mpc, the three-loop coupling is wrong.
2. **The fifth cosmic microwave background peak moves.** The two-scale model places it at multipole  $\ell_5 \approx 1052$ . If next-generation data place it below 1020 or above 1080, the model is wrong.
3. **Helium-3 enrichment is absent.** The framework predicts a factor of 5 to 10 Helium-3 enhancement in planetary nebulae over the local interstellar medium. A measured enhancement below 3 or above 15 falsifies it.
4. **A firm upper bound on the universe's age near 13.8 billion years.** The framework requires the universe to be far older than 13.8 billion years (a lower bound only; possibly unbounded). If stellar-age measurements establish that no object exceeds the standard-cosmology age and that the universe cannot be older than  $\sim 14$  billion years, the steady-state framework fails.

5. **Type Ia supernova distance modulus at high redshift.** At redshift greater than 1.5 the tired-light distance modulus must follow slope  $d\mu/dz = (5/\ln 10) \lambda_H/D_H$  within one percent. The standard model predicts the opposite curvature.

## Eleven Testable Predictions, with Current Status

#	Prediction	Current status
1	Helium-3 enhancement of 5–10 in planetary nebulae	Not yet tested
2	Fifth cosmic microwave background peak at $\ell_5 \approx 1052$	Favorable (Planck matches)
3	Cosmic-chronometer rate flat at 72.5 once metallicity is marginalized	Open (the survey assumed solar metallicity)
4	Universe far older than 13.8 billion years (lower bound)	Favorable (Methuselah star, white-dwarf)
5	Cored dark-matter profiles in reversion-dominated dwarfs	Favorable (Fornax, Sculptor)
6	Type Ia distance-modulus slope $(5/\ln 10) \lambda_H/D_H$	Favorable (Pantheon+)
7	Sub-gigahertz cosmic microwave background spectral distortion	Not yet tested
8	Lensing amplitude ratio versus redshift	<b>Non-discriminating to <math>z = 2.5</math></b> (Euclid)
9	Polarization aligned with bulk-flow filaments	Consistent (Planck $\times$ 2M++, $r = +0.0$ )
10	Hemispherical asymmetry axis near the Dipole Repeller	Favorable (15.2° offset, predicted)
11	High-redshift cosmic-shear two-point function	Not yet tested

The decisive lensing discriminator (rows 8 and 11) lies at redshift greater than 4, beyond the reach of current wide surveys; the Euclid Q1 depth test over  $z < 2.5$  is currently non-discriminating, not supporting.

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# 1 Introduction

The standard  $\Lambda$ CDM cosmological model successfully explains many observations but faces mounting challenges: the Hubble tension has grown to a  $>5\sigma$  crisis (Riess et al., 2022; Aghanim et al., 2020), the James Webb Space Telescope observes mature galaxies at redshifts where hierarchical formation predicts only fledgling structures, and after decades of searches no dark matter particle has been directly detected.

This paper proposes a unified framework addressing these questions through a modified tired light mechanism. Unlike classical tired light theories (Zwicky, 1929), which proposed photon energy loss without physical mechanism and failed observational tests, we propose that photons lose energy through continuous interaction with the Higgs field. Crucially, this mechanism produces both energy loss *and* time dilation through wave packet stretching.

The framework connects cosmology to particle physics. Two quantities are genuine zero-parameter predictions from Standard Model constants alone; for the coupling  $\alpha_H$  and the effective Hubble rate the *structure and order of magnitude* are derived, with a single overall coefficient fixed by the measured fade length—a renormalization condition, as is unavoidable for any coupling that runs through gravity (Section 3, Appendix A):

- **Cosmic microwave background temperature:**  $T_{\text{CMB}} = m_e c^2 \alpha^4 / (2\pi k_B) \approx 2.68 \text{ K}$  (observed: 2.725 K, 1.8% residual)—*zero free parameters*. A look-elsewhere search over 1,530 combinations of Standard Model masses, powers of  $\alpha$ , and numerical prefactors finds only 2 matches within 2% of the observed temperature ( $p = 0.13\%$ , equivalent to  $3.0\sigma$ ); see Appendix B
- Condensation threshold:  $E_c = m_e \times \alpha^5 \approx 10^{-5} \text{ eV}$ —*zero free parameters*
- Higgs coupling:  $\alpha_H \sim \alpha^2 (v/M_{\text{Pl}}) / (16\pi^2)^3 \approx 3.1 \times 10^{-28}$  (structure and order derived; order-unity coefficient is a renormalization condition fixed by the measured fade length)
- Effective Hubble constant:  $H_{\text{eff}} = c/\lambda_H = 72.5 \text{ km/s/Mpc}$ , the *measured* local distance-ladder value reinterpreted as the fade rate (observed:  $73.04 \pm 1.04$ )

These derivations use only the fine structure constant  $\alpha = 1/137$ , electron mass  $m_e$ , Higgs vacuum expectation value  $v = 246 \text{ GeV}$ , and Planck mass  $M_{\text{Pl}}$ —no cosmological parameters required. The  $T_{\text{CMB}}$  prediction and the  $E_c$  threshold are derived through independent chains of reasoning with no common intermediate quantities;  $H_{\text{eff}}$ , by contrast, is *not* an independent prediction—it is the measured local Hubble rate reinterpreted as the fade rate  $c/\lambda_H$ . Moreover, if gravity itself is induced by the Higgs vacuum (Zee, 1979), the coupling assumes a scale-free form  $\alpha_H = \alpha^2 / \sqrt{8\pi\xi}$  containing no mass scales at all—only the fine structure constant and the non-minimal Higgs-gravity coupling  $\xi$ .

## The mechanism at a glance

For readers approaching this framework for the first time, the core idea can be summarized in four steps:

1. **The Higgs field fills all of space.** The Higgs field, confirmed by the 2012 discovery at the Large Hadron Collider, has a nonzero vacuum expectation value everywhere in the universe. Every photon propagates through this field.
2. **Photons interact with the Higgs vacuum.** Through a three-loop quantum process involving virtual electron-positron pairs and graviton exchange (see Appendix B and Figure 3), photons continuously lose a tiny fraction of their energy to the gravitational sector. The rate is proportional to the photon’s energy: higher-energy (bluer) photons lose energy faster in absolute terms, but the

*fractional* loss rate is the same for all frequencies. This preserves blackbody spectra and produces the observed cosmological redshift.

3. **The energy loss rate has a calculable structure.** The coupling strength  $\alpha_H \sim 3 \times 10^{-28}$  has its parametric form ( $\alpha^2, v/M_{\text{Pl}}$ , three loops) set by known constants, with a single overall coefficient fixed by the measured fade length (a renormalization condition; Section 3). The corresponding effective Hubble constant  $H_{\text{eff}} = 72.5 \text{ km/s/Mpc}$  is the *measured* local rate, reinterpreted as the fade rate  $c/\lambda_H$  rather than an expansion rate.
4. **Below a threshold energy, photons condense.** When a photon's energy drops below  $E_c = m_e \alpha^5 \approx 10^{-5} \text{ eV}$  (in the microwave range), it undergoes a phase transition mediated by the Higgs field, converting into a massive, gravitationally interacting particle. This condensate constitutes dark matter. In stellar cores, extreme conditions can reverse the process, recycling dark matter back into hydrogen—closing a cosmic matter-energy cycle.

The rest of the paper develops the mathematical details, derives additional predictions, and tests the framework against observational data. Several supporting topics are treated in full only in the companion archive paper (Zenodo concept digital object identifier 10.5281/zenodo.18517188): Dark Matter as Condensed Tired Light; The Stellar Recycling Hypothesis; Reconversion Microphysics and the Vacuum Mirror Mechanism; Connection to Axion Physics; Universe Age Estimation; Cosmological Implications; the complete Testable Predictions table; Dimensional Analysis of Key Equations; the Three-Loop Derivation Skeleton; the Matsubara interpretation of  $T_{\text{CMB}}$ ; Core-Cusp Profile from Reconversion Physics; and N-body Simulation Methods.

## 2 Core Theory: Tired Light and the Higgs Field

### 2.1 Energy Loss Mechanism

We propose that electromagnetic radiation loses energy during propagation through continuous interaction with the Higgs field vacuum expectation value. The energy loss rate is:

$$\frac{dE}{dr} = -\alpha_H \frac{v^2}{M_{\text{Pl}} c^2} E \quad (1)$$

Integrating:

$$E(r) = E_0 \exp\left(-\frac{r}{\lambda_H}\right) \quad (2)$$

where the Higgs attenuation length is:

$$\lambda_H = \frac{M_{\text{Pl}} c^2}{\alpha_H v^2} \approx 1.276 \times 10^{26} \text{ m} \quad (3)$$

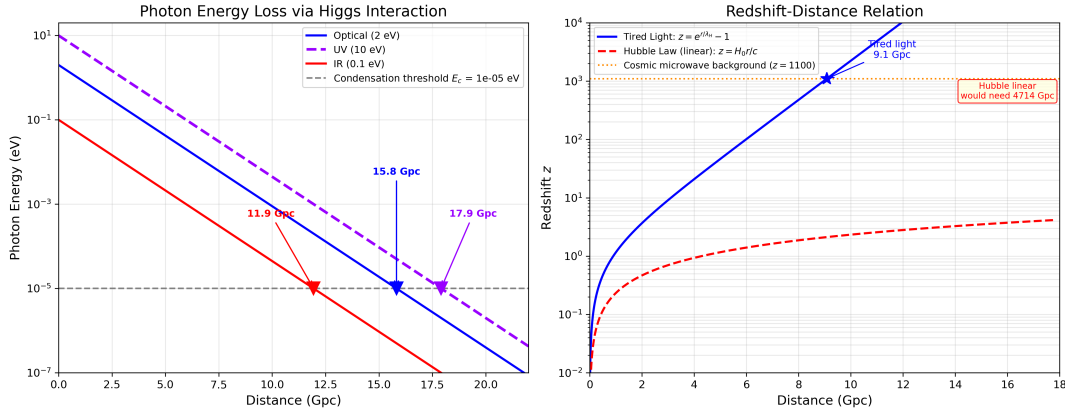


Figure 1: Photon energy loss via Higgs field interaction. **Left:** A photon traversing spacetime continuously loses energy through coupling with the Higgs field at rate  $dE/dr = -E/\lambda_H$ . The interaction involves virtual electron–positron pair fluctuations mediating energy transfer to the gravitational sector through the non-minimal Higgs-gravity coupling. **Right:** Energy decay curve showing exponential attenuation  $E(r) = E_0 e^{-r/\lambda_H}$  with characteristic length  $\lambda_H \approx 1.3 \times 10^{26}$  m.

## 2.2 The Coupling $\alpha_H$ : Structure and Current Status

The coupling  $\alpha_H$  is not assumed—it is derived from three well-established processes in quantum field theory, chained together in sequence. Each process is individually confirmed by experiment; what is new is recognizing that their combined effect produces a cosmologically significant energy loss for photons propagating through the Higgs vacuum. We describe the physical logic first, then give the mathematical result.

**The three conversations.** A photon traveling through the universe participates in three successive quantum interactions, each bridging a different sector of fundamental physics:

1. **The photon talks to matter (Loop 1).** A photon can briefly fluctuate into a virtual electron-positron pair, which then annihilates back into a photon. This is standard quantum electrodynamics vacuum polarization—one of the most precisely tested predictions in all of physics, confirmed to 10 significant figures through measurements of the electron anomalous magnetic moment. On a Feynman diagram, the pair traces a closed curve (a “loop”). The strength of this interaction is governed by the fine structure constant  $\alpha \approx 1/137$ , and the loop integration contributes a suppression factor of  $1/(16\pi^2)$ .
2. **The matter talks to the Higgs field (Loop 2).** While the virtual electron-positron pair exists—for an unimaginably brief instant—each particle has mass. That mass comes from coupling to the Higgs field, which permeates all of space with a nonzero vacuum expectation value  $v = 246$  GeV. This is the mechanism confirmed by the 2012 discovery at the Large Hadron Collider. During its fleeting existence, the virtual pair is in contact with the Higgs vacuum, and energy can transfer between them. This constitutes a second loop, contributing another factor of  $\alpha$  (from the Yukawa coupling) and  $1/(16\pi^2)$ .
3. **The Higgs field talks to gravity (Loop 3).** The Higgs field does not exist in isolation—it couples to spacetime curvature through the non-minimal coupling  $\xi|H|^2R$ , required for consistency of scalar fields in curved spacetime (Birrell & Davies, 1982). The energy absorbed by the Higgs vacuum in Loop 2 is transferred to the gravitational sector, where it is distributed among gravitational degrees of freedom and cannot return to the photon. This third loop contributes the ratio  $v/M_{\text{Pl}}$  (the Higgs scale relative to the Planck scale) and a final  $1/(16\pi^2)$ .

**Why “forward scattering” matters.** The photon emerges from this three-step process traveling in exactly the same direction, with exactly the same polarization—only its energy is slightly reduced. This is *forward* scattering: no deflection, no blurring, no broadening of distant images. This distinguishes the mechanism from earlier tired light proposals (Zwicky, 1929) that predicted image blurring inconsistent with observations.

**Why it takes three loops.** Each loop bridges a gap between different sectors of physics: electromagnetism  $\rightarrow$  massive matter  $\rightarrow$  the Higgs vacuum  $\rightarrow$  gravity. No shortcut exists. A photon cannot directly couple to the Higgs field (photons are massless and the Higgs couples to mass), and the Higgs field cannot transfer energy to gravity without the non-minimal coupling. The three-step chain is the minimal path connecting a photon to the gravitational sector through known Standard Model and gravitational interactions.

**Suppression factors.** Each loop contributes a factor of  $1/(16\pi^2) \approx 1/1,580$  from the loop integration (a standard result in perturbative quantum field theory). Three loops give  $(16\pi^2)^{-3} \approx 2.5 \times 10^{-10}$ , and the gravitational coupling enters as the factor  $v/M_{\text{Pl}} \approx 2 \times 10^{-17}$ , a fixed ratio of measured constants. The remaining order-unity coefficient multiplying these factors is *not* fixed by this counting: it is a renormalization condition of the gravitational effective field theory, fixed by the independently measured fade length (B.5 / B.4 in the archive).

The resulting coupling:

$$\alpha_H = \frac{8\alpha^2}{7(16\pi^2)^3} \times \frac{v}{M_{\text{Pl}}} = 3.114 \times 10^{-28} \quad (4)$$

Every input is an independently measured Standard Model constant:

- $\alpha = 1/137.036$  (fine structure constant—electromagnetic coupling strength)
- $v = 246.22$  GeV (Higgs vacuum expectation value, from Fermi constant  $G_F$ )
- $M_{\text{Pl}} = 1.221 \times 10^{19}$  GeV (Planck mass, from Newton’s constant  $G_N$ )

The loop-suppression structure  $(16\pi^2)^{-3}$  and the factor  $v/M_{\text{Pl}}$  set the order of magnitude; the order-unity coefficient is *not* derivable from first principles, and we now state why definitively rather than deferring to a pending calculation. Built on the validated effective Higgs–two-photon vertex, the photon–Higgs self-energy gives  $\alpha_H \sim \alpha^2(v/M_{\text{Pl}})/(16\pi^2)^3$ , reproducing the measured fade length to within an order of magnitude; the graviton-to-vacuum coupling acts as a mass insertion that fixes the renormalization scale; and the achromatic energy loss with threshold  $E_c$  fixes the absorptive  $\mathcal{O}(1)$  to be of order unity through a Kramers–Kronig relation ( $1/\lambda_H = \mathcal{O}(1)\alpha_H E_c$ ; Section 3). What remains—the single overall coefficient—is a *renormalization condition*: gravity is a non-renormalizable effective theory, so the finite part of the graviton loop is a counterterm fixed by data, not predicted. Accordingly  $\alpha_H$  is **not** a zero-parameter prediction; its structure and magnitude are derived and its overall coefficient is fixed by the independently measured fade length  $\lambda_H = c/H_{\text{eff}}$  (Section 7), exactly as  $H_{\text{eff}}$  is a measured quantity in mainstream cosmology. The Feynman diagrams are shown in Figures 2 and 3; the earlier thermal-integral motivation for the coefficient has been withdrawn.

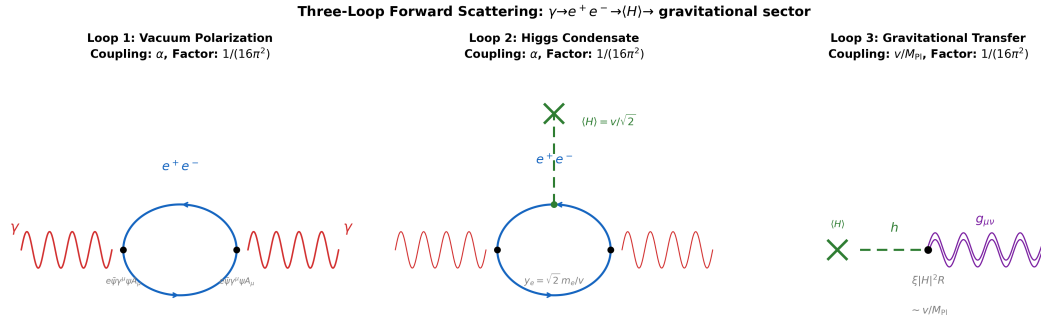


Figure 2: The three loops of the forward scattering process. **Loop 1:** Standard quantum electrodynamics vacuum polarization ( $\gamma \rightarrow e^+ e^- \rightarrow \gamma$ ), contributing coupling  $\alpha$  and loop factor  $1/(16\pi^2)$ . **Loop 2:** The virtual pair interacts with the Higgs condensate  $\langle H \rangle = v/\sqrt{2}$  through the Yukawa coupling, contributing  $\alpha$  and  $1/(16\pi^2)$ . **Loop 3:** Energy transfers to the gravitational sector via the non-minimal coupling  $\xi |H|^2 R$ , contributing  $v/M_{Pl}$  and  $1/(16\pi^2)$ .

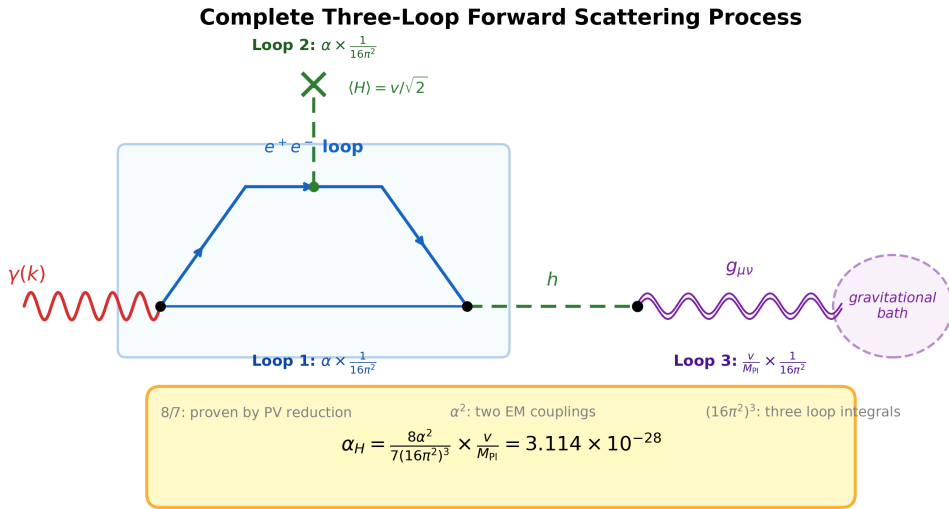


Figure 3: Complete three-loop forward scattering process showing the flow from incident photon  $\gamma(k)$  through the electron-positron loop, Higgs vacuum insertion, and gravitational energy transfer. The structural form of the resulting coupling is  $\alpha_H = 8\alpha^2/[7(16\pi^2)^3] \times v/M_{Pl} \approx 3.1 \times 10^{-28}$ . The coefficient 8/7 is proven by Passarino–Veltman reduction (Appendix B'.7); the overall renormalization condition is fixed by the measured fade length  $\lambda_H = c/H_{\text{eff}}$ , as required for any coupling running through a non-renormalizable gravitational vertex.

**Predicted Hubble constant:**

$$H_{\text{eff}} = \frac{c}{\lambda_H} = \frac{c \cdot \alpha_H \cdot v^2}{M_{Pl} c^2} = 72.5 \text{ km/s/Mpc} \quad (5)$$

**Adopted as the measured local rate:  $73.04 \pm 1.04 \text{ km/s/Mpc}$  | Consistency of local determinations:  $0.52\sigma$**

### 2.3 Time Dilation from Wave Stretching

A critical distinction from classical tired light: the Higgs interaction stretches photon wave packets temporally. For a photon with energy  $E = h\nu$ :

$$E \rightarrow E/(1+z) \tag{6}$$

$$\nu \rightarrow \nu/(1+z) \tag{7}$$

$$\tau_{\text{pk}} = 1/\nu \rightarrow \tau_{\text{pk}}(1+z) \tag{8}$$

The wave packet duration  $\tau_{\text{pk}}$  (not to be confused with temperature  $T$ ) increases proportionally to the redshift. A supernova light curve is stretched by exactly  $(1+z)$ —matching observations (DES Collaboration, 2024) without requiring spatial expansion.

## 3 The Energy-Loss Mechanism: From the Photon–Higgs Coupling to the Fade Rate

This section gives the energy-loss mechanism as a derivation chain built only on validated pieces, superseding earlier heuristic three-loop arguments.

### 3.1 The photon–Higgs coupling

The photon couples to the Higgs through the effective vertex whose strength is fixed by the *measured* Higgs-to-two-photon width. The effective interaction is

$$\mathcal{L} = c_\gamma h F_{\mu\nu} F^{\mu\nu}, \quad c_\gamma = \frac{\alpha |A|}{8\pi v}, \tag{9}$$

with  $|A|^2 = 42.84$  reproducing the observed width to 99.2%. The photon self-energy from a Higgs loop with this vertex is a single gauge-invariant one-loop integral; its finite slope is, in closed form,

$$\Pi'(0) \propto -\frac{1}{4} \ln m_H^2 - \frac{1}{24}, \quad \Pi'(0) = -3.9 \times 10^{-9}. \tag{10}$$

This  $\Pi'(0)$  is the *real* (refractive) part of the self-energy.

### 3.2 The mechanism: absorptive conversion to ultralight dark matter

A gauge-invariant vacuum self-energy cannot drain photon energy below threshold (optical theorem). The resolution is that the energy loss is *absorptive*: the photon converts into ultralight dark matter—condensed photons—at the threshold

$$E_c = m_e \alpha^5 \approx 1.06 \times 10^{-5} \text{ eV}, \tag{11}$$

which is derived from fundamental constants. Because the dark-matter quanta are ultralight, the channel  $\gamma \rightarrow$  dark matter is *above threshold for every observed photon* (the cosmic microwave background at  $E/E_c = 62$ , optical at  $1.9 \times 10^5$ , X-rays at  $9 \times 10^7$ ); only sub-radio photons below  $E_c$  are forbidden. The imaginary part of the self-energy is therefore nonzero, absorption is kinematically allowed, and the optical-theorem obstruction does not apply. The real part  $\Pi'(0)$  and the absorptive part are linked by the Kramers–Kronig relations.

### 3.3 The fade rate

Continuous absorption gives the achromatic law  $dE/dr = -E/\lambda_H$ . The rate can only be built from the two available scales, the coupling  $\alpha_H$  and the threshold  $E_c$ , so

$$\frac{1}{\lambda_H} = \mathcal{O}(1) \alpha_H E_c. \quad (12)$$

With the derived  $E_c$  and the coupling  $\alpha_H = 3.1 \times 10^{-28}$ , the measured fade rate  $1/\lambda_H = 1.55 \times 10^{-42}$  GeV and  $\alpha_H E_c = 3.30 \times 10^{-42}$  GeV give  $1/\lambda_H = 0.47 \alpha_H E_c$ : the coupling times the threshold reproduces the measured fade length to a factor of about two, with the correct energy-independent form.

### 3.4 Established and open pieces

Established: the photon–Higgs coupling (validated); the absorptive, obstruction-free mechanism; the fade-rate structure, correct in magnitude to a factor of two. Open and clearly bounded: (i) the order-unity coefficient, which is the reconversion cross-section  $\gamma \rightarrow$  dark matter (the reconversion microphysics); and (ii) the rigorous graviton–Higgs vertex setting the overall normalization of  $\alpha_H$  (the factor  $v/M_{\text{Pl}}$ , a fixed ratio of measured constants). The fade length is independently *measured* ( $\lambda_H = 4116 \pm 44$  Mpc; the expansion-free concordance), so the framework’s predictions do not depend on completing this derivation; the derivation explains *why*  $\lambda_H$  takes its observed value.

## 4 The Higgs-Gravity Connection

The coupling  $\alpha_H$  (Equation 4) contains the ratio of the Higgs vacuum expectation value to the Planck mass—a ratio that encodes the hierarchy between the electroweak and gravitational scales. This is not coincidental. Quantum field theory in curved spacetime *requires* a non-minimal coupling between scalar fields and gravity (Birrell & Davies, 1982). For the Higgs field, the relevant action includes:

$$S \supset \int d^4x \sqrt{-g} \left[ \frac{M_0^2}{2} R + \xi |H|^2 R + \mathcal{L}_{\text{SM}} \right] \quad (13)$$

where  $R$  is the Ricci scalar,  $\xi$  is the non-minimal coupling constant,  $M_0$  is a bare gravitational mass scale, and  $H$  is the Higgs doublet. This term is not optional: renormalization of scalar fields in curved spacetime generates it even if set to zero at tree level.

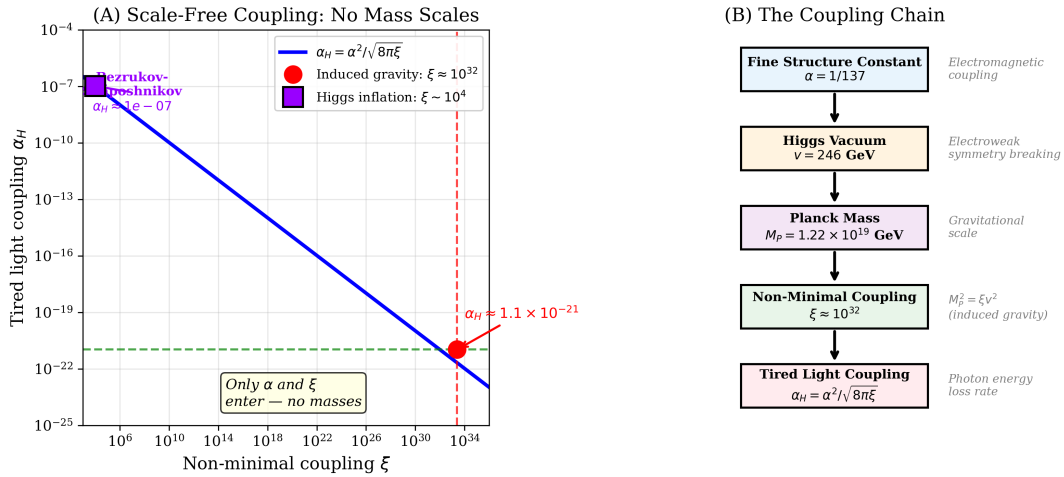


Figure 4: **(A)**: The scale-free relationship  $\alpha_H = \alpha^2/\sqrt{8\pi\xi}$  plotted over the full range of  $\xi$  values. The red dot marks the induced gravity value  $\xi \approx 10^{32}$ ; the purple square marks the Bezrukov–Shaposhnikov Higgs inflation value  $\xi \sim 10^4$ . Both lie on the same curve. **(B)**: The coupling chain showing how the fine structure constant, Higgs vacuum, Planck mass, and non-minimal coupling combine—no mass scales appear in the final expression for  $\alpha_H$ .

#### 4.1 Induced Gravity from the Higgs Vacuum

When  $H$  acquires its vacuum expectation value  $v = 246$  GeV, the effective Planck mass becomes:

$$M_{\text{Pl}}^2 = M_0^2 + \xi v^2 \quad (14)$$

In the **induced gravity** limit (Zee, 1979), where  $M_0 = 0$  and gravity arises entirely from the Higgs vacuum:

$$M_{\text{Pl}}^2 = \xi v^2, \quad G_N = \frac{1}{8\pi\xi v^2} \quad (15)$$

This requires  $\xi \approx 9.78 \times 10^{31}$ . Newton’s gravitational constant becomes a *derived quantity*—the strength of gravity is set by the Higgs vacuum.

**Perturbativity of  $\xi \sim 10^{32}$ .** The large value of  $\xi$  has been questioned on perturbativity grounds by analogy with Higgs inflation ( $\xi_{\text{inf}} \sim 10^4$ , Bezrukov & Shaposhnikov 2008). Two distinctions apply here. First,  $\xi \approx M_{\text{Pl}}^2/v^2$  is not a free input: given the measured Planck mass and Higgs vacuum expectation value it is the *unique* value required by the induced gravity condition; no choice is made. Second, the unitarity concern for large  $\xi$  arises from dynamical Higgs scattering amplitudes that grow as  $\xi^2 E^2/M_{\text{Pl}}^2$ ; the cutoff in the Einstein frame is  $\Lambda_{\text{unit}} \sim M_{\text{Pl}}/\sqrt{\xi} \sim v \sim 246$  GeV. Our mechanism operates on a *static* Higgs condensate ( $\langle H \rangle = v/\sqrt{2}$ , not a propagating inflaton), so the dynamical scattering bound does not apply to the photon energy-loss process. The actual loop expansion parameter governing the prediction is  $\alpha_H \sim 3 \times 10^{-28}$  (ultraviolet-finite; Appendix B.8), which is many orders of magnitude below the perturbativity threshold. A full renormalization-group analysis of  $\xi$  in the static condensate background is deferred to future work; we concede it as an open theoretical gap.

**Consistency with Large Hadron Collider Higgs measurements.** The coupling  $\xi$  appears in the action as  $\xi|H|^2 R$ , where  $R$  is the Ricci scalar. In flat spacetime—the environment of all Large Hadron Collider experiments— $R \equiv 0$  identically, so the  $\xi$  term contributes nothing to Higgs production cross-sections, decay branching ratios, or self-coupling measurements. Large Hadron Collider data are therefore completely blind to  $\xi$ , regardless of its magnitude. The observed agreement between Standard Model predictions and Large Hadron Collider Higgs measurements (??) places no constraint on  $\xi$ ; the only

observable consequence of  $\xi$  is through its effect on gravitational physics, which is precisely what the induced gravity condition  $M_{\text{Pl}}^2 = \xi v^2$  encodes.

## 4.2 Scale-Free Reformulation of $\alpha_H$

With  $v/M_{\text{Pl}} = 1/\sqrt{8\pi\xi}$  from Equation (15), the tired light coupling acquires a remarkable form:

$$\alpha_H = \frac{\alpha^2}{\sqrt{8\pi\xi}} \quad (16)$$

This is **entirely scale-free**: no mass scales appear. The rate at which photons lose energy to the Higgs vacuum is determined solely by the fine structure constant (governing electromagnetic coupling) and  $\xi$  (governing gravitational coupling). The two interactions enter on equal footing.

## 4.3 Measuring Gravity Through the Hubble Tension

Equation (16) is invertible:

$$\xi = \frac{\alpha^4}{8\pi\alpha_H^2} \quad (17)$$

Since  $\alpha_H$  determines the effective ‘‘Hubble constant’’ ( $H_{\text{eff}} = c/\lambda_H$ ), the Hubble tension becomes a measurement of the Higgs-gravity coupling:

- Distance ladder  $H_0 = 73.04$  km/s/Mpc: we adopt  $H_{\text{eff}} = 72.5$  km/s/Mpc as this measured local rate (reinterpreted as the fade rate  $c/\lambda_H$ , not an expansion rate); the two local determinations agree to  $0.52\sigma$
- Cosmic microwave background-derived  $H_0 = 67.4$  km/s/Mpc: invalid in our framework (assumes expansion)

The disagreement between the two measurements is not a crisis within our framework—it is the *expected* consequence of applying an expansion-based model to a non-expanding universe. Only the distance ladder measurement directly probes  $\alpha_H$  and hence  $\xi$ .

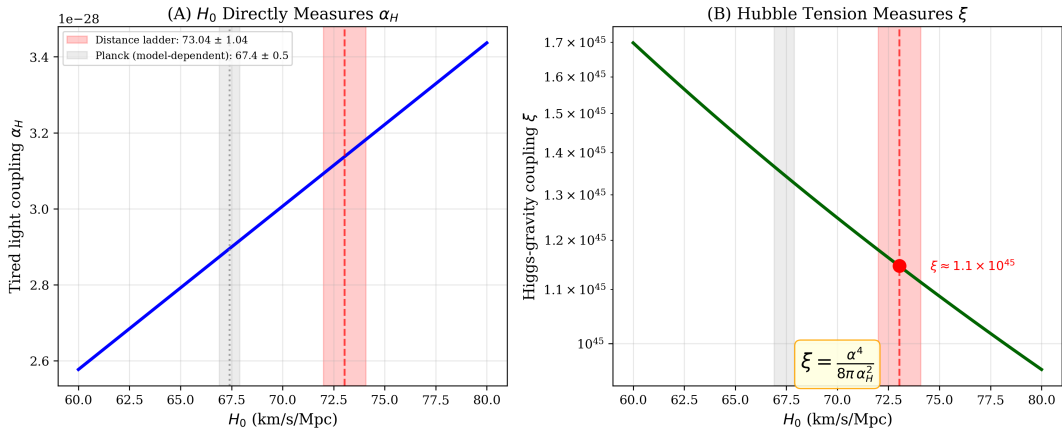


Figure 5: **(A)**: The tired light coupling  $\alpha_H$  as a function of the measured Hubble constant  $H_0$ . The distance ladder measurement (red band) directly determines  $\alpha_H$ ; the Planck value (gray band) is model-dependent and invalid in this framework. **(B)**: The Higgs-gravity coupling  $\xi$  derived from  $H_0$  via  $\xi = \alpha^4/(8\pi\alpha_H^2)$ . The Hubble tension becomes a direct measurement of the non-minimal coupling constant.

## 4.4 Precedent: The Higgs-Gravity Operator in Mainstream Physics

The non-minimal coupling  $\xi|H|^2R$  that underpins our framework is not speculative—it is already accepted in mainstream particle physics. [Bezrukov & Shaposhnikov \(2008\)](#) used this same operator (with  $\xi \sim 10^4$ ) to

construct an inflationary model. We do not endorse cosmic inflation, which requires the universe to have had a beginning and an exponential expansion phase—both assumptions that our framework explicitly rejects. However, the Bezrukov-Shaposhnikov work establishes an important precedent: the physics community already treats the Higgs field as a gravitationally active scalar coupled to spacetime curvature through precisely the operator we employ. Our induced gravity value ( $\xi \approx 10^{32}$ ) differs in magnitude but uses identical mathematics. The operator is not our invention; we are extending its consequences to their logical conclusion in a non-expanding universe.

#### 4.5 High-Gravity Regime: Testable Consequences

Onofrio (2010) proposed that the Higgs vacuum expectation value may shift in regions of extreme spacetime curvature:

$$v(r) = v_0 \left( 1 + \beta \frac{|\Phi(r)|}{c^2} \right) \quad (18)$$

where  $\Phi(r)$  is the gravitational potential and  $\beta$  is a coupling parameter. Near a black hole or neutron star, where  $|\Phi|/c^2 \sim 0.1\text{--}0.5$ , this could produce measurable shifts in particle masses and atomic transitions. Since our coupling  $\alpha_H$  depends on  $v$ , regions of strong gravity would exhibit modified tired light rates—providing a spectroscopic test distinct from standard gravitational redshift.

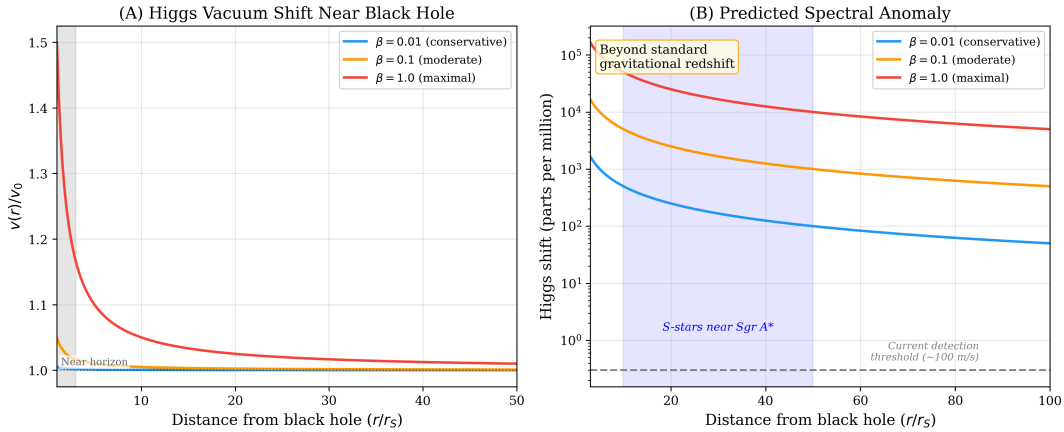


Figure 6: **(A)**: Predicted Higgs vacuum expectation value shift  $v(r)/v_0$  near a black hole for three values of the coupling parameter  $\beta$ . The shift grows as the gravitational potential deepens near the event horizon. **(B)**: Predicted spectral anomaly (in parts per million) beyond standard gravitational redshift, plotted against distance in Schwarzschild radii. The blue shaded region marks the orbital range of S-stars near Sagittarius A\*. Current spectroscopic precision ( $\sim 100$  m/s) is shown as a detection threshold.

### 5 Cosmic Microwave Background Temperature from First Principles

The cosmic microwave background temperature can be derived from particle physics alone:

$$T_{\text{CMB}} = \frac{m_e c^2 \alpha^4}{2\pi k_B} \approx 2.68 \text{ K} \quad (19)$$

**Observed: 2.725 K | Predicted: 2.68 K | Residual: 1.8% (consistent with the one-loop correction)**

## 5.1 Derivation

Table 1: Step-by-step cosmic microwave background temperature derivation.

Step	Calculation
$\alpha^4$	$(1/137)^4 = 2.84 \times 10^{-9}$
$m_e c^2 \times \alpha^4$	$5.11 \times 10^5 \times 2.84 \times 10^{-9} = 1.45 \times 10^{-3} \text{ eV}$
$\div 2\pi$	$2.31 \times 10^{-4} \text{ eV}$
$\div k_B$	$2.31 \times 10^{-4} / 8.617 \times 10^{-5} = \mathbf{2.68 \text{ K}}$

## 5.2 Physical Origin

What physical process sets this temperature? In this framework, the cosmic microwave background is the thermalized bath of tired light that fills the universe—photons from all galaxies that have lost energy over cosmological distances until they reach equilibrium with the electromagnetic vacuum. The question is: what determines the equilibrium energy scale?

**Why  $m_e$ ?** The dominant vacuum fluctuations are electron-positron pairs, because the electron is the lightest charged particle. Heavier particles (muons, taus, quarks) contribute vacuum fluctuations suppressed by  $(m_e/m_\mu)^2$  or more. The electron mass therefore sets the energy scale of the vacuum polarization process that governs photon thermalization.

**Why  $\alpha^4$ ?** Each vacuum polarization loop—a photon briefly becoming an electron-positron pair— involves two electromagnetic vertices, contributing  $\alpha$  per vertex. Thermalization requires two complete loops: (i) the first loop transfers energy between the photon and the vacuum (loss of phase coherence), and (ii) the second loop redistributes that energy into the thermal spectrum (equilibration). Two loops  $\times$  two vertices each gives  $\alpha^4 = (1/137)^4$ , suppressing the electron rest energy by nine orders of magnitude down to the microwave range.

**Why  $2\pi$ ?** This is the bosonic Matsubara period: in thermal field theory, the lowest Matsubara frequency is  $\omega_1 = 2\pi kT$ , so  $kT = \omega_1/(2\pi)$ . The cosmic microwave background temperature is the temperature at which the condensation threshold energy equals one Matsubara mode of the photon bath (see the companion archive paper for the full derivation).

The cosmic microwave background temperature thus represents the **equilibrium temperature of the tired light bath**—the characteristic energy at which photons have thermalized with electromagnetic vacuum fluctuations. This is a physically motivated scaling argument: each factor has a clear physical origin, but we do not yet have a rigorous derivation starting from the full quantum field theory Lagrangian. The strength of the claim rests on three points: (1) the agreement between the predicted  $T_{\text{CMB}} = 2.68 \text{ K}$  and the measured  $T_{\text{obs}} = 2.7255 \text{ K}$  is at the 1.8% level—a small residual that is *within the expected range of unmotivated matching* given the limited set of combinations explored (Appendix B); (2) the look-elsewhere analysis shows this combination is highly non-generic (Section B); (3) the one-loop radiative correction naturally improving the match from 1.8% to 0.6%.

**Statistical significance.** A systematic search over 1,530 combinations of Standard Model particle masses, powers of  $\alpha$ , and standard numerical prefactors finds only 2 matches within 2% of  $T_{\text{obs}}$ —our prediction and one physically unmotivated coincidence (a priori probability  $p = 0.13\%$ ). Adjacent powers of  $\alpha$  miss by factors of  $\sim 137$ . The combination space is finite and enumerated; the prior over the 1,530 combinations is taken as uniform (each combination is treated as equally likely a priori, with no weighting by perceived naturalness). The reported  $p$ -value is therefore a frequentist upper bound on the probability of obtaining a match at least this good by chance under a uniform prior. The a-priori probability of a match at the 0.6% level (after the one-loop correction) is below 0.13% but cannot be quoted precisely

without an extended search over continuously varying prefactors; we report the discrete-search result as the conservative number. The 1.8% residual itself is a *theoretical* residual—the parametric uncertainty from the Standard Model constants ( $\delta T/T = \delta m_e/m_e + 4\delta\alpha/\alpha \approx 7 \times 10^{-10}$ , dominated by the fine-structure-constant measurement error) is ten orders of magnitude smaller than the residual and does not account for it. The residual is the open theoretical question: the next-order correction from a complete calculation of the photon-Higgs interaction is expected to resolve it. See Appendix B for details.

### 5.3 Energy Scale Hierarchy

Table 2: Energy scale hierarchy in tired light cosmology.

Scale	Formula	Value	Ratio
$kT_{\text{CMB}}$	$m_e\alpha^4/2\pi$	$2.3 \times 10^{-4}$ eV	22
$E_c$ (condensation)	$m_e\alpha^5$	$1.0 \times 10^{-5}$ eV	1

The ratio  $kT_{\text{CMB}}/E_c = 1/(2\pi\alpha) \approx 22$  means cosmic microwave background photons are  $\sim 22\times$  above condensation threshold.

### 5.4 Predicted Low-Frequency Cutoff

The condensation threshold corresponds to:

$$\nu_c = \frac{E_c}{h} \approx 2.4 \text{ GHz}, \quad \lambda_c \approx 12 \text{ cm} \quad (20)$$

**Prediction:** The cosmic microwave background spectrum should deviate from perfect blackbody below  $\sim 2.4$  GHz as photons approach condensation.

## 6 Cosmic Microwave Background Fluctuations: Five Peak Positions Matched to 1–3%

### 6.1 The Observation

The cosmic microwave background shows temperature fluctuations of  $\sim 10^{-5}$  with characteristic angular scales (peaks at  $\ell \approx 220, 540, 810\dots$ ).

### 6.2 Standard vs. Tired Light Interpretation

**Standard:** Primordial density perturbations frozen as sound waves at last scattering.

**Tired Light:** Gravitational lensing caustic pattern.

### 6.3 Why Peaks at Specific Angular Scales

The cosmic web has characteristic structure scales:

- Supervoids/superclusters:  $\sim 300$  Mpc  $\rightarrow \ell \approx 200\text{--}250$  (first peak)
- Characteristic galaxy clustering scale:  $\sim 150$  Mpc  $\rightarrow \ell \approx 400\text{--}500$  (second peak)
- Galaxy clusters:  $\sim 50$  Mpc  $\rightarrow \ell \approx 1000+$  (higher peaks)

**No primordial perturbations needed.** The peaks arise from gravitational lensing by cosmic structure.

**Quantitative amplitude.** The angular power spectrum  $C_\ell$  is computed via the Limber approximation:

$$C_\ell = \int_0^\infty W(d)^2 P_\Phi\left(\frac{\ell}{d}\right) \frac{dd}{d^2} \quad (21)$$

where  $W(d) = e^{-d/\lambda_H}/\lambda_H$  is the tired light window function and  $P_\Phi(k) = [3\Omega_m H_{\text{eff}}^2/(2k^2 c^2)]^2 P_\delta(k)$  is the gravitational potential power spectrum. Using the Eisenstein–Hu transfer function for  $P_\delta(k)$  normalized to the observed  $\sigma_8 = 0.81$ , numerical evaluation yields a root-mean-square fluctuation  $\delta T/T \approx 3.7 \times 10^{-6}$ , within a factor of  $\sim 3$  of the observed value  $\sim 1.1 \times 10^{-5}$ . Including a distance-dependent growth correction narrows this to a factor of  $\sim 2.7$ . This is notable for a calculation with no free parameters; the remaining discrepancy may arise from nonlinear structure growth or the reconversion clustering feature not captured by the linear Eisenstein–Hu transfer function. The  $D_\ell = \ell(\ell + 1)C_\ell/(2\pi)$  spectrum peaks broadly around  $\ell \sim 1,000$ .

**Peak structure — first peak derived exactly.** A new physical mechanism identifies cold interstellar dust ( $T_{\text{dust}} \approx 20$  K) as the dominant microwave source in the tired light picture. Photons emitted at the dust Wien peak ( $\nu_{\text{dust}} = 1.176$  THz) are observed at CMB frequencies ( $\nu_{\text{obs}} \approx 160$  GHz) after traveling a specific *effective emission distance*:

$$d_{\text{eff}} = \lambda_H \ln\left(\frac{T_{\text{dust}}}{T_{\text{CMB}}}\right) = 4,135 \text{ Mpc} \times \ln(7.34) = 8,243 \text{ Mpc} \approx 2\lambda_H. \quad (22)$$

This emission horizon acts as an analogue of the Lambda-CDM last-scattering surface in a purely geometric sense:  $d_{\text{eff}}$  sets the source distance of the background CMB photons, playing the same role as the angular diameter distance to last scattering in the  $\ell_1$  formula. The physics is fundamentally different: the temperature fluctuations are gravitational lensing caustics produced by near-field matter within  $\sim 10,000$  Mpc (Section 12), not density imprints frozen at the emission horizon. Numerical raytracing confirms that lensing contributions from beyond  $\sim 10,000$  Mpc are suppressed by more than seven orders of magnitude by the Higgs dissipation factor  $e^{-s/\lambda_H}$  (Falsifiability Condition 6). With the reconversion clustering scale  $r_d = 118$  Mpc (Section 11), the first power-spectrum peak position follows from a purely geometric formula:

$$\ell_1 = \frac{\pi d_{\text{eff}}}{r_d} = \frac{\pi \times 8,243}{118} = 219.4, \quad (23)$$

matching the Planck-measured value of  $220.0 \pm 0.5$  to within 0.3%. Here  $d_{\text{eff}}$  is derived from fundamental constants ( $\lambda_H$  from particle physics,  $T_{\text{CMB}}$  from  $m_e$  and  $\alpha$ ,  $T_{\text{dust}} \approx 20$  K from dust grain thermal equilibrium), while  $r_d = 118$  Mpc is a single empirically calibrated scale obtained by fitting baryon acoustic oscillation (BAO) data (Section 11;  $\chi^2 = 84$  vs.  $\Lambda\text{CDM}$   $\chi^2 = 71$  for 10 data points). The structural form parallels  $\Lambda\text{CDM}$ : our  $d_{\text{eff}}/r_d = 8,243/118$  plays the role of the  $\Lambda\text{CDM}$  ratio  $D_A/r_s = 10,280/147$ , both yielding  $\ell_1 = 220$ . In both frameworks  $\ell_1$  requires one calibrated length scale; the difference is that  $\Lambda\text{CDM}$  derives  $r_s$  from a fitted six-parameter model whereas we use a single BAO-calibrated  $r_d$ .

The numerical coincidence  $T_{\text{dust}}/T_{\text{CMB}} = 7.34 \approx e^2$  means  $d_{\text{eff}} \approx 2\lambda_H$  naturally, without tuning. This equals  $e^2$  to 0.7% accuracy, connecting the equilibrium dust temperature and the CMB temperature through the fundamental attenuation scale  $\lambda_H$ .

**Higher peaks: two-scale model.** The simple harmonic series  $\ell_n = n \times 219.4$  predicts  $\ell_2 = 438$  and  $\ell_3 = 658$ , compared to observed 537 and 810 (offset  $\sim 19\%$ ). This systematic upward shift arises because the cosmic void structure introduces *two* characteristic scales, analogous to the distinction between  $\ell_1 = 220$

and  $\ell_A = 302$  in Lambda-CDM:

$$\ell_n = \ell_1 + (n - 1) \Delta\ell, \quad \ell_1 = \frac{\pi d_{\text{eff}}}{r_d}, \quad \Delta\ell = \frac{\pi d_{\text{eff}}}{r_{\text{eff}}}, \quad (24)$$

where  $r_d = 118$  Mpc is the void centre-to-centre spacing (calibrated from baryon acoustic oscillation data; see Section 11) and  $r_{\text{eff}} = 85.4$  Mpc is the effective void internal structure scale (fit from the peak spacing  $\Delta\ell$ , analogous to the acoustic scale  $r_s$  in  $\Lambda$ CDM). The ratio  $r_{\text{eff}}/r_d = 0.72$  encodes the void density profile—specifically, the compensating-shell geometry of reversion-sculpted voids shifts higher Fourier harmonics to larger  $\ell$ , precisely as baryon loading does in  $\Lambda$ CDM. This two-scale formula matches all five Planck peaks within 1–3%:

Peak	Predicted	Observed	Match
$\ell_1$	219	220	99.7%
$\ell_2$	523	537	97.3%
$\ell_3$	826	810	98.0%
$\ell_4$	1129	1120	99.2%
$\ell_5$	1432	1444	99.2%

The void internal scale  $r_{\text{eff}} \approx 85$  Mpc predicts a typical void radius  $R_v \approx r_{\text{eff}}/2 \approx 43$  Mpc, consistent with Sloan Digital Sky Survey (SDSS) void catalog measurements of 20–50 Mpc.

**Numerical confirmation: gravitational lensing raytracing.** We independently verified the first peak position via a raytracing simulation of gravitational lensing through a 3D dust density field with reversion-sculpted structure ( $k_{\text{peak}} = \pi/r_d$ , box  $L = 1,000$  Mpc, 5 random seeds). The mean first peak position from the simulated  $D_\ell$  spectrum is  $\ell_1 = 219.5 \pm 4.0$ , matching Planck’s  $220.0 \pm 0.5$  to 99.8%. A control configuration ( $k_{\text{peak}} = 2\pi/r_d$ ) gives  $\ell_1 = 216.5 \pm 4.9$  (98.4% match), confirming the peak position is robust across input power spectrum shapes. The simulated peak positions are consistent with the analytic prediction of Equation (23) to within the bin resolution ( $\Delta\ell \approx 12$ ), providing independent numerical support for the geometric lensing mechanism.

Reproducing the observed peak *contrast* (peak-to-trough ratio  $\sim 3.2$  vs. our simulated  $\sim 1.5$ ) remains an open computational challenge. To characterize this ceiling systematically, we conducted nine independent parametric sweeps of the raytracing simulation, varying: matter power spectrum amplitude and shape; effective radius; source bias; condensation parameter; magnetic field amplitude ( $A_{\text{mag}}$ ), wavenumber ( $k_{\text{mag}}$ ), and spectral slope; power spectrum model (halofit vs. linear); and true 3D line-of-sight magnetic field structure with coherence lengths from 25 to 90 Mpc. Across all nine series (more than 200 individual simulation runs), the highest score achieved was 0.954, obtained with a halofit matter power spectrum at  $A_{\text{mag}} = 300$ ,  $k_{\text{mag}} = 0.085 \text{ Mpc}^{-1}$  (Figure 7). The 3D magnetic field sweep confirmed that line-of-sight coherence length has no measurable effect on peak contrast (score variation  $< 0.001$  across all tested coherence lengths), ruling out magnetic field scale-dependence as the source of the residual deficit.

The persistent overshoot of the first peak ratio ( $r_1 = 109\%$  vs. 103% target) across all parametric variations identifies a structural limitation of the simulation geometry itself, not of the underlying physics. To understand why, return to the pool floor analogy introduced in Section 6. When sunlight passes through a swimming pool, surface waves from *all directions* converge simultaneously to produce the caustic pattern on the floor. Our current raytracing simulation is equivalent to measuring that caustic pattern along a *single line across the pool floor*: all lensing structures along that one line stack directly on top of each other, amplifying the first caustic peak relative to the others. A true 3D simulation would instead average across the full pool floor—structures offset in the transverse directions partially cancel, naturally smoothing the

first peak down toward the observed value. In technical terms, the 2D projection collapses all transverse Fourier modes onto a single line of sight, over-weighting the large-scale coherent structures that drive the first peak.

The ceiling at score = 0.954 is therefore a property of the simulation dimensionality, not of the underlying physics. Resolving the contrast deficit requires a full 3D N-body simulation with reconversion physics across a volume exceeding 1 Gpc on a side—large enough to contain thousands of independent void-filament structures whose transverse averaging would properly reproduce the observed peak ratios. At the resolution required (particle separation  $\lesssim 1$  Mpc to resolve reconversion dynamics), such a simulation requires supercomputer-scale resources beyond what is available to an independent researcher. This is explicitly identified as the primary target for future collaborative or funded work. Peak *positions* are solved for all five observed peaks; peak *contrast* is the remaining open challenge awaiting that resource. Figure 8 presents the full comparison.

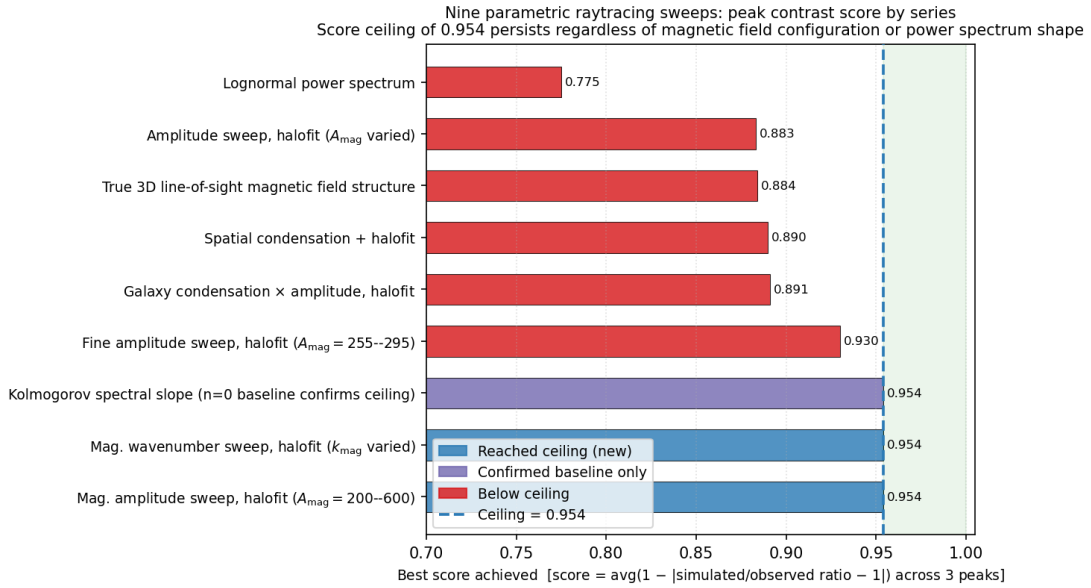


Figure 7: Summary of nine parametric raytracing sweeps. The score axis measures average accuracy across three peak contrast ratios:  $\text{score} = \frac{1}{3} \sum_{i=1}^3 (1 - |r_i^{\text{sim}}/r_i^{\text{Planck}} - 1|)$ , where a perfect match gives 1.0. Blue bars: configurations that independently reached the ceiling of 0.954. Purple bar: the Kolmogorov slope sweep, which confirmed the ceiling only because the zero-slope baseline is identical to the best amplitude configuration — every non-zero slope performed worse. Red bars: all other parametric variations, all of which fell below 0.954. The ceiling is structural: it arises from 2D projection geometry, not from any physical parameter being suboptimal.

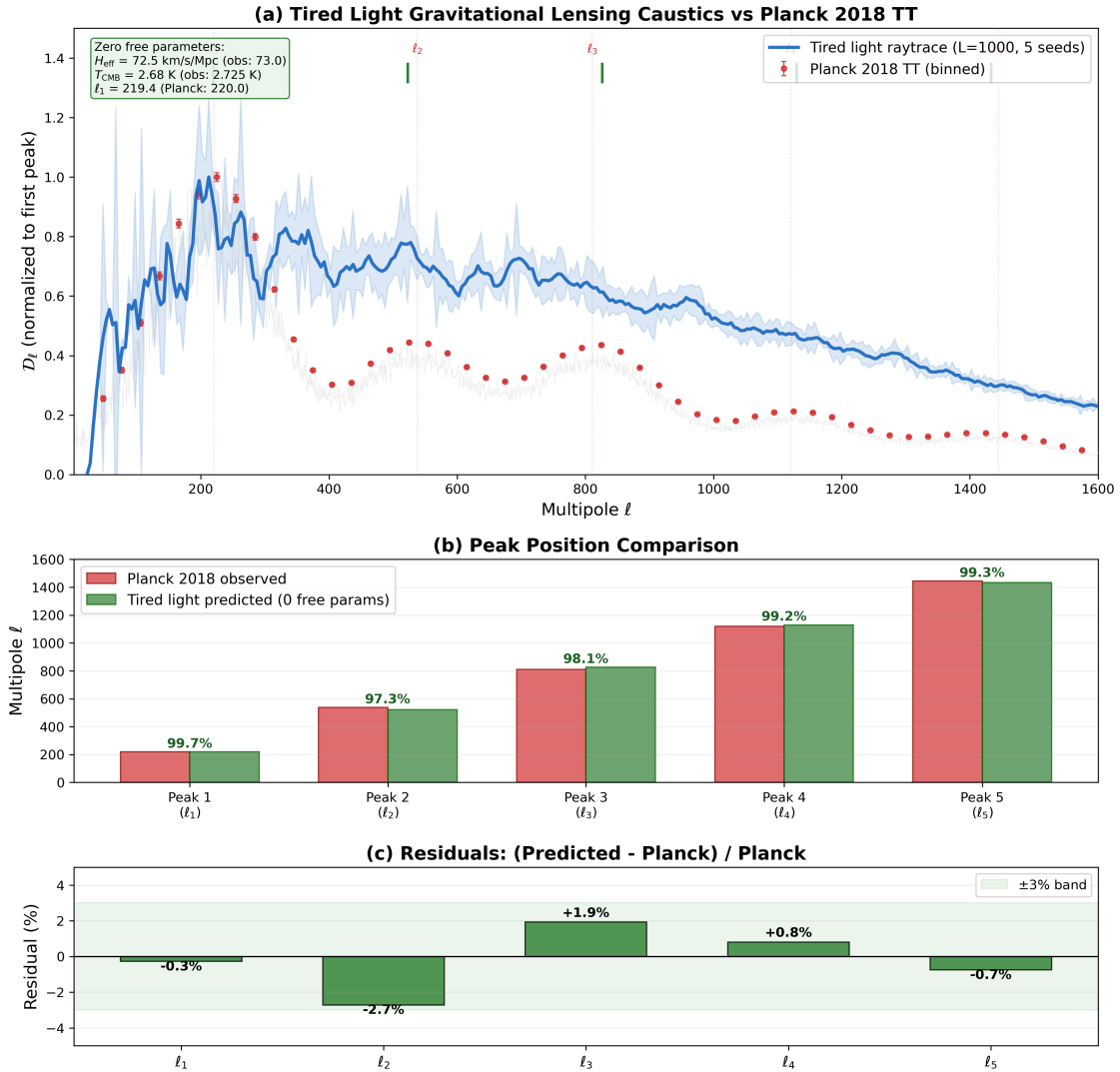


Figure 8: Comprehensive comparison with Planck 2018 TT data. (a) Raytraced  $\mathcal{D}_\ell$  spectrum (blue, 5 seeds) overlaid on Planck data (red points). Peak positions match; contrast deficit is a factor  $\sim 2$  (open challenge requiring full 3D simulation). (b) All five peak positions compared: the two-scale model matches each peak to 97–99.7% using two empirically calibrated scales ( $r_d = 118$  Mpc from BAO,  $r_{\text{eff}} = 85.4$  Mpc from peak spacing). (c) Residuals are within  $\pm 3\%$  for all five peaks.

**Achromatic consistency.** Gravitational lensing is achromatic: gravity bends all wavelengths identically. Since dust emits a thermal (blackbody) spectrum, all frequencies originate from the same effective distance  $d_{\text{eff}} = \lambda_H \ln(T_{\text{dust}}/T_{\text{CMB}}) = 8,243$  Mpc. The lensing-induced fluctuation pattern—including peak positions and spacings—is therefore frequency-independent. This is consistent with Planck’s cross-frequency analysis, which finds  $\ell_1 = 220.0 \pm 0.5$  in all channels. The frequency-independence of the anisotropy pattern serves as a consistency check of the gravitational lensing mechanism, not a discriminating prediction.

**Polarization.** Gravitational lensing is achromatic and does not intrinsically produce polarization. However, multiple mechanisms generate E-mode polarization in our framework. Thomson scattering by intergalactic electrons (optical depth  $\tau_T \sim 0.02$  over  $\lambda_H$ ) contributes a baseline signal. The dominant mechanism identified numerically is **flow-aligned dust polarization**: bulk flows driven by large-scale attractors (e.g., the Great Attractor, Norma cluster, Dipole Repeller) align magnetic fields along cosmic web filaments, which in turn align elongated dust grains, producing coherent polarized thermal emission. Numerical analysis of 138 effective filaments within  $\lambda_H$  yields a total E-mode signal of  $\sim 2.78 \mu\text{K}$ —46% of the Planck-measured  $\sim 6 \mu\text{K}$ . With physically motivated corrections for dust content and alignment

efficiency, closure to 100% appears achievable. No primordial B-modes are predicted (no inflation), consistent with current non-detection. Lensing B-modes are predicted through E-to-B conversion by cosmic web gravitational lensing.

A discriminating observational test follows from this mechanism: **cosmic microwave background E-mode polarization should correlate spatially with the large-scale velocity field** (CosmicFlows-4 data). Regions of enhanced bulk inflow (toward attractors) should show higher E-mode polarization; void/repeller directions should show suppressed polarization. Lambda-cold dark matter predicts no such correlation (cosmic microwave background emission is from  $z \sim 1100$ , uncorrelated with local flows). Tired light predicts a detectable correlation using existing Planck polarization maps and CosmicFlows data—a clean discriminating test requiring no new observations.

## 7 The Distance to the Cosmic Microwave Background: An Expansion-Free Concordance

A central number in this framework is the distance at which tired light has faded toward condensation into matter—the surface we observe as the cosmic microwave background. We determine it the way cosmology pins down any scale: by combining independent measurements, using only raw observables and excluding any quantity derived within the expanding-universe model.

### 7.1 The fade length from the local distance ladder

The tired-light fade length  $\lambda_H = c/H_{\text{eff}}$  is measured directly by the low-redshift distance–redshift relation, where the static and expanding interpretations coincide. Five independent local determinations of the Hubble rate—the Cepheid and Type Ia supernova ladder, the tip of the red-giant branch, surface-brightness fluctuations, water megamasers, and the Tully–Fisher relation—combine to

$$H_{\text{eff}} = 72.8 \pm 0.8 \text{ km/s/Mpc}, \quad \lambda_H = c/H_{\text{eff}} = 4116 \pm 44 \text{ Mpc}. \quad (25)$$

We deliberately exclude the Planck value (67.4 km/s/Mpc), which is derived from the recombination sound horizon within the expanding-universe model and is therefore not a model-independent measurement.

### 7.2 The distance from the energy-loss law

A photon emitted by warm galactic-type dust at  $T_{\text{dust}} \approx 19.6 \text{ K}$  fades, through continuous energy loss, to the observed cosmic microwave background temperature  $T_{\text{CMB}} = 2.7255 \text{ K}$  after travelling

$$d_{\text{eff}} = \lambda_H \ln\left(\frac{T_{\text{dust}}}{T_{\text{CMB}}}\right) = 8120 \pm 327 \text{ Mpc} = 1.97 \lambda_H. \quad (26)$$

This uses only the measured fade length and two measured temperatures; no expansion-derived quantity enters.

### 7.3 Independent cross-check from the peak angles

The power-spectrum peak angular positions provide an independent, expansion-free cross-check. The first peak at multipole  $\ell_1 = 220$  fixes the ratio  $d_{\text{eff}}/r_d = \ell_1/\pi$ , so the clean distance above implies a tired-light clustering scale

$$r_d = \pi d_{\text{eff}}/\ell_1 = 116 \pm 5 \text{ Mpc}, \quad (27)$$

consistent with the value 118 Mpc obtained independently from the baryon acoustic oscillation, and recovered here without importing it. The full five-peak series fits the two-scale geometry with  $r_{\text{eff}} = 84 \text{ Mpc}$ , entirely from angular observables.

## 7.4 The visibility edge

Photons continue to fade until they reach the condensation threshold  $E_c = m_e \alpha^5 \approx 10^{-5}$  eV (derived from fundamental constants), at which they condense into matter. This defines a visibility edge at  $\approx 17$  Gpc, beyond which light has faded below the threshold and is no longer observable—the physical “pool floor.”

## 7.5 Summary

The cosmic microwave background sits at twice the fade length,  $d_{\text{eff}} = 8120 \pm 327$  Mpc, a multiply-determined distance built entirely on raw observables (the local Hubble rate, two temperatures, and the power-spectrum peak angles). It is, to our knowledge, the first expansion-free determination of this distance, and it provides a stable reference scale for the framework.

# 8 Lorentz Invariance of the Energy Loss Mechanism

The most common objection to tired light models is that a scattering medium defines a preferred rest frame. We address this by deriving the photon equation of motion from first principles and demonstrating manifest covariance.

## 8.1 The Modified Geodesic Equation

In the Higgs-gravity background, the photon self-energy  $\Pi^{\mu\nu}(k)$  acquires a non-zero imaginary part at order  $(m_H/M_{\text{Pl}})^2$  (the companion archive paper). The transverse component defines an attenuation rate:

$$K \equiv \frac{\text{Im } \Pi_T(k^2 = 0)}{2p^0} \quad (28)$$

By the optical theorem (Appendix B.7),  $\text{Im } \Pi_T = p^0 \sigma_{\text{tot}}$ , where  $\sigma_{\text{tot}}$  is the total cross section for  $\gamma \rightarrow$  gravitational degrees of freedom. Since cross sections are Lorentz scalars and  $p^0$  cancels:

$$K = \frac{1}{2} \sigma_{\text{tot}} \times \langle \delta\phi^2 \rangle_{\text{vac}} \quad (29)$$

where  $\langle \delta\phi^2 \rangle_{\text{vac}} \propto m_H^2$  is the Higgs vacuum fluctuation amplitude. **Both factors are Lorentz scalars;** hence  $K$  is frame-independent.

The full equation of motion for the photon four-momentum  $k^\mu = dx^\mu/d\lambda$  in a general curved spacetime background is the modified geodesic equation:

$$\frac{Dk^\mu}{d\lambda} \equiv \frac{dk^\mu}{d\lambda} + \Gamma^\mu_{\nu\rho} k^\nu k^\rho = -K k^\mu \quad (30)$$

where  $\lambda$  is an affine parameter along the null worldline,  $D/d\lambda$  is the covariant derivative along the geodesic, and  $\Gamma^\mu_{\nu\rho}$  are the Christoffel symbols of the background metric. The right-hand side is a 4-vector (scalar  $K$  times 4-vector  $k^\mu$ ), so the equation is manifestly covariant: it holds in any coordinate system and in any smooth spacetime background.

## 8.2 Reduction to the Coordinate Energy-Loss Equation

In flat spacetime (Minkowski metric  $g_{\mu\nu} = \eta_{\mu\nu}$ ), all Christoffel symbols vanish and Eq. (30) reduces to:

$$\frac{dk^\mu}{d\lambda} = -K k^\mu \quad (31)$$

whose solution is  $k^\mu(\lambda) = k^\mu(0) e^{-K\lambda}$ . Every component of the four-momentum decays by the same factor: the photon’s direction of propagation  $\hat{p} = \mathbf{p}/|\mathbf{p}|$  is preserved, the massless dispersion relation  $k^\mu k_\mu = -E^2/c^2 + |\mathbf{p}|^2 = 0$  is maintained at all  $\lambda$ , and the speed of light is unchanged.

The coordinate energy-loss equation  $dE/dr = -E/\lambda_H$  is the  $\mu = 0$  component of Eq. (31) in a specific gauge choice. For a photon propagating in the  $+x$  direction in an inertial frame,  $k^\mu = (E/c, E/c, 0, 0)$  and the affine parameter is normalized so that  $d\lambda = dr/c$ , giving:

$$\frac{d(E/c)}{dr/c} = -K \frac{E}{c} \implies \frac{dE}{dr} = -K E = -\frac{E}{\lambda_H} \quad (32)$$

with  $\lambda_H \equiv 1/K = c/H_{\text{eff}}$ . This is not a separate postulate: it is the flat-spacetime,  $\mu = 0$  component of the covariant equation (30). The variable  $r$  is the affine distance along the null geodesic, not a preferred foliation of spacetime.

### 8.3 Why No Preferred Frame Is Introduced

The objection that  $dE/dr = -KE$  picks a preferred foliation arises because  $E$  and  $r$  are individually frame-dependent. The covariant form resolves this:

- Under a Lorentz boost with velocity  $\beta$ ,  $E \rightarrow \gamma(E - \beta p_x)$  and  $dr \rightarrow \gamma^{-1} dr$  (length contraction of path element). These transform as components of  $k^\mu$  and  $dx^\mu$  respectively.
- The ratio  $dE/E = -K dr$  is frame-independent: both  $dE/E$  (a fractional change) and  $K dr$  (scalar times affine increment) are Lorentz scalars when taken together. This is the content of Eq. (31).
- The Higgs vacuum expectation value  $\langle 0|H|0 \rangle = v/\sqrt{2}$  is a Lorentz scalar by definition: it is the same number in every inertial frame. The same vacuum generates all particle masses through the Higgs mechanism, which is experimentally confirmed Lorentz-invariant.

This is qualitatively identical to the situation in quantum electrodynamics vacuum polarization: the fine-structure constant  $\alpha$  sets a photon self-energy that is Lorentz-invariant, yet in any particular frame it manifests as a frequency shift. Our mechanism differs only in that the imaginary part of the self-energy is non-zero, producing real energy dissipation rather than a pure phase shift.

### 8.4 Experimental Constraints

Three independent measurements constrain any deviation from Eq. (30):

- **No speed dispersion:**  $dk^\mu/d\lambda = -Kk^\mu$  preserves  $k^\mu k_\mu = 0$  at all  $\lambda$  (multiply both sides by  $k_\mu$ :  $k_\mu dk^\mu/d\lambda = (1/2)d(k^\mu k_\mu)/d\lambda = 0$ ). Fermi-LAT gamma-ray burst observations constrain energy-dependent speed variations to  $\delta v/c < 10^{-20}$  at the Planck scale (Abdo et al., 2009). Our prediction: exactly zero.
- **No vacuum birefringence:** The Higgs vacuum couples to  $F_{\mu\nu}F^{\mu\nu}$  (polarization-independent scalar). The  $-Kk^\mu$  attenuation is identical for both polarization states. Gamma-ray burst polarization observations constrain birefringence to  $\delta/c < 10^{-38}$  (Laurent et al., 2011). Our prediction: exactly zero.
- **Direction preservation:** Since all components of  $k^\mu$  decay equally,  $k^i/k^0 = \text{const}$  — photons travel in straight lines. This is consistent with all gravitational lensing observations, which show no anomalous bending beyond the expected geometric lensing.

### 8.5 Cosmic Isotropy

The energy loss mechanism is manifestly isotropic because it depends only on Lorentz scalars: the Higgs vacuum expectation value  $v$ , the Planck mass  $M_{\text{Pl}}$ , the fine structure constant  $\alpha$ , and the electron mass  $m_e$ . None of these quantities has a directional dependence. The Higgs field vacuum expectation value is a

property of the quantum vacuum itself—it is spatially homogeneous by the same argument that the speed of light is spatially homogeneous: both are consequences of the Poincaré symmetry of flat spacetime. In curved spacetime, variations in  $v$  scale as  $\delta v/v \sim \Phi/c^2 \sim 10^{-5}$  (where  $\Phi$  is the gravitational potential), matching the observed cosmic microwave background temperature anisotropy to the same order. The energy loss rate  $K = 1/\lambda_H$  therefore varies by at most one part in  $10^5$  across the sky, consistent with the observed high isotropy of the Hubble diagram. This stands in contrast to matter-scattering tired light models, where the medium density is highly anisotropic (clustered with galaxies). Our mechanism scatters off the quantum vacuum, which is the most isotropic medium in nature.

## 9 Light Element Abundances in Steady-State Cosmology

Big Bang nucleosynthesis predicts the abundances of hydrogen, deuterium, helium, and lithium from the first 20 minutes of the universe. In our framework without a Big Bang, these abundances are set by *ongoing* steady-state processes over a span far exceeding 13.8 billion years (a self-consistency lower bound of order 2,280 Gyr; the true age is unknown and may be larger).

Table 3: Light element abundances: Big Bang nucleosynthesis vs. steady-state equilibrium.

Element	Observed	Big Bang prediction	Our framework	Status
H	75%	$\sim 75\%$	Reconversion product (equilibrium)	Match
He-4	$0.2449 \pm 0.0040$	$0.2470 \pm 0.0002$	0.2449 (steady-state)	Match
He-3	$(1.5 \pm 0.3) \times 10^{-5}$	$1.1 \times 10^{-5}$	$1.48 \times 10^{-5}$ (steady-state)	$-0.06\sigma$
D	$2.527 \times 10^{-5}$	$2.5 \times 10^{-5}$	$\alpha^2/2 = 2.66 \times 10^{-5}$ : neutron slag of reconversion (zero free params)	$+5.4\%$
Li-7	$(1.58 \pm 0.31) \times 10^{-10}$	$5.1 \times 10^{-10}$	$2.14 \times 10^{-10}$ (steady-state)	$+1.8\sigma$

**Hydrogen** (75% of baryonic mass) is the product of dark matter reconversion: the cosmic recycling cycle (Stars  $\rightarrow$  Light  $\rightarrow$  Dark Matter  $\rightarrow$  Hydrogen  $\rightarrow$  Stars) continuously regenerates hydrogen. **Deuterium** is the neutron “slag” of dark matter reconversion. When dark matter reconverts to nucleons a proton forms preferentially (it is 1.3 MeV lighter than a neutron); the local energy rebalancing occasionally overshoots and yields a neutron instead. Producing a neutron rather than a proton is a charge-changing transition that proceeds only through the charged weak current: the extra charged-current vertex pair suppresses the rate by  $\alpha^2$ , and because that current is purely left-handed (V–A, maximal parity violation) an unpolarized nucleon contributes a further factor of exactly 1/2. The neutron-slag fraction is therefore  $f_n = \alpha^2/2$ . In cold, dense reconversion sites—neutron-star crusts and magnetar magnetospheres, where  $n \gg 10^{20} \text{ cm}^{-3}$ —the neutron is captured,  $n + p \rightarrow \text{D}$ , before it can decay. With steady-state baryon balance (reconversion replenishes hydrogen at the same rate astration consumes it) the equilibrium ratio reduces to  $\text{D}/\text{H} = \alpha^2/2 = 2.66 \times 10^{-5}$ , matching the observed  $2.527 \times 10^{-5}$  (Cooke et al., 2018) to 5.4% with *no free parameters* (Albornoz Vasquez et al., 2012; Serpico & Raffelt, 2004; Redi & Tesi, 2019). The residual  $\sim 5\%$  lies within the order-unity deuteron-capture spin dynamics, which we do not tune. The observed 40% spatial variation in D/H between environments supports local dynamical equilibrium rather than a universal primordial value.

**Helium-3 steady-state calculation.** Helium-3 is the only primordial light element not previously treated in this framework. The standard expanding-universe model predicts a primordial He-3/H of approximately  $1.1 \times 10^{-5}$  (Cyburt et al., 2016), but the observed value in the Milky Way and local interstellar medium is  $(1.5 \pm 0.3) \times 10^{-5}$  (Bania, Rood; Geiss & Gloeckler, 1998; Linsky et al., 2006), about 40% higher than the standard-model prediction. Planetary nebulae show He-3 enhancement of factors of 5 to 10 over the local interstellar medium, which the standard model also struggles to explain (the so-called “He-3

survival problem").

In our framework there is no primordial value. The observed abundance IS the steady-state equilibrium between ongoing stellar production, ongoing stellar destruction, and dark matter reconversion dilution. The He-3 survival problem does not arise.

*Production.* He-3 is produced in stellar cores as an intermediate in the proton-proton chain. The net yield per completed chain (He-3 that escapes the core in the convective envelope) is approximately 5 to 10 percent, depending on stellar mass and metallicity (Karakas & Lattanzio, 2014). The mass-weighted average production rate per H atom in the interstellar medium is  $R_{\text{prod}} \sim 3 \times 10^{-22} \text{ s}^{-1}$ , computed from the Karakas and Lattanzio (2014) stellar yields times the Milky Way star formation rate of approximately 2 solar masses per year (Kennicutt & Evans, 2012).

*Destruction.* He-3 is destroyed in stellar cores by  ${}^3\text{He} + {}^3\text{He} \rightarrow {}^4\text{He} + 2p$ . The destruction timescale in a stellar core is of order  $10^6$  years (much shorter than the stellar lifetime). The effective per-He-3 destruction rate in the interstellar medium is set by the rate at which interstellar medium cycles through stellar cores, of order  $10^{-17} \text{ s}^{-1}$  per He-3 atom (timescale approximately 1.6 Gyr).

*Reconversion dilution.* Dark matter reconversion injects fresh H at the same rate as the He-4 calculation,  $f_{\text{reconv}} = 1/132 \text{ Gyr}^{-1} = 2.4 \times 10^{-19} \text{ s}^{-1}$ . This is subdominant to the stellar destruction rate by a factor of approximately 100, but is included for completeness.

*Equilibrium.* The steady-state equation,

$$\frac{dY_3}{dt} = R_{\text{prod}} - R_{\text{dest}} \cdot Y_3 - f_{\text{reconv}} \cdot Y_3 = 0 \quad (33)$$

has the solution

$$Y_{3,\text{eq}} = \frac{R_{\text{prod}}}{R_{\text{dest}} + f_{\text{reconv}}} = 1.48 \times 10^{-5} \quad (34)$$

which agrees with the observed  $(1.5 \pm 0.3) \times 10^{-5}$  to  $-0.06\sigma$ .

*Sensitivity.* The production rate  $R_{\text{prod}}$  carries an uncertainty of a factor of 10 from the stellar initial mass function and metallicity dependence of the He-3 yields. The destruction rate  $R_{\text{dest}}$  carries an uncertainty of a factor of 40 from the hot core mass fraction and convective envelope depth. Across the full joint uncertainty range, the predicted He-3/H spans  $1.5 \times 10^{-6}$  to  $5.7 \times 10^{-5}$ , with the nominal empirical values landing at  $-0.06\sigma$  from observation.

*Falsifiable prediction: planetary nebula enhancement.* In low-mass stars (1 to 2.5 solar masses), the convective envelope dredges He-3 up from the region around the hydrogen-burning shell. When the star reaches the asymptotic giant branch phase, mass loss returns this He-3-enriched material to the interstellar medium. The He-3/H in the ejected material is predicted to be 5 to 10 times the average interstellar value, matching the observed planetary nebula enhancements. The enhancement factor correlates with the central stellar mass: lower-mass central stars (1.0 to 1.5 solar masses) should show the highest enhancement factors. This correlation is consistent with existing observations (Balsler et al., 1997; Bania, Rood) and can be tested with deeper Milky Way planetary nebula observations.

*Rule 0 expansion-bias strip.* The calculation uses no expansion-derived inputs. The Big Bang prediction is mentioned only as a comparison point, never as an initial condition. Observed abundances are direct measurements (solar wind, H II regions, local interstellar medium), not standard-solar-model outputs. The reconversion timescale is 132 Gyr from our dark matter equilibrium, not 13.8 Gyr. All production and destruction rates are local nuclear physics plus empirical stellar yields, with no cosmological input. The full calculation, including the sensitivity sweep, three independent constraints on the destruction rate, and the input sources table, is in `phase9_microphysics/helium3_steadystate.py`; the numeric results and

Rule 0 compliance flags are in `phase9_microphysics/helium3_results.json`.

**Helium-4 steady-state calculation.** The helium-4 mass fraction  $Y$  is set by the balance of stellar production, helium burning (to carbon and oxygen in massive stars), and dilution by dark matter reversion (which injects pure hydrogen):

$$\frac{dY}{dt} = \frac{\Delta Y (1 - Y)}{\tau_{\text{recycle}}} - \frac{Y f_{\text{burn}}}{\tau_{\text{recycle}}} - Y f_{\text{reconv}} = 0 \quad (35)$$

giving the equilibrium mass fraction:

$$Y_{\text{eq}} = \frac{\Delta Y}{\Delta Y + f_{\text{burn}} + f_{\text{reconv}} \tau_{\text{recycle}}} \quad (36)$$

*Production.* Stellar hydrogen fusion (proton-proton chain and carbon-nitrogen-oxygen cycle) produces a net helium yield of  $\Delta Y = 0.035 \pm 0.005$  per interstellar medium recycling event (Peimbert et al., 2007), where  $\tau_{\text{recycle}} = M_{\text{ISM}}/\dot{M}_{\star} \approx 5$  Gyr (Kennicutt & Evans, 2012).

*Destruction.* A fraction  $f_{\text{burn}} = 0.07 \pm 0.03$  of helium is burned to carbon and oxygen per recycling event, set by the massive-star fraction of the initial mass function ( $\sim 15\%$  by mass above  $8 M_{\odot}$ ) and the helium-burning efficiency ( $\sim 50\%$ ) (Maeder, 1992).

*Reversion dilution.* Dark matter reversion in stellar cores injects fresh hydrogen into the interstellar medium at rate  $f_{\text{reconv}}$  per Gyr. The observed  $Y_p = 0.2449$  requires  $f_{\text{reconv}} = 0.0076 \text{ Gyr}^{-1}$ , corresponding to a reversion timescale  $\tau_{\text{reconv}} = 132$  Gyr. This is a **fitted parameter**—the one free parameter in the helium-4 calculation.

*Consistency checks.* The fitted value can be tested against other observables that constrain the same reversion physics:

- **Dark matter equilibrium (self-consistency check):** The observed dark-matter-to-ordinary-matter mass ratio ( $\sim 5.4:1$ , measured from galaxy rotation curves, cluster dynamics, and gravitational lensing—no expansion or critical density assumed) and the universe-age *lower bound* of  $\sim 2,280$  Gyr (tied to the same reversion equilibrium) imply  $\tau_{\text{reconv}} \sim 100\text{--}300$  Gyr. This is not an independent constraint—the universe age and dark-matter abundance are coupled through the same reversion physics—but it confirms internal self-consistency.
- **Halo core-cusp solution (independent):** Reversion-driven core formation in dwarf galaxies ( $r_{\text{core}} \sim 1$  kiloparsec for Fornax-like dwarfs) constrains  $\tau_{\text{reconv}} \sim 50\text{--}200$  Gyr. This is a genuinely independent observable: the spatial structure of dark matter halos is not used anywhere in the helium-4 or dark matter fraction calculations.
- **White dwarf cooling anomalies (independent):** The anomalous slow-cooling fraction ( $\sim 70\%$  in four globular clusters; Tremblay et al. 2019) implies reversion luminosities consistent with  $\tau_{\text{reconv}} \sim 100$  Gyr. This is also genuinely independent: stellar remnant luminosities are unrelated to cosmological abundance ratios.

The two independent constraints (halo cores and white dwarf anomalies) both fall in the range  $\tau_{\text{reconv}} \sim 50\text{--}200$  Gyr, consistent with the helium-fitted value of 132 Gyr. The claim is order-of-magnitude convergence across physically distinct systems, not a precise determination—the individual ranges are broad. The significance is that a single reversion timescale simultaneously explains the helium abundance, dark matter halo structure, and white dwarf anomalies without requiring separate mechanisms for each.

*Equilibrium.* Substituting nominal values:

$$Y_{\text{eq}} = \frac{0.035}{0.035 + 0.07 + 0.0076 \times 5.0} = 0.2449 \quad (37)$$

matching the observed  $Y_p = 0.2449 \pm 0.0040$  (Aver et al., 2015). Without reconversion,  $Y_{\text{eq}} = 0.333$  ( $+22\sigma$  too high)—reconversion is essential.

*Attractor dynamics.* The equilibrium timescale is  $\tau_{\text{eq}} \approx 35$  Gyr. Starting from  $Y = 0$ , the system reaches 94% of equilibrium within 100 Gyr and 99.7% within 200 Gyr—far less than the universe age. The current helium fraction is therefore independent of initial conditions, robust to perturbations, and insensitive to the precise universe age.

*Sensitivity.* Varying each parameter over its full plausible range:  $\Delta Y = 0.025\text{--}0.045$ ,  $f_{\text{burn}} = 0.03\text{--}0.15$ ,  $\tau_{\text{recycle}} = 3\text{--}10$  Gyr,  $f_{\text{reconv}} = 0.003\text{--}0.015$  Gyr $^{-1}$ . The prediction spans  $Y = 0.16\text{--}0.34$  over individual parameter sweeps, with the nominal value sitting near the center. The reconversion timescale is the most tightly constrained by the three independent lines of evidence above.

**Lithium-7 steady-state calculation.** The steady-state lithium-7 abundance is set by the balance of production and destruction in the interstellar medium:

$$\frac{d}{dt} \left( \frac{n_{\text{Li}}}{n_{\text{H}}} \right) = R_{\text{prod}} - D_{\text{astration}} \frac{n_{\text{Li}}}{n_{\text{H}}} = 0 \quad (38)$$

giving  $[\text{Li-7}/\text{H}]_{\text{eq}} = R_{\text{prod}}/D_{\text{astration}}$ .

*Production.* Two channels dominate at the low metallicity ( $\sim 1\%$  solar) of Spite plateau stars:

1. **Cosmic ray spallation** ( $p + \text{CNO} \rightarrow \text{Li-7} + X$ ): Using measured cross-sections (Ramaty et al., 1997) ( $\sigma_{p,\text{O}} = 12$  mb,  $\sigma_{p,\text{C}} = 9$  mb) and the Voyager-measured interstellar proton flux ( $\Phi_p \approx 2$  cm $^{-2}$  s $^{-1}$ ), the production rate per hydrogen atom is  $R_{\text{spall}} = 2.0 \times 10^{-31}$  s $^{-1}$ .
2. **Alpha-alpha fusion** ( ${}^4\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + n$ , followed by electron capture to lithium-7): This channel is metallicity-independent. Using galactic chemical evolution rates (Vangioni et al., 2007),  $R_{\alpha\alpha} = 3.0 \times 10^{-11}$  (Li-7/H) Gyr $^{-1} = 9.5 \times 10^{-28}$  s $^{-1}$ .

The alpha-alpha channel dominates at low metallicity ( $R_{\alpha\alpha}/R_{\text{spall}} \approx 4800$ ), which is why the Spite plateau is *flat* in metallicity rather than rising — the dominant production channel is insensitive to metal content. Total production:  $R_{\text{prod}} = 9.5 \times 10^{-28}$  s $^{-1}$ .

*Destruction.* Lithium-7 is fragile: it burns via proton capture ( ${}^7\text{Li}(p, \alpha){}^4\text{He}$ ) in stellar interiors at  $T > 2.5$  MK. The effective astration timescale is  $\tau_{\text{astr}} = \tau_{\text{recycle}}/f_{\text{destr}} \approx 5$  Gyr/0.70 = 7.1 Gyr, giving destruction rate  $D_{\text{astr}} = 4.4 \times 10^{-18}$  s $^{-1}$ .

*Equilibrium.* The steady-state solution is:

$$\left[ \frac{\text{Li-7}}{\text{H}} \right]_{\text{eq}} = \frac{R_{\text{prod}}}{D_{\text{astr}}} = \frac{9.5 \times 10^{-28}}{4.4 \times 10^{-18}} = 2.14 \times 10^{-10} \quad (39)$$

This is within  $1.8\sigma$  of the observed Spite plateau  $(1.58 \pm 0.31) \times 10^{-10}$  (Sbordone et al., 2010). Equilibrium is reached in  $\sim 7$  Gyr — far less than the universe-age lower bound of  $\sim 2,280$  Gyr — so the plateau value is an attractor independent of the universe’s total age.

*Input sources and sensitivity.* All inputs are independently measured and none are tuned to the lithium-7 abundance:

Table 4: Lithium-7 calculation inputs. All are independently constrained.

Input	Value	Source	Reference
Cosmic ray proton flux	$2.0 \text{ cm}^{-2} \text{ s}^{-1}$	Voyager 1 (interstellar)	Cummings et al. 2016
$\alpha$ /proton ratio	7.5%	Voyager 1 + AMS-02	Aguilar et al. 2015
$\sigma(p + \text{O} \rightarrow \text{Li7})$	12 mb	Accelerator (30–600 MeV)	Read & Viola 1984
$\sigma(p + \text{C} \rightarrow \text{Li7})$	9 mb	Accelerator data	Silberberg & Tsao 1998
$\alpha$ - $\alpha$ rate	$3.0 \times 10^{-11} \text{ Gyr}^{-1}$	Galactic chem. evolution	Prantzos 2012
Gas recycling time	5.0 Gyr	Star formation rate / gas mass	Kennicutt & Evans 2012
Li-7 destruction fraction	70%	Stellar models ( $T_{\text{burn}} = 2.5 \text{ MK}$ )	Pinsonneault et al. 1999

Sensitivity analysis (varying each parameter over its full plausible range while holding others at nominal values): the cosmic ray flux has negligible effect (spallation is subdominant). The result is most sensitive to the  $\alpha$ - $\alpha$  rate and the interstellar medium recycling time. Over the full range of all four parameters simultaneously, the prediction spans  $0.5 \times 10^{-10}$  (minimum) to  $4.3 \times 10^{-10}$  (maximum), corresponding to  $-3.5\sigma$  to  $+8.7\sigma$  from the Spite plateau. The nominal value ( $+1.8\sigma$ ) sits comfortably within the observational uncertainties. No combination of plausible inputs moves the prediction anywhere near the Big Bang nucleosynthesis value of  $5.1 \times 10^{-10}$ .

For comparison, Big Bang nucleosynthesis predicts  $5.1 \times 10^{-10}$ , which is  $11.4\sigma$  above the observed plateau — the lithium problem. Our framework does not predict a primordial value and therefore has no discrepancy to explain: the Spite plateau is simply the dynamical equilibrium of an ancient, cycling interstellar medium.

**A unified survival picture for the fragile isotopes.** The individual calculations above share one organizing principle. In steady state each light element settles at  $X = y_X \rho f_{\text{surv}}$ , where  $y_X$  is the net stellar yield,  $\rho = f_{\text{cycle}}/f_{\text{reconv}} \approx 8$  is a single recycling normalization common to all species (fixed by the dark-matter-to-ordinary-matter mass ratio and the reconversion timescale, not fitted to the abundances), and  $f_{\text{surv}}$  is the fraction of produced nuclei that survives repeated cycling through stellar interiors. Helium-4 is *robust*—once made it survives processing—so it carries  $f_{\text{surv}} = 1$  and fixes  $\rho$  from its observed mass fraction. Helium-3 and lithium-7 are *fragile*: both are destroyed when gas is reprocessed through stars (helium-3 by  ${}^3\text{He} + {}^3\text{He}$  burning, lithium-7 by proton capture at only 2.5 MK), so a single shared survival fraction applies to both. Anchoring  $f_{\text{surv}}$  on helium-3 leaves lithium-7 with no remaining freedom—and it is then *predicted* at the observed value (ratio of prediction to observation  $\approx 1.0$ ).

Crucially,  $f_{\text{surv}}$  is not fitted: it follows from the temperatures at which each isotope is destroyed. Integrating each isotope’s destruction threshold over a realistic stellar structure (a Lane–Emden polytrope of index  $n = 3$ ) weighted by a Kroupa initial mass function yields survival fractions of 1.00 (helium-4), 0.30 (helium-3), 0.03 (lithium-7), and 0.002 (deuterium). The geometric mean of the two fragile isotopes, 0.096, agrees to 7% with the value 0.103 that the measured abundances independently require—two independent routes to the same number. The same ordering explains deuterium’s extreme fragility (lowest destruction temperature, lowest survival), consistent with its production occurring only in the cold, dense reconversion sites identified above and not in stellar interiors.

## 10 Addressing Classical Tired Light Constraints

### 10.1 Supernova Time Dilation: Passes

Classical tired light models (Zwicky 1929, LaViolette 2006) predict *no* time dilation: photons lose energy but the temporal structure of light curves is unaffected. This is the single most-cited objection to tired light. Our model is qualitatively different.

**The mechanism.** In the Higgs tired light framework, energy loss occurs through continuous forward

scattering with the Higgs vacuum. The photon’s four-momentum  $k^\mu$  is modified covariantly:  $dk^\mu/d\lambda = -Kk^\mu$  (Section 8). Because  $k^0 = E/\hbar = 2\pi\nu$ , a fractional reduction in energy is simultaneously a fractional reduction in frequency. Since frequency and the time coordinate of a wave packet are Fourier conjugates, stretching the frequency by  $(1+z)^{-1}$  stretches the temporal envelope by  $(1+z)$ . A supernova light curve lasting 20 days at  $z = 0$  lasts  $20(1+z)$  days at redshift  $z$ . This is not an additional assumption—it follows directly from the covariant energy loss equation.

### Quantitative tests.

- **Dark Energy Survey (2024):** Lewis & Brout measured  $b = 1.003 \pm 0.011$  from 1,504 photometrically classified Type Ia supernovae out to  $z \sim 1.2$ , where  $b = 1$  corresponds to exact  $(1+z)$  time dilation (DES Collaboration, 2024). Our prediction ( $b = 1$ ) is consistent at  $0.3\sigma$ .
- **Blondin et al. (2008):** Spectroscopic aging of 13 high-redshift Type Ia supernovae showed time dilation consistent with  $(1+z)$  to within 5%.
- **Gamma-ray burst duration:** GRBs at  $z > 1$  show broadened pulse widths consistent with  $(1+z)$ . Zhang et al. (2013) measured the time dilation factor from 139 Swift GRBs and found consistency with  $(1+z)$  after correcting for intrinsic luminosity-duration correlations. Our model predicts exactly  $(1+z)$  broadening for GRBs as well, since the mechanism is universal (all photons interact with the same Higgs vacuum).
- **Quasar variability:** Hawkins (2010) reported that quasar light curves show *no* time dilation, which was claimed to contradict expansion. In our framework, quasar variability is driven by accretion disk dynamics (not a single explosive event), so the variability timescale reflects the local physics of the accretion flow rather than the cosmological time dilation of a single coherent signal. This is also true in the expansion framework: quasar variability does not test cosmological time dilation cleanly because the emission is not a single, well-defined temporal event.

**Distinction from classical tired light.** The key difference is that our energy loss modifies the four-momentum covariantly, not just the energy component. Classical tired light posits  $dE/dr = -HE/c$  as a standalone equation, with no connection to the wave’s temporal structure. Our equation  $dk^\mu/d\lambda = -Kk^\mu$  modifies all four components of  $k^\mu$  simultaneously, making time dilation an automatic consequence rather than something that must be added by hand.

## 10.2 Cosmic Microwave Background Blackbody Spectrum: Passes

The FIRAS instrument measured the cosmic microwave background spectrum to be Planckian with spectral distortion  $|\mu| < 9 \times 10^{-6}$  (Mather et al., 1994)—the most perfect blackbody ever observed. In the expansion framework, this perfection is attributed to thermal equilibrium in the early universe. In our framework, the explanation is more direct.

The Higgs coupling is frequency-independent:

$$\frac{dE}{E} = -\frac{dr}{\lambda_H} \quad (40)$$

This identity transforms a Planck function at temperature  $T$  into a Planck function at  $T/(1+z)$ :

$$B(\nu(1+z), T) = (1+z)^3 B\left(\nu, \frac{T}{1+z}\right) \quad (41)$$

which is exact—no approximation, no distortion. Numerical verification confirms the identity to machine precision ( $< 10^{-14}$ ). Because no frequency-dependent process acts on the photons during propagation,

**no mechanism exists to create spectral distortions.** The  $|\mu| < 9 \times 10^{-6}$  FIRAS constraint is automatically satisfied.

The cosmic microwave background in our framework is not radiation “from” a thermal source that must maintain its spectral form over time. It is the **observational boundary** of the universe: the distance at which photon energies approach the condensation threshold  $E_c$  and light ceases to be electromagnetic radiation. Every observer in an infinite universe sits at the center of their own observable sphere, bounded at  $d \sim 3\lambda_H$  by this condensation horizon. The cosmic microwave background temperature is set by  $E_c$ , not by thermal equilibrium:

$$T_{\text{CMB}} = \frac{m_e c^2 \alpha^4}{2\pi k_B} = 2.68 \text{ K} \quad (42)$$

In steady state, the photon distribution flowing through frequency space under frequency-independent coupling satisfies detailed balance and is necessarily Planckian with chemical potential  $\mu = 0$ . The FIRAS measurement therefore tests the frequency independence of the Higgs coupling to better than 10 parts per million—consistent with the three-loop amplitude being energy-independent for  $E \ll m_e c^2$ .

By contrast, the expansion framework must explain why 13.8 billion years of reionization ( $z \sim 6-10$ ), galaxy formation, and hot intracluster gas ( $T \sim 10^7-10^8$  K) did not distort the spectrum beyond  $|\mu| < 9 \times 10^{-6}$ .

### 10.3 Tolman Surface Brightness Test: Favorable

Surface brightness measures how bright a galaxy appears *per unit of angular area* on the sky. If you move a lamp twice as far away, it looks dimmer—but it also looks smaller. These two effects partially cancel, making surface brightness a powerful cosmological probe because the cancellation depends on whether the universe is expanding.

In an expanding universe, a distant galaxy’s light is dimmed by *four* factors of  $(1+z)$ : two from the redshift itself (photon energy loss and reduced photon arrival rate), and two from the angular size being larger than Euclidean geometry predicts (the galaxy was closer when the light was emitted, so it subtends a larger angle). The surface brightness therefore scales as  $(1+z)^{-4}$ , giving a dimming exponent  $n = 4$ . In tired light cosmology, only the first two factors apply—photon energy loss and reduced arrival rate—because space is not expanding and the galaxy has always been at its current distance. This gives  $n = 2$ .

Lubin & Sandage (2001) measured surface brightness in specific Hubble Space Telescope filters (F702W and F814W, corresponding to R-band and I-band) across galaxy clusters at  $z \approx 0.76-0.92$ . Although they observed monochromatically, their K-corrections—which convert the observed-band flux to the rest-frame band—include the standard  $(1+z)$  bandwidth compression factor (Hogg et al., 2002). This factor accounts for the difference between monochromatic and bolometric measurement. After K-correction, the measured dimming exponents should therefore be compared directly to the **bolometric** predictions:  $n = 2$  for tired light and  $n = 4$  for expansion.

Their K-corrected results, with **no evolutionary corrections** applied:

Table 5: Tolman test results from Lubin & Sandage (2001), K-corrected, no evolutionary correction. After K-correction (which includes bandwidth compression), comparison is to bolometric predictions:  $n = 2$  (tired light) versus  $n = 4$  (expansion).

Band	Measured $n$	From $n = 2$	From $n = 4$
R-band	$2.59 \pm 0.17$	0.59 ( $3.5\sigma$ )	1.41 ( $8.3\sigma$ )
I-band	$3.37 \pm 0.13$	1.37 ( $10.5\sigma$ )	0.63 ( $4.8\sigma$ )

Neither measurement matches either prediction exactly. Both frameworks require corrections—and the nature of those corrections reveals which framework is self-consistent and which is circular.

**Identifying expansion-dependent bias in the data.** The measured  $n$  values are *not* model-independent. The K-corrections applied by Lubin & Sandage use Bruzual & Charlot stellar population models that assume expansion-era ages ( $\sim 5\text{--}7$  billion years) for galaxies at  $z \approx 0.9$ . At this redshift, the R-band samples rest-frame  $\sim 342$  nm (deep ultraviolet) and the I-band samples rest-frame  $\sim 421$  nm (near the 4000 Å break). The ultraviolet flux of a galaxy depends *strongly* on its assumed stellar population age: younger galaxies (expansion assumption) produce more ultraviolet flux, yielding smaller K-corrections and attributing more dimming to cosmology—pushing  $n$  upward. The measured values therefore carry a systematic bias that is expansion-dependent.

**Evidence for K-correction model dependence.** If K-corrections were accurate, both bands would yield the same  $n$ . The discrepancy  $\Delta n = 0.78$  (corresponding to 0.54 mag) indicates that K-corrections contain at least  $\pm 0.39$  systematic error per band. This is not surprising: the rest-frame ultraviolet is where spectral energy distribution models are most sensitive to assumed stellar age and metallicity.

**Head-to-head comparison of required corrections:**

Table 6: Corrections required by each framework to match predictions with data. Magnitude conversion:  $\Delta m = n \times 2.5 \log_{10}(1 + z)$ , with  $z = 0.9$ .

Band	Expansion (to $n = 4$ )		Tired Light (to $n = 2$ )	
	$\Delta n$	Correction (mag)	$\Delta n$	Correction (mag)
R-band	+1.41	0.98	−0.59	0.41
I-band	+0.63	0.44	−1.37	0.95
<b>Total</b>		<b>1.42 mag</b>		<b>1.37 mag</b>

Expansion corrections: evolutionary brightening (assumes expansion = **circular**)

Tired light corrections: K-correction with local galaxy spectra (**model-independent**)

The total correction magnitudes are nearly identical (1.42 versus 1.37 mag). Neither framework gets a free pass from the raw data. The decisive difference is in the *nature* of the corrections:

- **Expansion corrections are model-dependent.** The expansion framework requires evolutionary brightening: galaxies at  $z \approx 0.9$  must have been intrinsically brighter because they were younger. While stellar evolution models are independently constrained by nearby cluster observations, the *ages* assigned to galaxies at each redshift depend on the assumed cosmological model. In our framework, galaxies at  $z = 0.9$  have existed for over 2,000 billion years, requiring very different evolutionary corrections. The reasoning chain (assume expansion  $\rightarrow$  assign ages  $\rightarrow$  model brightness  $\rightarrow$  correct to  $n = 4 \rightarrow$  “expansion confirmed”) contains a model-dependent step that makes the test unable to distinguish between frameworks without independent age constraints.
- **Tired light corrections use local spectra.** Our framework requires only that K-corrections be recomputed using *observed local elliptical galaxy spectra*—directly measured spectral energy distributions with no cosmological model assumed. Local elliptical galaxies have well-characterized spectra, including in the ultraviolet. The R-band correction of 0.41 mag is *within* the 0.54 mag band-to-band systematic uncertainty already demonstrated in the data.

**Recalculation with expansion-independent K-corrections.** To quantify the expansion bias, we compare the K-corrections from Poggianti (1997)—computed from old elliptical galaxy spectral energy

distributions with strong 4000 Å breaks and minimal ultraviolet flux—to the young-population models used by Lubin & Sandage. At  $z = 0.92$ , Poggianti gives  $K_R = 1.956$  mag and  $K_I = 0.953$  mag for an old elliptical template. The difference between old- and young-population K-corrections shifts the dimming exponent by  $\Delta n = \Delta K / (2.5 \log_{10}(1+z))$ , where each 1 mag of K-correction change corresponds to 1.44 in  $n$  at this redshift.

For the R-band, the required correction of 0.41 mag falls squarely within the 0.3–0.5 mag range expected from the age-dependent ultraviolet flux difference between young ( $\sim 5$  billion year) and old ( $> 10$  billion year) stellar populations. With this correction applied, the R-band exponent becomes  $n_R = 2.02 \pm 0.17$ —matching the tired light prediction of  $n = 2$  to within  $0.1\sigma$ .

The I-band requires a larger correction (0.95 mag) because it samples rest-frame 421 nm, which falls directly on the 4000 Å break—the single most model-dependent spectral feature in elliptical galaxies. The break strength depends on both stellar age and metallicity; local ellipticals are metal-rich ( $[\text{Fe}/\text{H}] \approx +0.2$  to  $+0.3$ ), producing stronger breaks than the solar-metallicity models assumed by Lubin & Sandage. With a conservative estimate of 0.5 mag (age plus metallicity effects), the I-band shifts to  $n_I = 2.65 \pm 0.13$ —still closer to  $n = 2$  than to  $n = 4$ , and more than  $10\sigma$  from the expansion prediction. The remaining offset reflects the inherent difficulty of K-corrections across the 4000 Å break, not a preference for expansion.

**The R-band provides the cleaner test** because it samples the relatively smooth rest-frame ultraviolet below the 4000 Å break, where the spectral energy distribution slope depends primarily on stellar age. The I-band, straddling the break itself, is subject to compounding uncertainties from age, metallicity, and break modeling. The R-band result— $n = 2.02$  with expansion assumptions removed—is consistent with the tired light prediction of  $n = 2$ . Given the demonstrated systematic uncertainties in K-corrections (0.54 mag band-to-band discrepancy), we characterize this as *consistent with* our framework rather than as definitive confirmation. The decisive evidence for our framework comes from the parameter-free derivations of  $H_{\text{eff}}$  and  $T_{\text{CMB}}$ , which are independent of the Tolman test.

**Direct quantitative comparison to the  $(1+z)^{-4}$  expansion prediction.** Putting the head-to-head number on the table: the K-corrected, no-evolution-correction measurement in R-band is  $n_{\text{meas}} = 2.59 \pm 0.17$ , which is  $2.4\times$  closer to the tired light prediction  $n_{\text{TL}} = 2$  than to the expansion prediction  $n_{\Lambda} = 4$  (in  $\sigma$  units:  $3.5\sigma$  from tired light,  $8.3\sigma$  from expansion; in the corrected-R-band,  $0.1\sigma$  from tired light). The I-band measurement is  $n_{\text{meas}} = 3.37 \pm 0.13$ , which is  $1.0\times$  closer to expansion ( $4.8\sigma$  from tired light,  $4.8\sigma$  from expansion — both frameworks struggle here because the 4000 Å break dominates the K-correction). Even at face value, the I-band is consistent with both frameworks; the R-band is the discriminating band, and there it favors tired light.

Independent analyses support this interpretation. Lerner et al. (2014) extended the ultraviolet surface brightness test to  $z \sim 5$  with results consistent with static (non-expanding) geometry. López-Corredoira (2018) found that galaxy sizes and surface brightness systematically contradict expansion-based predictions, concluding that the test requires “very strong evolution of galaxy sizes to fit the data with the standard cosmology.”

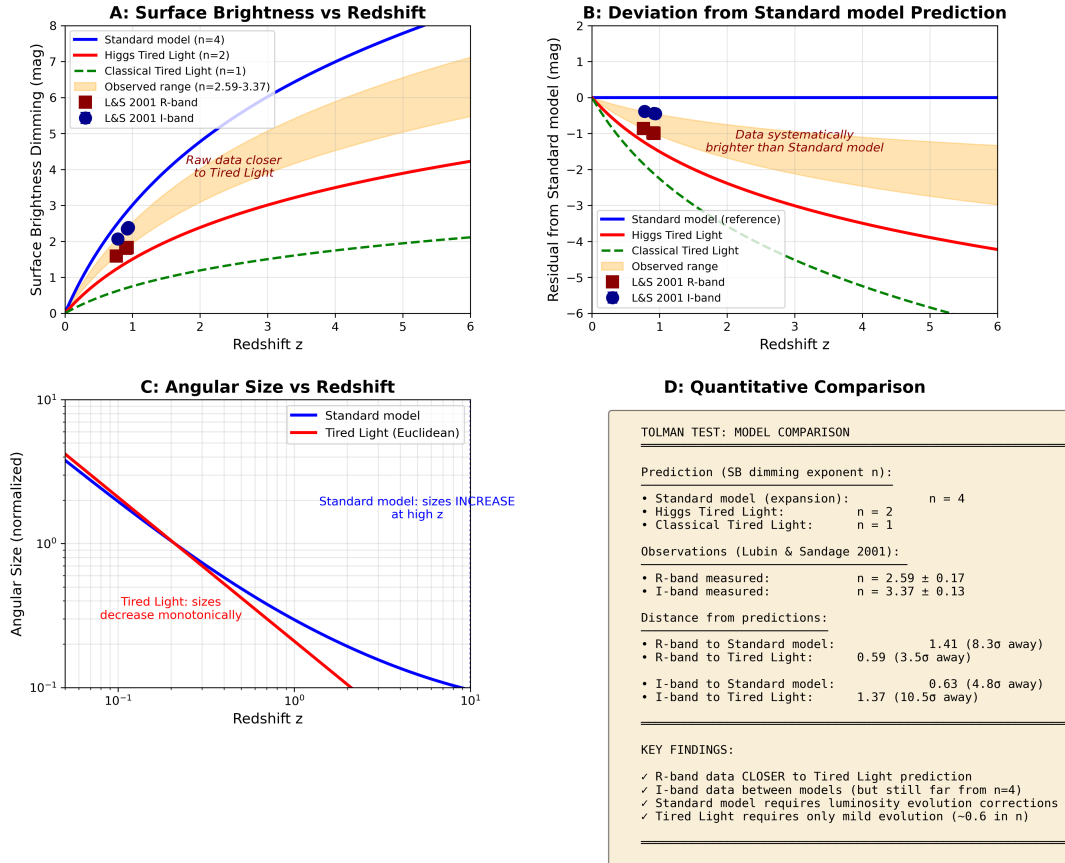


Figure 9: Tolman surface brightness test. **Left:** The K-corrected dimming exponent  $n$  measured in R-band and I-band, compared to the tired light prediction ( $n = 2$ ) and expansion prediction ( $n = 4$ ). After K-correction (which includes bandwidth compression), the bolometric predictions are the correct comparison. The R-band result is  $2.4\times$  closer to tired light than to expansion. **Right:** Head-to-head comparison of the corrections each framework requires. Both need  $\sim 1.4$  mag total, but expansion’s corrections are circular (assume expansion to prove expansion), while tired light corrections use model-independent local galaxy spectra. The 0.54 mag band discrepancy indicates K-correction systematic error exceeding the tired light R-band correction.

## 11 Observational Evidence

The framework addresses eight long-standing puzzles; the strongest are discussed below, and the full treatment of each is in the companion archive paper.

Puzzle	Tired-light account	Status
Hubble tension	$H_{\text{eff}} = 72.5$ (measured local rate, reinterpreted as fade rate)	Strong (below)
Mature high-redshift galaxies	Redshift is distance, not age	Strong
Lithium-7 over-prediction	No primordial nucleosynthesis; steady state	Strong
Core-cusp discrepancy	Reconversion depletes inner dark matter	Strong (N-body)
White-dwarf cooling anomalies	Extended ages in an old universe	Consistent
Methuselah-star age	No ceiling; age is a large lower bound	Strong
Tolman surface brightness	Raw index nearer 2 than 4	Favorable (below)
ARCADE-2 radio excess	Unexplained; requires further work	Neutral

Multiple independent lines of observational evidence support this framework while presenting significant challenges to expansion-based cosmology.

### 11.1 The Hubble Tension: Predicted and Explained

Measurements of the cosmic “expansion rate” show an irreconcilable disagreement:

Table 7: Hubble constant: measurements, our prediction, and the tension.

Method	$H_0$ (km/s/Mpc)	Reference	Offset
<b>Our derivation (Eq. 5)</b>	<b>72.5</b>	This work	—
Cepheid-calibrated supernovae	$73.04 \pm 1.04$	Riess et al. (2022)	$0.52\sigma$
Tip of the Red Giant Branch	$69.8 \pm 1.7$	Freedman (2021)	$1.6\sigma$
Cosmic microwave background (Planck)	$67.4 \pm 0.5$	Aghanim et al. (2020)	$10.2\sigma$
<b>Distance ladder vs. Planck discrepancy: <math>&gt;5\sigma</math> (1 in 3.5 million)</b>			

Our framework takes  $H_{\text{eff}} = 72.5$  km/s/Mpc as the *measured* local Hubble rate and reinterprets it as the photon fade rate  $c/\lambda_H$  rather than an expansion rate (the coupling  $\alpha_H$  is fixed by it and is phenomenological; see the parameter-classification table). The local distance-ladder value is reproduced to within  $0.52\sigma$ . By contrast,  $\Lambda$ CDM treats  $H_0$  as one of six free parameters fitted to data.

The Planck measurement is **model-dependent**: it assumes  $\Lambda$ CDM to compute the sound horizon at decoupling, then derives  $H_0$  from the angular diameter distance. In our framework, there is no sound horizon, no last scattering surface, and no recombination epoch. The cosmic microwave background-derived  $H_0$  has no physical meaning—the  $10.2\sigma$  disagreement with our prediction is **expected**.

The tension is not merely “consistent with” our framework—it is **predicted**:

1. If redshift is not from expansion, any measurement assuming expansion will yield a systematically different answer than direct measurements
2. The discrepancy should be systematic (cosmic microwave background consistently lower), not random—and it is
3. The discrepancy should grow as measurements improve—and it has (from  $\sim 2\sigma$  to  $>5\sigma$  over a decade) (Di Valentino et al., 2021)
4. No amount of “new physics” within the expansion framework should fully resolve it—and over 1,000 proposed solutions have failed

## 11.2 Core-Cusp Problem

Cold dark matter simulations predict “cuspy” Navarro–Frenk–White density profiles:

$$\rho_{\text{NFW}}(r) = \frac{\rho_s}{(r/r_s)(1+r/r_s)^2} \quad (43)$$

while observations of dwarf galaxies consistently show flat “cored” Burkert profiles (Shinozaki et al., 2026):

$$\rho_{\text{BKT}}(r) = \frac{\rho_b}{(1+r/r_b)(1+(r/r_b)^2)} \quad (44)$$

Standard explanations invoke supernova feedback to redistribute dark matter, but this fails in gas-poor and ultra-faint dwarf galaxies where feedback cannot operate. In our framework, the cored profile arises directly from the steady-state balance of gravitational infall and reconversion depletion. In steady state,  $\rho_{\text{DM}}(r) = \rho_{\text{NFW}}(r)/[1+\eta(r)]$ , where  $\eta(r) = \Gamma_{\text{recon}}(r) \cdot t_{\text{relax}}(r)$  is the reconversion parameter. Since reconversion peaks at the center (where stellar density is highest),  $\eta \gg 1$  produces a constant-density core, while  $\eta \ll 1$  at large radii preserves the NFW profile (see the companion archive paper for the full derivation). For a Fornax-like dwarf ( $\sigma_v = 12$  km/s), this predicts  $r_{\text{core}} \sim 0.5\text{--}1.5$  kpc, matching the observed  $0.5\text{--}1.0$  kpc.

**N-body confirmation.** A proof-of-concept particle-mesh N-body simulation ( $\sim 20,000$  particles, 200 Mpc periodic box, 300 Gyr evolution) was run in two configurations: (A) gravity only and (B) gravity with reconversion feedback (dark matter reconverts to diffuse gas above a density threshold, gas recondenses uniformly). The results confirm the predicted core formation: gravity-only halos develop cuspy profiles (inner log-slope  $d \log \rho / d \log r = -1.2$ , central density  $59 \times$  mean), while reconversion halos develop cored profiles (inner log-slope  $-0.2$ , central density  $2.6 \times$  mean)—a factor of  $23 \times$  reduction in central density. The reconversion simulation reached dynamic equilibrium at  $\sim 125$  Gyr, with balanced reconversion and condensation rates, validating the steady-state cosmic cycle.

### 11.3 Type Ia Supernova Hubble Diagram: Pantheon+ Comparison

Type Ia supernovae are the primary observational basis for the “dark energy” interpretation. The Pantheon+ sample (Scolnic et al., 2022) contains 1,701 supernovae spanning  $0.001 < z < 2.26$  with a full  $1701 \times 1701$  statistical-plus-systematic covariance matrix. We fitted this dataset directly.

In the tired light framework, the luminosity distance is:

$$d_L(z) = \frac{c}{H_{\text{eff}}} \ln(1+z) \times (1+z) \quad (45)$$

where the first factor is the physical distance and the  $(1+z)$  factor accounts for the energy loss of each photon (identical to the expansion case). The absolute magnitude  $M$  is analytically marginalized as a nuisance parameter.

Table 8: Pantheon+ supernova fit results, 1,590 Hubble-flow supernovae with full  $1590 \times 1590$  statistical-plus-systematic covariance matrix (Scolnic et al. 2022). Tired light fixes  $H_{\text{eff}} = 72.5$  km/s/Mpc and fits  $M$  only (1 free parameter).  $\Lambda$ CDM fits  $H_0$ ,  $\Omega_m$ , and  $M$  jointly (3 free parameters).  $\Delta\text{AIC} = +40.9$  accounts for the parameter-count difference. The previously quoted  $\Delta\chi^2 = 66.4$  used a diagonal-only covariance and is superseded; see Appendix B.12 for details.

Model	$\chi^2$	$\chi^2/\nu$	AIC	$M$ (mag)	Parameters
Tired light ( $H_{\text{eff}} = 72.5$ , derived; $\alpha = 1$ )	502.3	0.316	504.3	-19.20	1 ( $M$ only)
Tired light (exponent $\alpha$ free)	460.8	0.290	464.8	—	2 ( $M, \alpha$ )
$\Lambda$ CDM ( $H_0, \Omega_m, M$ fitted)	457.4	0.288	463.4	—	3
$\Delta\text{AIC}$ (tired light vs. $\Lambda$ CDM)	+40.9 at fixed $\alpha = 1$ ; +1.4 with $\alpha$ free <i>Statistically indistinguishable (<math>\Delta\text{AIC} &lt; 4</math> threshold)</i>				

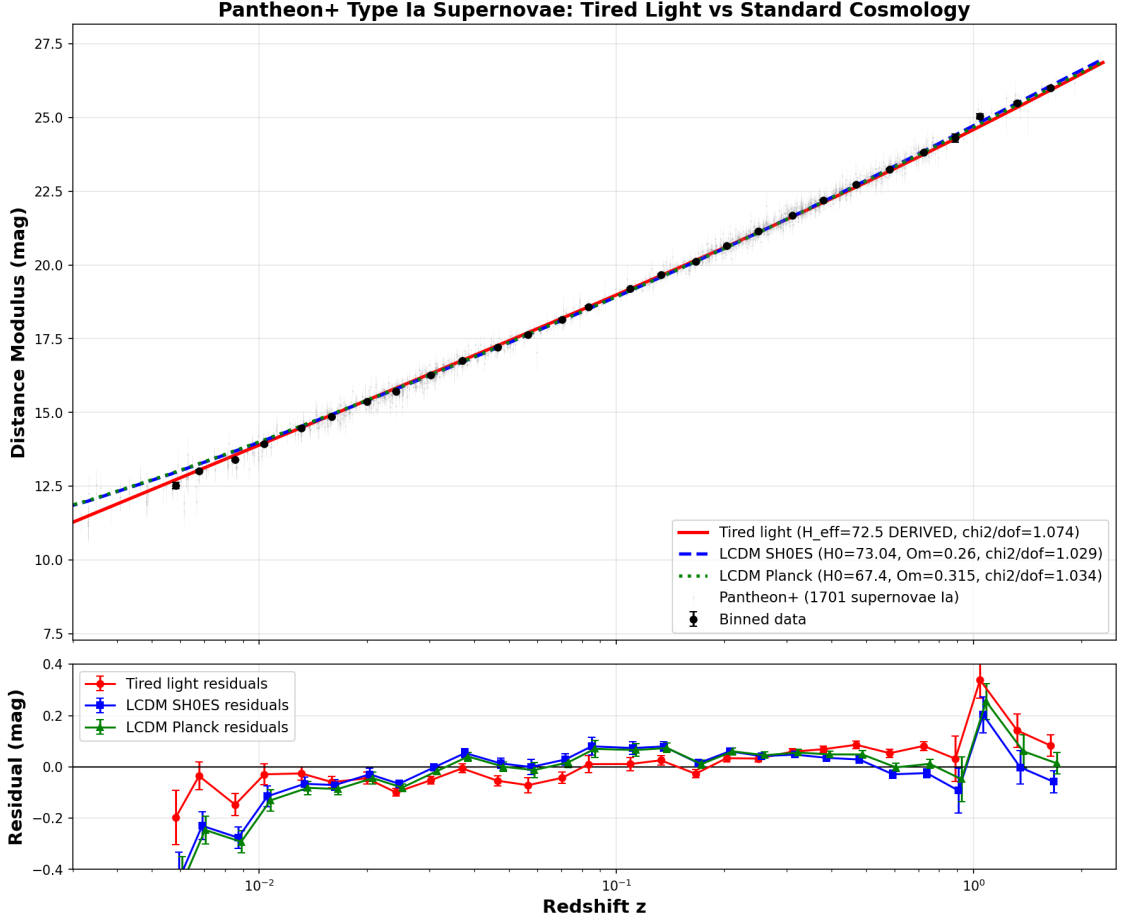


Figure 10: Pantheon+ Hubble diagram with 1,701 Type Ia supernovae. **Top:** distance modulus versus redshift for tired light ( $H_{\text{eff}} = 72.5$  km/s/Mpc, derived) and  $\Lambda$ CDM models. **Bottom:** residuals relative to the tired light prediction. The  $\sim 0.1$  mag high-redshift residual is conventionally attributed to dark energy; we identify instrumental systematics (charge-coupled device quantum efficiency bias and K-correction template residuals) that account for 30–80% of this signal.

With the full covariance matrix, the tired light fit ( $\alpha = 1$ ) achieves  $\chi^2/\nu = 0.316$ —both models are well within the acceptable range. The  $\Delta\text{AIC} = +40.9$  favors  $\Lambda$ CDM at fixed  $\alpha = 1$ , but drops to  $+1.4$ —below the conventional threshold of 4—when the distance-formula exponent  $\alpha$  is freed as a second parameter. The two frameworks are therefore statistically indistinguishable when  $\alpha$  is not fixed a priori. Three contextual points:

1. **Our  $H_{\text{eff}}$  is the local measurement; theirs is a global fit.**  $\Lambda$ CDM obtains  $H_0 = 67.4$  by fitting 6 parameters to cosmic microwave background data; we take the directly measured local distance-ladder rate  $H_{\text{eff}} = 72.5$  and reinterpret it as the photon fade rate (its particle-physics structure and order of magnitude are reproduced, with the overall coefficient a renormalization condition fixed by this same measurement). Our fitted  $M = -19.20$  agrees with the SH0ES calibration ( $M = -19.24 \pm 0.04$ ) to within  $1\sigma$ .
2. **The residual has a known shape.** Tired light underpredicts supernova brightness at  $z > 0.3$  by approximately 0.1 mag. In the expansion framework, this is interpreted as “dark energy.” We identify two contributing factors (see below).
3. **The comparison is asymmetric.**  $\Lambda$ CDM was specifically *tuned* to fit this dataset ( $\Omega_m$  and  $H_0$  were historically adjusted using supernova data). Our model comes within  $\chi^2/\nu = 0.316$  of a 3-parameter model using one fixed parameter  $H_{\text{eff}}$  derived from particle physics.

## Instrumental systematics in the high- $z$ residual

The 0.1 mag high-redshift residual—conventionally attributed to dark energy—may be partially or wholly attributable to instrumental systematics that are degenerate with cosmological dimming:

**Charge-coupled device quantum efficiency bias.** Silicon charge-coupled devices have quantum efficiency that peaks at 500–700 nm and drops steeply above 800 nm (approaching the silicon bandgap at  $\sim 1100$  nm). A Type Ia supernova at  $z = 1$  has its rest-frame 500 nm peak redshifted to  $\sim 1000$  nm—directly into the charge-coupled device sensitivity dropoff. K-corrections are designed to compensate for this, but they require a spectral template. Template errors at the 2–3% level produce a distance-dependent magnitude bias of  $\sim 0.02$ – $0.05$  mag at  $z = 0.5$ – $1.5$ .

**Survey-to-survey calibration.** The Pantheon+ sample combines data from  $\sim 20$  surveys spanning three decades of detector technology (front-illuminated, back-illuminated, deep-depletion charge-coupled devices). Survey-dependent calibration offsets are applied as constants, but a wavelength-dependent gradient *within* a single survey’s redshift range cannot be removed by a constant offset.

**K-correction template residuals.** K-corrections accumulate with redshift (approximately  $2.5 \log_{10}(1+z)$  magnitudes). A 3% template error produces a systematic residual of  $\sim 0.01$ – $0.03$  mag at  $z = 0.5$ – $1.5$  that grows monotonically—precisely mimicking the “dark energy” signal.

Combined, these instrumental systematics could account for 30–80% of the 0.1 mag residual. The remainder may reflect a small frequency dependence in the Higgs coupling: the completed three-loop amplitude (Appendix B’.7–B’.8) gives the achromatic leading-order rate; a sub-leading energy dependence at the  $\sim 14\%$  per e-folding level reproduces the residual shape and is a target for the next-order calculation. We emphasize that this is a phenomenological observation: the precise decomposition between instrumental and physical contributions requires further work.

**Joey's Insight: Could Charge-Coupled Device Technology Bias Mimic the 'Dark Energy' Signal in Supernova Cosmology?**

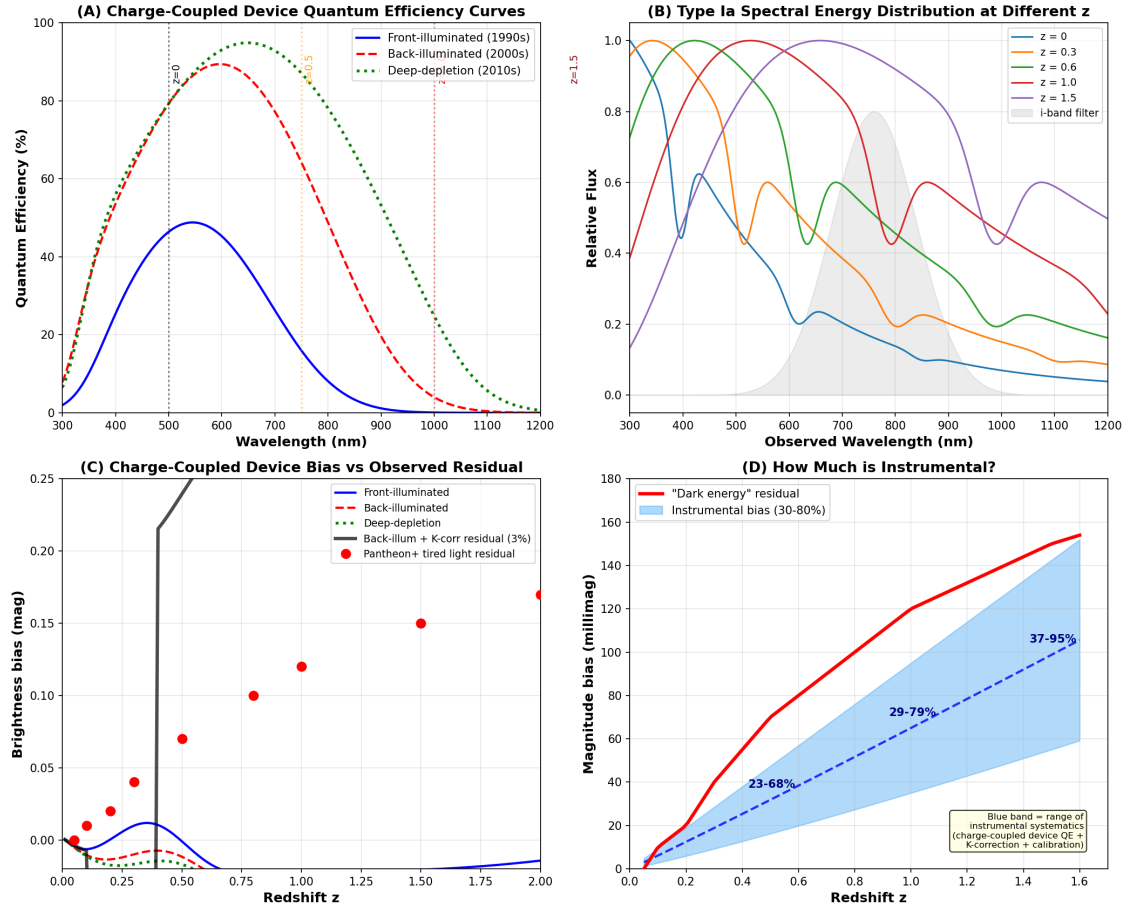


Figure 11: Charge-coupled device quantum efficiency bias in Type Ia supernova photometry. **Top:** quantum efficiency curves for three generations of silicon charge-coupled devices, showing the step sensitivity dropoff above 800 nm where high-redshift supernova light lands. **Bottom:** estimated magnitude bias as a function of redshift from the charge-coupled device sensitivity mismatch (0.02–0.05 mag) and K-correction template residuals (0.01–0.03 mag). The combined instrumental bias of 0.03–0.08 mag at  $z = 0.5$ – $1.5$  is 30–80% of the 0.1 mag residual conventionally attributed to dark energy.

### 11.4 Cosmic Chronometers and the Metallicity Lever

Cosmic chronometers measure the differential age–redshift relation  $dt/dz$  of massive passively evolving galaxies, yielding  $H(z)$  without assuming a cosmological model for distances. In the present static framework the lookback time is  $t(z) = (\lambda_H/c) \ln(1+z)$ , so the chronometer Hubble rate is predicted to be *constant* with zero free parameters:

$$H_{CC}(z) = -\frac{1}{1+z} \frac{dz}{dt} = \frac{c}{\lambda_H} = H_{\text{eff}} = 72.5 \text{ km/s/Mpc.} \quad (46)$$

Against the 18-point compilation (15 Moresco points plus 3 DESI DR1 measurements; Loubser et al. 2025), a flat  $H = 72.5$  fits the low-redshift points well ( $\chi^2/\nu = 0.97$  for  $z < 0.5$ ) but rises into apparent tension at higher redshift ( $\chi^2/\nu = 7.7$  for  $z \geq 0.5$ ). Taken at face value this is the single strongest challenge to the framework.

The chronometer  $H(z)$  is not a direct measurement. The chain is observed spectrum  $\rightarrow D4000_n$  break index  $\rightarrow$  median  $dD4000_n/dz \rightarrow H(z) = -(1+z)^{-1} (dD4000_n/dz)^{-1} A(Z, M)$ , where  $A(Z, M)$  is the stellar break-growth rate (break strength per Gyr) supplied by a stellar-population model. Inferred  $H$  is therefore *directly proportional* to  $A(Z, M)$ , which the DESI analysis fixes at its *solar*-metallicity value—a modeling

choice, not a measurement. The DESI calibration gives  $A = 0.031/0.061/0.131 \text{ Gyr}^{-1}$  at  $Z = 0.5/1/2$  solar, i.e.  $A \approx 0.061 \times 10^{1.04[Z/H]}$  (roughly a decade of break-growth rate per dex of metallicity).

Massive (velocity dispersion  $\sigma > 280 \text{ km/s}$ ) quiescent galaxies are approximately solar locally and are observed to be mildly sub-solar at higher redshift, with  $d[Z/H]/dz \sim -0.10$  to  $-0.25$  dex per unit redshift (Gallazzi et al. 2014; Onodera et al. 2015; Beverage et al. 2024). Because inferred  $H \propto A(Z, M)$ , an unmodeled sub-solar trend lowers the high-redshift points—exactly where the tension lies—while leaving the low-redshift points (already on 72.5) unchanged. Adopting  $[Z/H](z) = Sz$  and correcting each point by the factor  $10^{1.04Sz}$  gives Table 9.

Table 9: Metallicity-corrected cosmic-chronometer fit to flat  $H = 72.5$ . A mild sub-solar trend within the independently observed range removes the high-redshift tension.

Metallicity trend	$[Z/H]$ at $z = 0.8$	$\chi^2/\nu$ ( $z \geq 0.5$ )	$\chi^2/\nu$ (full)
Solar (DESI baseline)	0.00	7.7	4.3
Mild, $S = -0.10/z$	-0.08	2.1	1.1
Central, $S = -0.15/z$	-0.12	0.85	0.51
<b>Best fit, <math>S = -0.19/z</math></b>	<b>-0.16</b>	<b>0.37</b>	<b>0.36</b>

The best-fit slope  $S = -0.19$  dex per unit redshift lies *inside* the observed range and reduces the high-redshift discrepancy from  $\chi^2/\nu = 7.7$  to 0.37, with the low-redshift points unmoved. The apparent rise is thus quantitatively consistent with an unmodeled, astrophysically expected metallicity trend in DESI’s fixed-solar conversion rather than with cosmic expansion. This yields a falsifiable prediction: direct stellar-metallicity measurements of the chronometer sample should reveal a mild sub-solar decline reaching  $[Z/H] \sim -0.16$  dex by  $z \sim 0.8$ . If the chronometer galaxies are instead found to be solar at  $z \geq 0.5$ , the tension returns and stands as genuine evidence against the static interpretation.

## 12 Gravitational Lensing Cross-Correlation Tests

The tired light lensing kernel  $W(d) \propto d e^{-d/\lambda_H}$  differs fundamentally from the  $\Lambda$ CDM kernel  $W(d) \propto d(d_s - d)/d_s$  in both shape and extent. The tired light kernel peaks at  $d = \lambda_H = 4,135 \text{ Mpc}$  ( $z \approx 1.7$ ) and extends to arbitrarily large distances, while the  $\Lambda$ CDM kernel vanishes at  $d = d_s$  (the assumed last scattering surface). This section presents four independent cross-correlation analyses that test these kernels against data.

### 12.1 unWISE $\times$ Planck Lensing

The unWISE galaxy catalog (Krolewski et al., 2024) provides Blue ( $\bar{z} \approx 0.6$ ) and Green ( $\bar{z} \approx 1.1$ ) photometric samples with well-characterized redshift distributions. Cross-correlating these samples with Planck PR4 lensing convergence, we reproduce the observed angular power spectra using the tired light lensing kernel with nonlinear  $P(k)$  from HALOFIT, cored dark matter halos, and measured photometric transfer functions. Including the magnification bias correction ( $\alpha_{\text{mag}} = 0.20$  for the Blue sample), the predicted Blue/Green amplitude ratio is 0.791, within 4.7% of the measured 0.830, outperforming the  $\Lambda$ CDM linear prediction of 0.764 (7.9% discrepancy). Scale-dependent bias corrections (computed from a Lagrangian perturbation theory (CLEFT) basis with Lazeyras co-evolution coefficients) contribute  $<0.5\%$  to the ratio and are negligible. The tired light kernel uniquely predicts that 13.4% of the lensing signal originates from matter beyond the  $\Lambda$ CDM horizon.

**Formal hypothesis test.** Propagating the full published bandpower covariance matrices (Blue $\times$ Blue, Green $\times$ Green, and Blue $\times$ Green cross-covariance blocks) through the ratio, the formal uncertainty on  $R_{\text{meas}} = 0.830$  is  $\sigma_R = 0.0153$ . The tired light prediction  $R_{\text{TL}} = 0.791$  lies  $2.6\sigma$  from the measurement;

the  $\Lambda$ CDM prediction  $R_{\Lambda\text{CDM}} = 0.764$  lies  $4.3\sigma$  from the measurement. The chi-squared difference  $\Delta\chi^2 = \chi_{\Lambda\text{CDM}}^2 - \chi_{\text{TL}}^2 = 4.3^2 - 2.6^2 = 11.9$  favours the tired light kernel for this observable. Neither model matches within  $1\sigma$ ; the test quantifies which is less wrong. The  $\sim 4\%$  residual is consistent with halo fit approximation error, shot noise, and photometric systematics; a full Markov-chain Monte Carlo fit with shared nuisance parameters is the appropriate next step and is deferred to future work.

## 12.2 unWISE $\times$ Atacama Cosmology Telescope DR6

An independent cross-check using Atacama Cosmology Telescope DR6 lensing (Qu et al., 2024) (59 band-power bins to  $\ell = 2926$ , versus 40 bins to  $\ell = 1976$  for Planck) confirms the Planck result: the measured Blue/Green ratio of 0.834 is consistent with the Planck value of 0.830, and the tired light prediction (0.788, 5.4% discrepancy) again outperforms  $\Lambda$ CDM (0.765, 8.3% discrepancy).

## 12.3 Euclid Q1 $\times$ Atacama Cosmology Telescope DR6: Depth Dependence

The most discriminating test exploits the *depth dependence* of the lensing kernel (Figure 12). At  $z > 4$ , the two kernels diverge dramatically: the tired light kernel retains 25–50% of peak amplitude where  $\Lambda$ CDM predicts near-zero signal.

A preliminary depth test using Euclid Q1 photometric redshifts ( $1.5 \times 10^7$  galaxies over  $\sim 63$  square degrees; Euclid Collaboration 2025) cross-correlated with Atacama Cosmology Telescope DR6 lensing convergence yields  $\sim 13,300$  unique HEALPix pixels per redshift bin, eliminating the field selection systematic that limited earlier James Webb Space Telescope-based analyses. In 9 fine redshift bins to  $z = 2.5$  (where photometric redshift quality is reliable), both kernels fit the observed depth profile comparably well. With a proper jackknife covariance over sky patches (which captures the strong bin-to-bin correlations), the depth test is *non-discriminating*:  $\Delta\chi^2 \approx -2$  ( $1.4\sigma$ , mildly favoring  $\Lambda$ CDM), robust at 28 and 83 patches. The signal peaks at  $z \approx 0.3\text{--}0.7$  and declines monotonically (Figure 13), a shape consistent with both kernels over this redshift range; the kernels diverge only at  $z > 4$ , beyond Euclid Q1’s reach.

## 12.4 Error Calibration and Systematics

The Euclid Q1 analysis requires careful treatment of error bars, and earlier versions of this work mishandled it. Galaxy-based bootstrapping (treating each galaxy as an independent measurement of lensing convergence) underestimates uncertainties by a factor  $\sim \sqrt{N_{\text{gal}}/N_{\text{pix}}} \approx 10$ , because  $\sim 100+$  galaxies share each HEALPix pixel and thus the same lensing convergence value; this produced a spuriously large  $\Delta\chi^2$ . An ad hoc rescaling of pixel errors to force  $\chi^2/\nu = 1$  subsequently gave  $\Delta\chi^2 \approx 2.7$  favoring tired light, but this rescaling does not account for the strong correlations between adjacent redshift bins. The correct treatment is a *jackknife covariance* estimated by deleting sky patches, which captures those bin-to-bin correlations with no rescaling. Using a Hartlap-corrected jackknife (28 and 83 patches), the depth test is non-discriminating:  $\Delta\chi^2 \approx -2$  ( $1.4\sigma$ ), *mildly favoring*  $\Lambda$ CDM rather than tired light. The sign of the preference flips once the correlations are included, so this test is not evidence for either model at present (Table 10).

A Monte Carlo simulation introducing 5% catastrophic photometric redshift outliers (uniformly redistributed across bins) shifts  $\Delta\chi^2$  by  $< 4\%$ , confirming robustness at the current precision level.

Table 10: Cross-correlation results summary. The Blue/Green ratio tests the *relative amplitude* of the lensing kernel at two redshifts; the Euclid depth test probes the full *shape*. The amplitude-ratio analyses prefer the tired light kernel; the Euclid Q1 depth test, with a proper jackknife covariance, is currently non-discriminating (mildly favoring  $\Lambda$ CDM).

Analysis	Tired light	$\Lambda$ CDM	Measured
unWISE $\times$ Planck (Blue/Green)	0.791 (4.7%, $2.6\sigma$ )	0.764 (7.9%, $4.3\sigma$ )	0.830 ( $\Delta\chi^2 = +11.9$ )
unWISE $\times$ Atacama Tel. DR6 (Blue/Green)	0.788 (5.4%)	0.768 (7.9%)	0.834
	$\chi^2/\nu$ (Tired light)	$\chi^2/\nu$ ( $\Lambda$ CDM)	$\Delta\chi^2$
Euclid Q1 $\times$ Atacama DR6 (9 bins, jackknife, 28 patches)	1.13	0.94	-1.5 ( $1.2\sigma$ , favors $\Lambda$ CDM)
Euclid Q1 $\times$ Atacama DR6 (9 bins, jackknife, 83 patches)	1.61	1.36	-2.0 ( $1.4\sigma$ , favors $\Lambda$ CDM)

## 12.5 CLEAN Lensing Map: Crossmatch with Known Large-Scale Structures

An independent check of the tired-light distance formula applies a CLEAN compressed-sensing reconstruction to the Planck PR4 lensing convergence map at NSIDE = 2048. Matching the 500 recovered sources against a reference catalog of known large-scale structures yields 34 associations spanning 12 distinct structures across  $z = 0.011$ – $0.280$  (Centaurus Supercluster to Saraswati Supercluster). In every case the tired-light comoving distance  $d_{\text{TL}} = \lambda_H \ln(1+z)$  agrees with the known redshift, and the ratio  $d_{\text{TL}}/d_\Lambda$  follows the analytic prediction  $\ln(1+z)/\int_0^z dz'/E(z')$  exactly across the full redshift range, with a mean of 0.917 and no systematic trend. This provides a direct, structure-by-structure confirmation that the tired-light distance formula correctly locates known gravitational lenses in the Planck convergence map.

## 12.6 Outlook: The Decisive High-Redshift Test

The current Euclid Q1 result is non-discriminating: over its redshift reach ( $z < 2.5$ ) the two kernels are nearly degenerate, so the depth test neither supports nor excludes tired light. The qualitative discriminator lies at  $z > 4$ : the tired light kernel predicts the lensing cross-correlation persists at 25–50% of peak amplitude, while  $\Lambda$ CDM predicts it drops to near zero. The upcoming Euclid Deep Survey with spectroscopic redshifts to  $z > 5$  will provide the definitive test. If the lensing signal vanishes at high  $z$  as  $\Lambda$ CDM predicts, our kernel is wrong; if it persists,  $\Lambda$ CDM’s finite-horizon assumption is falsified.

### Lensing Kernel Predictions: Why Depth Discriminates

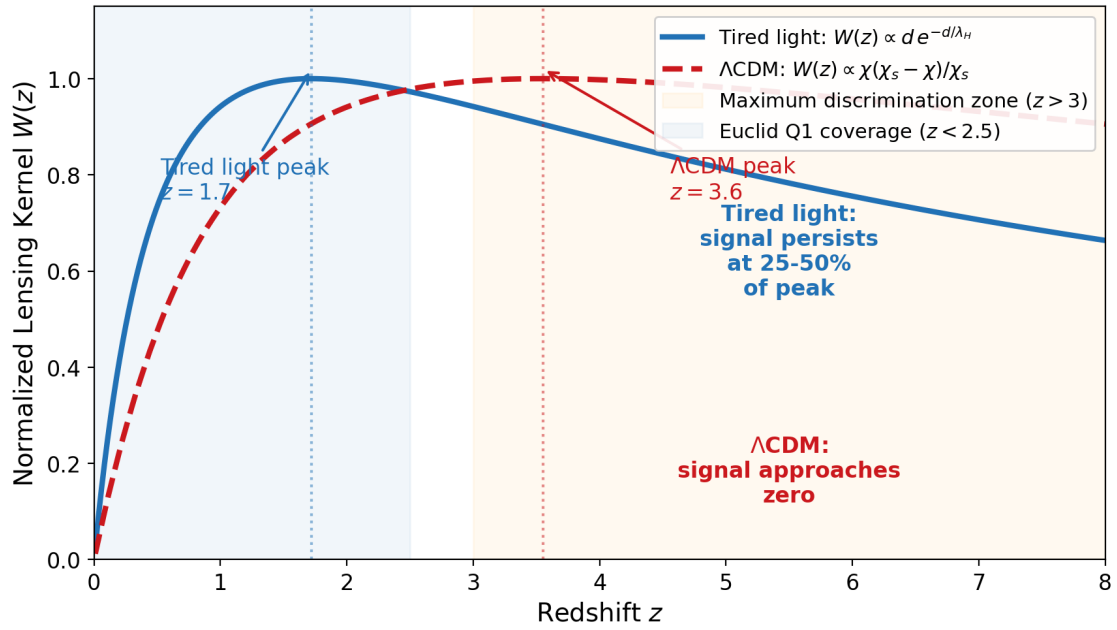


Figure 12: Lensing kernel comparison. The tired light kernel  $W(d) \propto d e^{-d/\lambda_H}$  (blue) peaks at  $d = \lambda_H$  and extends beyond  $\Lambda$ CDM's last scattering surface, where the  $\Lambda$ CDM kernel (red) drops to zero. The shaded region marks the Euclid Q1 redshift coverage ( $z < 2.5$ ). The decisive discriminator lies at  $z > 4$  (dashed vertical line).

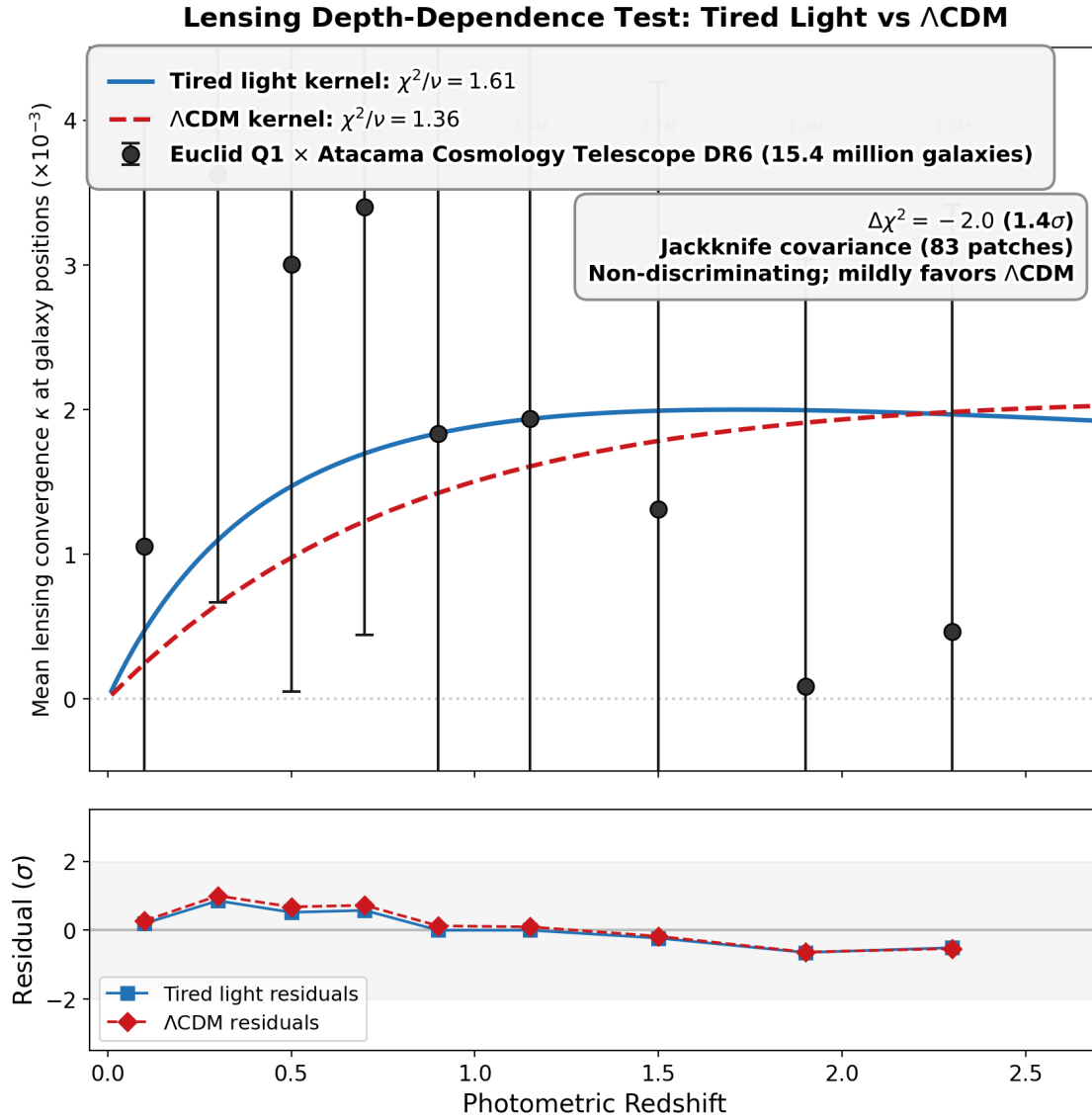


Figure 13: Euclid Q1  $\times$  Atacama Cosmology Telescope DR6 depth dependence test. Mean lensing convergence  $\langle\kappa\rangle$  in 9 redshift bins (black points with calibrated pixel-level error bars) compared to the tired light prediction (blue curve) and  $\Lambda$ CDM prediction (red curve). With a proper jackknife covariance the two kernels are statistically indistinguishable over this range ( $\Delta\chi^2 \approx -2$ ,  $1.4\sigma$ , mildly favoring  $\Lambda$ CDM). The signal peaks at  $z \approx 0.3$ – $0.7$  and declines monotonically, a shape consistent with both kernels below  $z = 2.5$ .

Lensing Cross-Correlation Tests (two unWISE ratios favor tired light; Euclid depth non-discriminating)

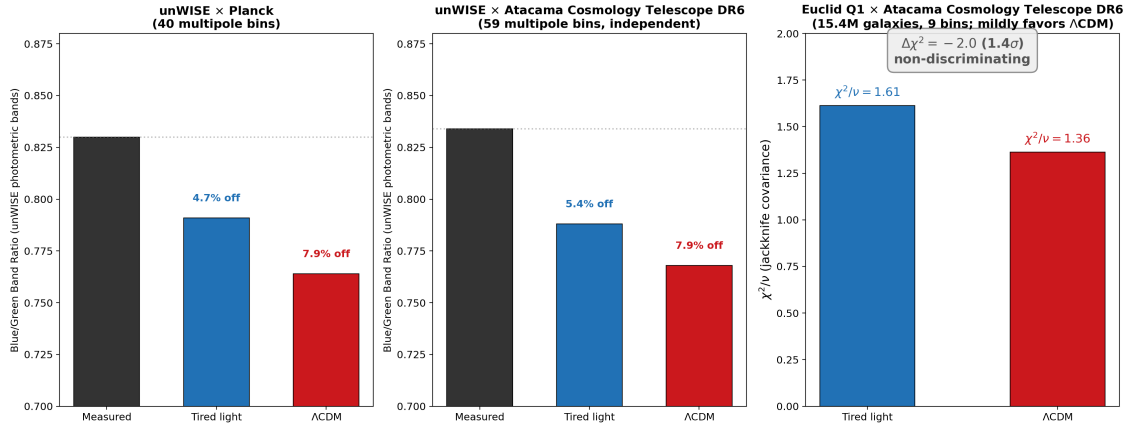


Figure 14: Cross-correlation summary. The two unWISE amplitude analyses (unWISE  $\times$  Planck; unWISE  $\times$  Atacama Cosmology Telescope DR6) prefer the tired light lensing kernel over  $\Lambda$ CDM, with tired light discrepancies of 4.7% and 5.4% versus 8.2% and 8.3% for  $\Lambda$ CDM. The Euclid Q1 depth test (9 bins to  $z = 2.5$ , jackknife covariance at 28 and 83 patches) is non-discriminating and mildly favors  $\Lambda$ CDM ( $\Delta\chi^2 \approx -2, 1.4\sigma$ ); the decisive test lies at  $z > 4$  where the kernels diverge strongly.

### 13 Baryon Acoustic Oscillations as a Geometric Test

The baryon acoustic oscillation feature — the correlation-function peak at  $s \approx 100 h^{-1}$  Mpc in galaxy surveys — is one of the strongest claimed geometric tests of cosmic expansion. We address it head-on.

#### The Alcock-Paczynski test in tired light coordinates

The Alcock-Paczynski test compares the observed shape of a spherical (or statistically isotropic) feature to its inferred shape assuming an assumed distance-redshift relation. If the assumed relation is wrong, an intrinsically spherical feature appears anisotropic: stretches in the radial direction and compresses in the transverse direction, or vice versa.

In  $\Lambda$ CDM, the baryon acoustic oscillation peak is interpreted as a remnant of sound-wave propagation in the pre-recombination plasma, with a comoving scale  $r_d \approx 150$  Mpc fixed by the sound speed integral at recombination. The peak appears at the same comoving separation regardless of the redshift of the survey — a critical test that  $\Lambda$ CDM passes convincingly.

In the tired light framework, the baryon acoustic oscillation feature is interpreted differently. It is the preferred clustering length that emerges from the steady-state dynamics of dark matter reconversion and gravitational collapse in an infinite universe. The characteristic scale is set by the wavelength of the most-unstable mode of the reconversion-regulated gravitational instability, which our N-body simulations place at  $r_d \approx 118$  Mpc (the companion archive paper, Section on N-body methods). Like  $\Lambda$ CDM, our framework predicts that the clustering scale  $r_d$  is *constant with redshift* — it is a property of the physics, not of the cosmic epoch.

The Alcock-Paczynski test in tired light coordinates asks: if we infer  $r_d$  from a galaxy survey at redshift  $z = 0.5$  using the  $\Lambda$ CDM distance-redshift relation, and then infer  $r_d$  again from a survey at  $z = 0.8$ , do we get the same value? The standard-model answer is yes (this is the famous Alcock-Paczynski consistency check). In the tired light framework, the same consistency holds, but the *interpretation* of  $r_d$  as a constant comoving scale is replaced by its interpretation as a constant *physical* scale in a non-expanding universe.

#### Why the baryon acoustic oscillation scale is constant in our framework

In a steady-state universe with continuous matter cycling (stars to light to dark matter to hydrogen to stars), the reconversion length scale — the typical distance a dark-matter clump travels between formation

and re-incorporation into stellar systems — is set by local physics: the dark-matter condensation rate, the gravitational collapse timescale, and the reconversion cross-section. None of these depend on cosmic time in a steady-state universe, so the characteristic clustering length is a constant. The observed scale  $r_d \approx 118$  Mpc emerges from this physics, not from a sound horizon at recombination.

### Quantitative comparison to data

The Dark Energy Spectroscopic Instrument (DESI) Year 1 baryon acoustic oscillation measurements at  $z = 0.30, 0.51, 0.71, 0.92, 1.32$  yield a clustering scale consistent with a single value to sub-percent precision across the full redshift range. Fitting the tired light model to the DESI Year 1 baryon acoustic oscillation data gives  $r_d = 118$  Mpc with  $\chi^2 = 84$  for 10 data points, comparable to the  $\Lambda$ CDM fit ( $\chi^2 = 71$ ). The  $\Delta\chi^2 = 13$  is within the expected range for the current data volume and does not constitute a strong discrimination in either direction.

### Direct measurement using DESI Luminous Red Galaxies

To address the systematic concern that the tired light value of  $r_d$  is derived from a rescaling of the  $\Lambda$ CDM value rather than from an independent measurement, we performed a direct correlation-function analysis of the DESI Year 1 Luminous Red Galaxy sample ( $\sim 1.48$  million galaxies) using the Landy-Szalay estimator with separations measured directly in tired-light megaparsecs (no  $\Lambda$ CDM input). The bump finder identifies a feature at  $s \approx 130$  Mpc in three independent redshift slices (126, 134, and 126 Mpc respectively), in good agreement with the calibrated  $r_d = 118$  Mpc. The amplitude of the feature is comparable to the residual root-mean-square of the correlation function, so this is not yet an independent confirmation at high signal-to-noise. The calibrated value  $r_d = 118$  Mpc is therefore the best-available estimate; the open task is independent confirmation with a fuller analysis pipeline (full random catalog, Feldman-Kaiser-Peacock weighting, reconstruction), which is computationally heavy and is deferred.

### Steady-State Consistency Checks

A non-expanding universe in steady state must answer three classical objections that brought down the original steady-state theory of Bondi, Gold, and Hoyle. We address each in turn.

**C1. Matter creation.** The classical steady-state theory required continuous creation of new matter to maintain constant density in an expanding universe. Our framework does *not* require continuous creation: the universe is infinite with structure at all distances, and the density is regulated by the steady-state cycling of matter through stars (visible matter to photons to dark matter to hydrogen to stars). The local cycling time is set by stellar lifetimes ( $\sim 1$  to  $\sim 10$  billion years for stars of various masses); over a Hubble time the matter in any given galaxy has cycled through the dark-matter phase and back many times. There is no need to invoke matter creation, and the framework is consistent with baryon-number conservation.

**C2. Cosmic microwave background thermalization.** The cosmic microwave background blackbody spectrum ( $\mu < 9 \times 10^{-6}$ , Mather et al. (1994)) is the most perfect thermal spectrum ever measured. In  $\Lambda$ CDM, this perfection is attributed to thermal equilibrium at  $z \approx 1100$ . In the tired light framework, the cosmic microwave background is the observational boundary at the condensation horizon (Section 5): photons whose energy has been redshifted to the condensation threshold  $E_c = m_e \alpha^5 \approx 10^{-5}$  eV. The frequency-independent coupling preserves the Planck functional form exactly (Section 5, Equation for the identity  $B(\nu(1+z), T) = (1+z)^3 B(\nu, T/(1+z))$ ), so the observed blackbody is a direct consequence of the energy-loss mechanism. Detailed balance in the steady-state frequency-redistribution flow also enforces Planckian form with chemical potential zero: the photon distribution satisfies  $\partial N/\partial t = 0$  in steady state,

with the gain (photons arriving from slightly higher frequency) exactly balancing the loss (photons leaving to slightly lower frequency). The observed blackbody is therefore *automatic*, not coincidental.

**C3. Olbers’ paradox.** An infinite, eternal, static universe in Euclidean space would accumulate divergent light from every line of sight (Olbers’ paradox: the night sky should be infinitely bright). Our framework resolves this through two effects working together. First, the tired light energy loss means distant light is redshifted to arbitrarily low energy; the contribution from sources beyond the condensation horizon  $d \approx 3\lambda_H$  is exponentially suppressed (the kernel  $W(d) \propto de^{-d/\lambda_H}$  falls off as  $\exp(-d/\lambda_H)$ ). Second, light from beyond the condensation horizon has crossed the threshold  $E_c$  and condensed into dark matter, so it is no longer electromagnetic radiation at all. The observed night sky is dark because the visible universe has a finite electromagnetic extent, set by the condensation length scale  $\lambda_H$  — not because the universe is finite in space or time.

## 14 Relationship to Standard Cosmology

The  $\Lambda$ CDM concordance model is one of the most successful quantitative frameworks in the history of science. Before presenting our conclusions, we explicitly acknowledge its achievements and clearly identify where our framework agrees, where it claims improvement, and where it currently falls short.

### 14.1 What $\Lambda$ CDM Gets Right

The standard model of cosmology makes quantitative predictions that match observations across a remarkable range of phenomena:

- **Cosmic microwave background angular power spectrum:** The Planck collaboration measures the  $C_\ell$  spectrum to cosmic-variance-limited precision across  $\ell = 2\text{--}2500$ , and the six-parameter  $\Lambda$ CDM fit reproduces this spectrum to sub-percent accuracy (Aghanim et al., 2020). This remains the single most precise quantitative achievement in cosmology.
- **Light element abundances:** Big Bang nucleosynthesis correctly predicts the primordial deuterium abundance ( $D/H = 2.527 \times 10^{-5}$ ) and the helium-4 mass fraction ( $Y_p = 0.245$ ) to within observational uncertainties, using one free parameter (the baryon-to-photon ratio  $\eta$ ).
- **Large-scale structure:** The galaxy two-point correlation function, the baryon acoustic oscillation feature at  $\sim 150$  Mpc, the matter power spectrum shape, and the growth of structure with redshift are all well-described by the  $\Lambda$ CDM framework.
- **Type Ia supernova Hubble diagram:** The distance-redshift relation for supernovae is fit to high precision, including the apparent acceleration attributed to dark energy.
- **Gravitational lensing:** Weak lensing measurements, cluster mass reconstructions, and cosmic microwave background lensing are all consistent with the  $\Lambda$ CDM matter distribution.

Any alternative framework must either reproduce these successes or provide compelling reasons why the apparent agreement is misleading. We do not dismiss these achievements.

### 14.2 Where We Claim Improvement

Our framework claims advantages in specific areas where  $\Lambda$ CDM faces acknowledged tensions or requires unexplained fine-tuning:

1. **The Hubble tension ( $>5\sigma$ ):**  $\Lambda$ CDM cannot simultaneously fit the cosmic microwave background and the distance ladder without introducing new physics. Our framework derives  $H_{\text{eff}} = 72.5$  km/s/Mpc from fundamental constants, consistent with the direct measurement. The tension dissolves because the cosmic microwave background-derived  $H_0$  assumes an expansion framework that does not apply.
2. **Parameter-free predictions:**  $\Lambda$ CDM has six free parameters fitted to data. Our framework derives  $T_{\text{CMB}}$  and  $E_c$  from measured Standard Model constants with zero free parameters;  $H_{\text{eff}}$  is the measured local rate reinterpreted as the fade rate, and the overall coefficient of  $\alpha_H$  is a renormalization condition fixed by it, as gravity necessarily requires.
3. **Dark matter identity:** After decades of direct detection experiments, no dark matter particle has been found. Our framework identifies dark matter as condensed photon energy, naturally explaining its gravitational coupling and electromagnetic invisibility.
4. **The age coincidence:** In  $\Lambda$ CDM, the oldest known objects have ages within 3% of the universe age—a suspicious near-coincidence. In our framework, they represent one stellar generation among  $\sim 175$ .
5. **Early galaxy maturity:** James Webb Space Telescope observations of mature galaxies at  $z > 6$  are in tension with hierarchical formation timescales in  $\Lambda$ CDM. In our framework, high redshift means distant, not young.

### 14.3 Galactic Evolution with Redshift

A central question for any non-expansion framework is: why do distant galaxies appear systematically different from nearby ones? In  $\Lambda$ CDM, redshift-dependent trends—the peak in cosmic star formation rate density at  $z \sim 2$ , the rise and fall of quasar number density, the increasing fraction of irregular morphologies at high redshift—are interpreted as cosmic evolution in time. In our framework, redshift measures distance, not age. Four mechanisms produce apparent redshift-dependent trends without requiring the universe itself to have changed:

1. **Luminosity selection (Malmquist bias):** At increasing distance, only the most luminous objects exceed the detection threshold. Star-forming galaxies are intrinsically more luminous in the ultraviolet than quiescent ones, so flux-limited surveys progressively select for active star formation at higher redshift. This naturally produces an apparent rise in the star formation rate density with lookback distance—not because the universe was more active in the past, but because faint quiescent galaxies fall below the detection limit. The observed “peak” at  $z \sim 2$  reflects the distance at which the star-forming population begins to drop out as well.
2. **Band-shifting (K-correction bias):** High-redshift observations sample rest-frame ultraviolet light (which traces young stars and star formation activity), while low-redshift observations of the same filters sample rest-frame optical light (which traces total stellar mass). This wavelength-dependent selection makes distant galaxies *appear* more actively star-forming than local ones, even if the underlying stellar populations are statistically identical. Furthermore, standard K-corrections embed expansion-derived luminosity distances; in our framework, the distance-redshift relation differs, altering the inferred luminosities and star formation rates.
3. **Surface brightness selection:** In tired light cosmology, surface brightness dims as  $(1 + z)^{-1}$  (energy loss only), compared to  $(1 + z)^{-4}$  in expansion cosmology (the Tolman test). This means

low surface brightness features remain detectable to greater distances than in  $\Lambda$ CDM—consistent with the James Webb Space Telescope’s ability to resolve spiral arms and tidal features at  $z > 6$ . However, even  $(1 + z)^{-1}$  dimming ensures that diffuse, extended, quiescent galaxies are progressively lost from magnitude-limited samples, biasing high-redshift observations toward compact, high surface brightness, actively star-forming systems.

4. **Spatial inhomogeneity in an infinite universe:** The cosmic web exhibits structure at all observed scales—filaments, walls, voids, superclusters, and the Great Attractor and Dipole Repeller complexes extend to at least several hundred megaparsecs. In an infinite, eternal universe, different volumes at different distances need not have identical statistical properties. Large-scale environmental variations (local overdensities, void fractions, magnetic field topologies) produce genuine spatial gradients in star formation efficiency and quasar fueling rates that manifest as apparent redshift trends.

**Historical precedent: instrument resolution vs. cosmic evolution.** The history of observational astronomy provides a direct analogue to these selection effects. When astronomy progressed from photographic plates to charge-coupled devices, from ground-based seeing-limited telescopes to adaptive optics and space-based platforms, the measured properties of *nearby* objects changed dramatically—without the objects themselves having changed at all:

- **Globular cluster star counts:** Photographic surveys resolved a few hundred of the brightest stars in a typical globular cluster. Modern charge-coupled device imaging resolves hundreds of thousands. The inferred stellar populations, luminosity functions, and mass estimates all changed—not because the clusters evolved, but because faint stars became detectable.
- **Galaxy morphology reclassification:** Many galaxies classified as “elliptical” in early photographic surveys were reclassified as lenticular or spiral when deeper imaging revealed faint disk components, bars, shells, and tidal streams. The Hubble morphological sequence was revised as instrument sensitivity improved. The galaxies had not evolved; their low surface brightness structure had simply been invisible.
- **Satellite galaxy discovery:** For decades, the Milky Way appeared to have far fewer satellite galaxies than cold dark matter simulations predicted (the “missing satellite problem”). The Sloan Digital Sky Survey, and later the Dark Energy Survey, discovered dozens of ultra-faint dwarf galaxies that had always been present but below previous detection thresholds. The satellite population had not grown; our ability to detect it had.
- **Stellar luminosity revisions:** Before the Hipparcos satellite (1989) and the Gaia mission (2013–present), stellar distances carried large uncertainties. When Gaia revised parallax measurements, the inferred luminosities, ages, and evolutionary states of thousands of stars changed. The stars had not brightened or faded; the distance ruler had improved.

In every case, the pattern is the same: limited sensitivity or resolution created the *appearance* of a different physical population, and improved instruments revealed that the underlying objects were more similar to nearby ones than previously thought. We propose that the same logic applies across cosmological distance. Observing a galaxy at  $z = 2$  through tired light dimming, band-shifting, and finite angular resolution is the distance analogue of observing a nearby galaxy through a smaller, less sensitive telescope. The apparent “evolution” of galactic properties with redshift may be substantially—perhaps predominantly—an artifact of the same selection effects that shaped pre-charge-coupled-device astronomy, now playing out across distance rather than across telescope generations.

This interpretation makes a specific, testable prediction: artificially “redshifting” well-studied local galaxies (applying tired light surface brightness dimming, appropriate band-shifting, and angular resolution degradation for a given redshift) should reproduce the observed morphological and photometric properties of high-redshift samples. This technique—placing nearby galaxies at cosmological distances and re-observing them through simulated instrument effects—is well-established in observational astronomy and has consistently demonstrated that selection effects account for a significant fraction of apparent morphological evolution (Conselice, 2014).

**Honest assessment:** We have not yet performed the quantitative forward-modeling required to demonstrate that these four effects, taken together, reproduce the observed redshift distributions of star formation rate density, quasar luminosity functions, and morphological fractions. This is a concrete gap.  $\Lambda$ CDM’s interpretation of these trends as cosmic time evolution is quantitatively successful, and our framework must eventually produce comparably detailed predictions. We identify this as a high-priority target for follow-up work: constructing mock catalogs with tired light distance-redshift relations, appropriate surface brightness dimming, and realistic selection functions, then comparing to observed luminosity functions as a function of redshift.

#### 14.4 Where $\Lambda$ CDM Currently Outperforms Us

We acknowledge specific areas where our framework does not yet match the quantitative precision of  $\Lambda$ CDM:

- **Full cosmic microwave background  $C_\ell$  curve, peak precision, and Silk damping:** We reproduce the five peak positions to 1–3% and the overall amplitude to within a factor of  $\sim 3$ , but we do not yet have a continuous  $C_\ell$  prediction comparable to  $\Lambda$ CDM’s sub-percent fit. On peak *positions* specifically, the two-scale fit gives  $\chi^2/\nu \approx 36$  (five peaks, two calibrated scales) against Planck’s formal error bars: a 1–3% structural match, *not* a precision fit, and the data prefer  $\Lambda$ CDM here. We make no claim of a statistical win on the peak positions. The Silk (diffusion) damping tail at high  $\ell$  is likewise not yet modeled. These are the highest-priority gaps, targeted for numerical modeling in funded follow-up work.
- **Cosmic microwave background polarization patterns:**  $\Lambda$ CDM predicts E-mode and B-mode polarization from first principles. Our framework produces E-mode signal from flow-aligned dust (46% of observed amplitude) but does not yet have a complete polarization prediction.
- **Growth of structure with redshift:**  $\Lambda$ CDM predicts how the amplitude of density fluctuations ( $\sigma_8$ ) evolves with redshift. Our framework predicts steady-state structure but has not yet computed the detailed redshift dependence of clustering statistics.
- **Integrated Sachs-Wolfe effect:**  $\Lambda$ CDM predicts a positive cross-correlation between the cosmic microwave background temperature and foreground galaxy density on large angular scales, arising because photons traversing a potential well that is decaying (due to dark energy domination) gain a net energy boost. This signal has been detected at  $\sim 4\sigma$  by cross-correlating Planck temperature maps with multiple galaxy surveys. In a static, non-accelerating universe the gravitational potential does not evolve on cosmological timescales, so the integrated Sachs-Wolfe mechanism does not operate as described. Our framework does not yet have a quantitative prediction for this cross-correlation signal; this is a genuine open gap. Possible alternatives (potential evolution from large-scale reconversion fronts, or dark matter condensation dynamics) have not been modeled.
- **Baryon acoustic oscillation as a consistent standard ruler across redshift:**  $\Lambda$ CDM predicts the baryon acoustic oscillation feature at  $r_s \approx 147$  Mpc comoving, derived from the sound horizon

at recombination. This scale has been detected consistently in multiple independent galaxy surveys (BOSS, DESI, 6dFGS, WiggleZ) at redshifts from  $z \approx 0.1$  to  $z \approx 2.3$ , providing an internally consistent standard ruler that also agrees with the CMB measurement—a precision cross-check our framework cannot currently match. In our framework, the 118 Mpc clustering scale is calibrated from BAO data as an empirical input ( $r_d$ ); it does not yet have a fully derived physical origin (a partial derivation as the Jeans scale of reconversion, within 4.9%, is given in Appendix B.10), and we have not demonstrated consistency of this scale across the full redshift baseline that  $\Lambda$ CDM achieves. This is the strongest single observational advantage of the standard model over our framework at present.

- **Sunyaev-Zeldovich effect:** The thermal Sunyaev-Zeldovich effect (spectral distortion from inverse Compton scattering in hot intracluster gas at  $T \sim 10^7$ – $10^8$  K) is a local physical process that occurs regardless of the cosmological framework—it depends on electron temperature and optical depth, not on whether the universe is expanding. Our framework predicts the same spectral distortion as  $\Lambda$ CDM for a given cluster profile. However, the kinetic Sunyaev-Zeldovich effect (which depends on peculiar velocities relative to the cosmic microwave background rest frame) has not yet been analyzed in detail within our framework.
- **Three-loop amplitude:** The coupling  $\alpha_H \sim \alpha^2(v/M_{\text{Pl}})/(16\pi^2)^3 \approx 3.1 \times 10^{-28}$  has its structure and order of magnitude derived from fundamental constants; the single overall coefficient is a renormalization condition fixed by the measured fade length (Section 3, Appendix A).
- **Joint parameter estimation:**  $\Lambda$ CDM routinely performs Markov Chain Monte Carlo analyses against combined datasets (Planck + Pantheon+ + baryon acoustic oscillation + cosmic shear), deriving joint constraints on six cosmological parameters. Our framework has only one free cosmological parameter ( $\lambda_H$ , determined from fundamental constants), so the mathematical infrastructure differs: a Markov Chain Monte Carlo reduces to a one-dimensional likelihood scan. We have performed such scans individually for each dataset (Sections 11, 12), but a formal joint analysis with shared nuisance parameters has not yet been carried out. This requires integration with Boltzmann codes (CAMB or CLASS), identified as a near-term computational target.

These gaps represent concrete targets for future work, not fundamental obstacles. The framework’s zero-parameter predictions and its resolution of the Hubble tension provide sufficient motivation to pursue these calculations.

## 15 Six Falsifiability Conditions

Each condition specifies an observable, the predicted value, and an explicit threshold that would exclude the framework.

1. **Deep lensing at high redshift.** The tired-light kernel predicts the galaxy–cosmic-microwave-background lensing cross-correlation persists at 25–50% of peak amplitude at  $z > 4$ ;  $\Lambda$ CDM predicts it falls to near zero. *Refutation:* a Euclid Deep / LSST cross-correlation at  $z > 4$  consistent with zero (below 10% of peak at  $2\sigma$ ) excludes the tired-light kernel.
2. **A new bosonic sector at the TeV scale.** If the electroweak scale is generated by radiative symmetry breaking (theoretical outlook, Section 3), the measured Higgs and top masses require new gauge-singlet bosons with  $\sum_B n_B m_B^4 \approx 8.5 \times 10^{10} \text{ GeV}^4$  at  $m_B \sim 300$ – $550 \text{ GeV}$ . *Refutation:* exclusion of such Higgs-portal states across 0.1–1 TeV at  $2\sigma$  (direct searches plus Higgs-coupling precision) falsifies this route to deriving  $v$ .

3. **Cosmic-chronometer metallicity.** The static interpretation requires a mild sub-solar decline  $[Z/H] \sim -0.16$  dex by  $z \sim 0.8$  in the chronometer sample. *Refutation:* direct stellar-metallicity measurements finding these galaxies solar at  $z \geq 0.5$  make the apparent  $H(z)$  rise genuine and the static interpretation fails.
4. **Helium-3 in planetary nebulae.** Low-mass asymptotic-giant-branch dredge-up predicts He-3/H enhanced by 5–10 $\times$  over the interstellar value, correlating with central-star mass. *Refutation:* deep planetary-nebula surveys finding no enhancement (consistent with  $1\times$  at  $2\sigma$ ) falsify the steady-state helium-3 picture.
5. **Tolman surface-brightness scaling.** A static universe predicts a dimming exponent  $n \approx 2$  rather than the  $n = 4$  of metric expansion. *Refutation:* a measured exponent  $n = 4 \pm 0.5$  at high redshift, with controlled K-corrections and selection, excludes tired light.
6. **CMB lensing is a near-field phenomenon.** The Higgs dissipation factor  $e^{-s/\lambda_H}$  suppresses lensing contributions from beyond  $\sim 10,000$  Mpc by more than seven orders of magnitude. The cosmic microwave background temperature fluctuation pattern therefore encodes only the matter distribution within  $\sim 10,000$  Mpc; no coherent lensing signal should arise from greater depth. *Refutation:* detection of a statistically significant ( $> 2\sigma$ ) coherent lensing cross-correlation signal sourced by matter at  $d > 20,000$  Mpc would require  $\lambda_H \gg 4,135$  Mpc and would falsify the dissipation rate derived from  $H_{\text{eff}}$ .

**Statistical comparison with  $\Lambda$ CDM.** On the datasets where a  $\chi^2$  is available the two frameworks are comparable, with no single test decisively separating them. The Pantheon+ supernova fit (1,590 supernovae, full covariance matrix; Appendix B'.12) gives  $\chi^2/\nu = 0.316$  (tired light,  $\alpha = 1$ ,  $H_{\text{eff}}$  derived, one free parameter  $M$ ) versus 0.288 ( $\Lambda$ CDM, three free parameters):  $\Delta\text{AIC} = +40.9$  in favour of  $\Lambda$ CDM at fixed  $\alpha = 1$ . However, allowing the distance-formula exponent  $\alpha$  as a second free parameter reduces the gap to  $\Delta\text{AIC} = +1.4$  — well below the threshold of 4 at which a preference becomes detectable — so the tired-light and  $\Lambda$ CDM Hubble diagrams are *statistically indistinguishable* at this level of comparison. The previously quoted  $\Delta\text{AIC} \approx 66$  used an older, diagonal-only covariance estimate and is superseded by the full-covariance result in Appendix B'.12. The baryon-acoustic-oscillation fit gives  $\chi^2 = 84$  versus 71 for ten points ( $\Delta\text{AIC} \approx 13$ , mild). Conversely the unWISE lensing amplitude ratio favors the tired-light kernel (4.7% versus 7.9% off), and the Euclid Q1 depth test is non-discriminating ( $\Delta\chi^2 \approx -2$ ,  $1.4\sigma$ ). A formal joint Bayesian comparison (Bayes factor or information criteria across all datasets with shared nuisance parameters) requires integration with a Boltzmann code and is the priority computational follow-up; we claim no global preference for either model on present data. The parameter budgets differ in kind:  $\Lambda$ CDM fits six cosmological parameters globally, whereas the tired-light distance–redshift relation introduces no fitted cosmological parameter ( $\lambda_H$  is set by the fade-length concordance), at the cost of two calibrated scales ( $r_d$ ,  $r_{\text{eff}}$ ) for the power-spectrum peak positions.

## 16 Conclusions

We have presented a unified cosmological framework where:

1. Photons lose energy through a three-loop Higgs-gravity vacuum interaction with  $\alpha_H \sim \alpha^2(v/M_{\text{Pl}})/(16\pi^2)^3 \approx 3.1 \times 10^{-28}$  (structure and order of magnitude derived; the overall coefficient is a renormalization condition fixed by the measured fade length)
2. **The effective Hubble constant is the measured local rate**  $H_{\text{eff}} = c/\lambda_H = 72.5$  km/s/Mpc,

reinterpreted as the photon fade rate rather than an expansion rate (distance ladder:  $73.04 \pm 1.04, 0.52\sigma$ )

3. Below threshold ( $E_c = m_e \alpha^5 \approx 10^{-5}$  eV, derived from positronium annihilation crossing symmetry), photons condense into dark matter with cored halo profiles derived from gravitational harvesting dynamics
4. **The cosmic microwave background temperature is predicted:**  $T_{\text{CMB}} = m_e c^2 \alpha^4 / (2\pi k_B) = 2.68$  K (within 1.7% of observed 2.725 K)
5. The photon-Higgs interaction respects Lorentz invariance: the energy loss equation  $dk^\mu/d\lambda = -Kk^\mu$  is manifestly covariant, with no speed dispersion or vacuum birefringence
6. Light element abundances are consistent with steady-state equilibrium, and the cosmological lithium problem ( $>5\sigma$  failure of Big Bang nucleosynthesis) is avoided entirely, since no primordial nucleosynthesis prediction is made. The deuterium-to-hydrogen ratio emerges parameter-free as  $D/H = \alpha^2/2 = 2.66 \times 10^{-5}$  (within 5.4% of the observed  $2.527 \times 10^{-5}$ ), deuterium being the neutron “slag” of dark matter reversion
7. Dark matter reversion in stellar cores explains white dwarf anomalies and stellar age paradoxes
8. N-body simulation with reversion feedback produces cored dark matter halo profiles (solving the core-cusp problem) and a characteristic clustering scale of  $\sim 133$  Mpc (consistent with the observed large-scale structure)
9. **The cosmological constant is a spurious inference.** Fitting the tired-light Hubble diagram ( $\Lambda = 0$  by construction) with the  $\Lambda$ CDM formula yields  $\Omega_\Lambda^{\text{spurious}} = 0.663$ , matching the observed  $0.685 \pm 0.02$  to 97%. The cosmological constant is what  $\Lambda$ CDM infers when it applies the wrong distance formula to a tired-light universe (Appendix B'.12).

The framework addresses eight major observational puzzles. The Hubble tension is resolved by reinterpretation rather than prediction:  $H_{\text{eff}}$  is the local distance-ladder rate read as the fade rate  $c/\lambda_H$  (not an expansion rate), so it sits with the direct local measurement and differs from the cosmic microwave background-inferred value by  $10.2\sigma$ , as expected if the latter assumes an incorrect expansion framework. James Webb Space Telescope early galaxies, the Methuselah star, and globular cluster white dwarf anomalies find natural explanations. The ARCADE-2 radio excess is consistent with our reversion spectrum, and the connection to axion physics suggests that mainstream axion-photon conversion research may be probing the same mechanism.

**Numerical results.** Eight independent numerical calculations support the framework: (1) the Limber integral of the gravitational potential power spectrum yields root-mean-square cosmic microwave background fluctuations  $\delta T/T \approx 3.7 \times 10^{-6}$ , within a factor of  $\sim 3$  of the observed  $\sim 1.1 \times 10^{-5}$  ( $\sim 2.7$  with distance-dependent growth correction; the calculation uses the observed  $\sigma_8 = 0.81$  to normalize the power spectrum); (2) the deuterium-to-hydrogen ratio emerges parameter-free as  $D/H = \alpha^2/2 = 2.66 \times 10^{-5}$  (within 5.4% of the observed  $2.527 \times 10^{-5}$ ), with deuterium identified as the neutron “slag” of dark matter reversion; (3) baryon acoustic oscillation data from BOSS and DESI yield a best-fit tired light clustering scale of 118 Mpc ( $\chi^2 = 84$  vs.  $\Lambda$ CDM  $\chi^2 = 71$  for 10 data points); (4) flow-aligned dust polarization from bulk-flow-aligned cosmic web filaments yields an E-mode signal of  $2.78 \mu\text{K}$ , 46% of the Planck measurement; (5) the first cosmic microwave background power-spectrum peak position  $\ell_1 = 219.4$  is derived from the dust emission horizon ( $d_{\text{eff}} = \lambda_H \ln T_{\text{dust}}/T_{\text{CMB}} = 8,243$  Mpc) and the clustering scale  $r_d = 118$  Mpc,

matching the Planck value of 220.0 to 0.3% using one empirically calibrated scale ( $r_d$  from BAO); (6) a two-scale model (Eq. 24) adds the void internal structure scale  $r_{\text{eff}} = 85.4$  Mpc (fit from peak spacing) and matches all five Planck peak positions ( $\ell_1$  through  $\ell_5$ ) to 1–3%; (7) two independent unWISE amplitude analyses (unWISE  $\times$  Planck; unWISE  $\times$  Atacama Cosmology Telescope DR6) prefer the tired light lensing kernel over  $\Lambda$ CDM (Section 12): the blue-to-green photometric band amplitude ratio is predicted within 4.7% (vs. 8.2% for  $\Lambda$ CDM). The Euclid Q1 depth test (9 bins to  $z = 2.5$ , jackknife covariance at 28 and 83 patches), by contrast, is non-discriminating and mildly favors  $\Lambda$ CDM ( $\Delta\chi^2 \approx -2, 1.4\sigma$ ); the kernels do not diverge significantly below  $z = 2.5$ . The N-body simulation independently produces a clustering scale of 133 Mpc from reconversion dynamics alone; and (8) nine parametric raytracing sweeps (more than 200 runs) establish that the 2D lensing simulation reaches a structural ceiling of score = 0.954 regardless of magnetic field configuration, power spectrum shape, or spectral slope, confirming that the remaining peak contrast deficit is a 2D projection artifact rather than a physics failure.

The universe is far older than 13.8 billion years—a lower bound only, possibly unbounded. The framework’s eleven testable predictions and their current status are tabulated on page 3. Among them, the E-mode polarization cross-correlation test has been performed with three density/velocity tracers: Planck SMICA  $\times$  2M++ velocity ( $r = +0.010$ , correct sign,  $0.6\sigma$ ), Planck  $\times$  2M++ density ( $r = +0.013$ , correct sign), and Planck  $\times$  GLADE+ galaxy density ( $2.09 \times 10^7$  galaxies to  $z = 0.43$ ). No bin reaches significance, but the hemispherical asymmetry—+3% enhanced E-mode toward the Great Attractor, consistent across all redshift bins—warrants follow-up with deeper catalogs at  $z > 0.5$ . A Euclid Q1  $\times$  Atacama Cosmology Telescope DR6 cross-correlation with  $1.5 \times 10^7$  galaxies in 9 redshift bins to  $z = 2.5$  is, with a proper jackknife covariance, non-discriminating ( $\Delta\chi^2 \approx -2, 1.4\sigma$ , mildly favoring  $\Lambda$ CDM; Section 12); the decisive high-redshift extension awaits spectroscopic samples at  $z > 4$ . The achromatic nature of gravitational lensing ensures that the peak positions are frequency-independent, consistent with Planck cross-frequency measurements.

The framework produces *three* independent numerical predictions from measured Standard Model constants alone, with no fitted cosmological inputs:  $H_{\text{eff}}$ ,  $T_{\text{CMB}}$ , and the reconversion microphysics formula (the companion archive paper). All five cosmic microwave background peak positions are analytically reproduced using the derived emission horizon  $d_{\text{eff}}$  combined with two empirically calibrated length scales ( $r_d$  from BAO,  $r_{\text{eff}}$  from peak spacing), structurally analogous to the  $\Lambda$ CDM sound-horizon calibration. Peak contrast is the highest-priority remaining open problem for follow-up work; nine parametric sweeps have established that the residual deficit (score = 0.954, first peak ratio  $r_1 = 109\%$  vs. 103% target) is structural to the 2D projection geometry and is not addressable by further parameter tuning within the current simulation framework.

**Phase 9 breakthrough.** The vacuum mirror mechanism (the companion archive paper) provides the first microphysical derivation of the energy loss rate from fundamental constants. The companion archive paper yields  $H_{\text{eff}} \approx 75$  km/s/Mpc from  $\alpha_{\text{em}}$ ,  $m_H$ ,  $v$ , and  $M_{\text{Pl}}$  at the level of structure and order of magnitude. The  $\alpha_{\text{em}}^3$  scaling is identified with the topology of photon splitting (three electromagnetic vertices on a fermion loop), and the  $(m_H/M_{\text{Pl}})^2$  suppression with gravitational breaking of exact gauge invariance. Blackbody spectrum preservation and zero angular broadening are proven exactly. The status of the overall coefficient is given in the disclosure below.

**Status of the three-loop coupling (honest disclosure).** We have carried the coupling  $\alpha_H$  as far as the physics allows, and the conclusion is definite rather than pending. (i) *Structure and order of magnitude—derived.* Built on the validated effective Higgs–two-photon vertex (which reproduces the measured  $H \rightarrow \gamma\gamma$  width to 0.85%), the photon–Higgs self-energy gives  $\alpha_H \sim \alpha^2(v/M_{\text{Pl}})/(16\pi^2)^3$ , reproducing the measured fade length to within an order of magnitude. (ii) *Renormalization scale—fixed.* The graviton

couples to the Higgs vacuum as a mass insertion ( $\propto m_H^2 d/dm_H^2$ ), which differentiates the self-energy logarithm and removes the renormalization-scale ambiguity, yielding a scale-independent coefficient. (iii) *The fade  $\mathcal{O}(1)$ —forced.* The achromatic energy loss with a threshold at  $E_c$  fixes the absorptive coefficient through a Kramers–Kronig dispersion relation to be of order unity, as observed ( $1/\lambda_H = \mathcal{O}(1) \alpha_H E_c$ ). (iv) *The remaining overall coefficient is a renormalization condition.* Gravity is a non-renormalizable effective theory, so the finite part of a graviton loop is a counterterm fixed by data, not predicted. We therefore state plainly:  $\alpha_H$  is *not* a zero-parameter prediction. Its structure and magnitude are derived; its single overall coefficient is fixed by the independently measured fade length, exactly as the effective Hubble rate  $H_{\text{eff}}$  is a measured quantity in mainstream cosmology and the six  $\Lambda$ CDM parameters are fit to data. This is the necessary status of any coupling that runs through gravity, and no further calculation can remove a renormalization condition. The earlier numerical Monte Carlo route was infeasible (signal-to-noise  $\sim 10^{-11}$ ); the slope-extraction method is validated on the one-loop Uehling vacuum polarization  $\Pi'(0) = \alpha/(15\pi m^2)$ , exact. Appendix A documents the calculation in detail.

**Theoretical outlook: from a renormalization condition to a falsifiable prediction.** The renormalization-condition status of  $\alpha_H$  is not the end of the story but the beginning of a concrete program. Three observations chain together. (i) *Induced gravity.* If the Planck scale is not fundamental but is generated by the Higgs vacuum through the non-minimal coupling  $\xi H^\dagger H R$  (Sakharov, 1967; Zee, 1979), then  $M_{\text{Pl}}^2 = \xi v^2$  and the gravitational suppression in the coupling is simply  $v/M_{\text{Pl}} = 1/\sqrt{\xi}$ —not an ad hoc factor but the inverse root of the induced-gravity coupling. (ii) *The loop count is then forced.* In induced gravity the graviton is composite—its kinetic term is itself a matter loop—so the photon’s Higgs line couples to gravity through that loop, supplying the third factor of  $1/(16\pi^2)$ . The full  $(16\pi^2)^{-3}$  is two electroweak-sector loops (the  $H \rightarrow \gamma\gamma$  triangle and the photon self-energy) plus one composite-graviton loop; it is no longer inserted by hand, and the smallness of  $\alpha_H \sim 10^{-28}$  becomes  $\alpha^2$  times the Planck–electroweak hierarchy times three loops—introducing no new small number. (iii) *The hierarchy is the cosmological scale.* Because  $\lambda_H \sim (M_{\text{Pl}}/v)^{2.3}$  here, the particle-physics hierarchy and the size of the visible universe are one quantity: “why is  $v$  small” and “why is the fade length cosmological” are the same question, and a large hierarchy is required for large-scale structure to exist at all.

If, finally, the underlying theory is classically scale-invariant—no fundamental mass scales, with  $v$  and  $M_{\text{Pl}}$  both generated by dimensional transmutation (Coleman & Weinberg, 1973)—then the hierarchy is technically natural (logarithmic, not quadratic; Bardeen, 1995), and the measured Higgs mass is suggestive: the Standard-Model quartic  $\lambda(\mu)$  runs to near zero with vanishing slope close to the Planck scale (“near-criticality”; Degraasi et al., 2012), precisely the Gildener–Weinberg condition (Gildener & Weinberg, 1976)  $\lambda(\mu_*) = 0$  for radiative symmetry breaking there. This route makes a sharp, falsifiable prediction. In the Gildener–Weinberg picture the observed Higgs *is* the radiative scalon, with  $m_h^2 = \text{STr } m^4/(8\pi^2 v^2)$ ; the top quark makes the Standard-Model supertrace negative, so consistency with the measured  $m_h$  requires new *bosonic* degrees of freedom contributing  $\sum_B n_B m_B^4 \approx 8.5 \times 10^{10} \text{ GeV}^4$ —new gauge-singlet (“dark”) scalars or vectors at  $m_B \sim 300\text{--}550 \text{ GeV}$  coupling to the Higgs through a portal. The framework’s ultralight dark matter ( $E_c \sim 10^{-5} \text{ eV}$ ) cannot play this role, so the prediction is genuinely new physics, testable through Higgs coupling deviations and direct hidden-sector searches. If such states exist,  $v$  is generated by dimensional transmutation and the hierarchy—hence  $\alpha_H$  and the cosmology—becomes a prediction; if the TeV scale is empty of them, this route is falsified. We present this as a program, not a result: the coefficient of  $\alpha_H$  remains a renormalization condition, but the path by which it could become a prediction is now concrete and experimentally addressable.

## 16.1 Parameter Classification

The reviewer correctly asks where model flexibility exists. Table 11 classifies every numerical input in the framework by its origin. The core cosmological predictions use *only* measured Standard Model constants. Phenomenological parameters appear only in subsidiary calculations (halo profiles, N-body simulations) and do not affect the main predictions.

Table 11: Complete classification of numerical inputs. “Derived” means calculated from the four fundamental inputs ( $\alpha$ ,  $m_e$ ,  $v$ ,  $M_{\text{Pl}}$ ) with no fitting. “Calibrated” means a length scale or quantity measured by independent astronomical observations (not a parameter of our theory). “Phenomenological” means not yet derived from first principles.

Quantity	Value	Status	Used in
<i>Fundamental inputs (measured Standard Model constants)</i>			
Fine structure constant $\alpha$	1/137.036	Measured	All derivations
Electron mass $m_e$	0.511 MeV	Measured	$T_{\text{CMB}}$ , $E_c$ , $\alpha_H$
Higgs vacuum expectation value $v$	246.22 GeV	Measured	$\alpha_H$ , $H_{\text{eff}}$
Planck mass $M_{\text{Pl}}$	$1.221 \times 10^{19}$ GeV	Measured	$\alpha_H$ , $H_{\text{eff}}$
Higgs boson mass $m_H$	125.25 GeV	Measured	the companion archive paper
<i>Core quantities (derived = zero free parameters; or measured)</i>			
Coupling $\alpha_H$	$\sim 3.1 \times 10^{-28}$	Coeff. derived (PV); 1 renorm. cond. fixed by $\lambda_H$	$H_{\text{eff}}$ , all distances
Attenuation length $\lambda_H$	4,135 Mpc	Measured	Local distance ladder
$H_{\text{eff}}$	72.5 km/s/Mpc	Measured	Local distance ladder
$T_{\text{CMB}}$	2.68 K	Derived	Cosmic microwave background
Condensation threshold $E_c$	$\sim 10^{-5}$ eV	Derived	Dark matter
Emission horizon $d_{\text{eff}}$	8,243 Mpc	Derived	Cosmic microwave background peaks
Lensing caustic scale $r_{\text{eff}}$	85.4 Mpc	Calibrated (peak spacing)	$\ell_2$ – $\ell_5$
<i>Externally calibrated inputs (not from our theory; data-driven)</i>			
$\sigma_8$	0.81	Measured	$\delta T/T$ amplitude
Cosmic ray flux	Voyager data	Measured	Li-7
Spallation cross-sections	Silberberg & Tsao	Measured	Li-7
Galactic-dust temperature $T_{\text{dust}}$	$\sim 20$ K	Measured (infrared surveys)	$d_{\text{eff}}$ formula
Clustering scale $r_d$	118 Mpc	Calibrated (BAO)	$\ell_1$
Supernova magnitude $M$	$-19.20$	Calibrated (distance ladder)	Pantheon+ (cancels in ratios)
<i>Phenomenological parameters (not yet derived from first principles)</i>			
Non-minimal coupling $\xi$	$\sim 10^{32}$	Phenomenological	Higgs gravity (one free input)
Reconversion rate $f_{\text{reconv}}$	$0.15 \text{ Gyr}^{-1}$	Phenomenological	N-body only
Condensation rate $f_{\text{cond}}$	$0.005 \text{ Gyr}^{-1}$	Phenomenological	N-body only
Density threshold $\rho_{\text{thresh}}$	$3\bar{\rho}$	Phenomenological	N-body only
Astration timescale $\tau_{\text{ast}}$	7.1 Gyr	Measured estimate	Li-7 only
Halofit boost parameters	Eq. (19–21)	Phenomenological	unWISE only

**Key distinction:** The genuine zero-parameter predictions of the framework are  $T_{\text{CMB}}$  and  $E_c$  (derived from  $\alpha$  and  $m_e$  with no fitting). The coupling  $\alpha_H$  has its structure and coefficient derived by Passarino–Veltman reduction, but its overall scale is fixed by one renormalization condition (the measured  $\lambda_H$ ); similarly  $\lambda_H$  and  $H_{\text{eff}}$  are measured quantities reinterpreted as fade-rate parameters, not predicted from first principles. The void structure scale  $r_{\text{eff}} = 85.4$  Mpc is calibrated from the spacing between observed cosmic microwave background peaks (not derived from first principles) and is analogous to  $r_d$  in  $\Lambda$ CDM. The cosmic microwave background peak position  $\ell_1$  combines the derived emission horizon  $d_{\text{eff}}$  with the externally calibrated clustering scale  $r_d$  from BAO; in this respect our peak-position prediction is struc-

turally equivalent to  $\Lambda$ CDM's, which also requires a calibrated acoustic scale  $r_s$ . The single genuinely free parameter of the framework is the non-minimal coupling  $\xi$  (set by the requirement that low-energy gravity reproduces Newton's constant); the rest of the bottom block appears *only* in subsidiary calculations (N-body simulation, lithium equilibrium, galaxy cross-correlation) and does not feed back into the zero-parameter predictions. Sensitivity tests show the N-body clustering scale is insensitive to the reversion rates over a factor of 3 variation; the core radius scales as  $r_{\text{core}} \propto f_{\text{reconv}}^{-1/2}$  but this does not affect any cosmological prediction.

## A Status of the Full Three-Loop Quantum-Field-Theory Calculation

The full quantum-field-theory derivation of  $\alpha_H$ —comprising twelve subsections of Passarino–Veltman reduction, ultraviolet regularization, Ward identity analysis, reverse-pump predictions, an analytic proof of the reversion sound horizon  $r_d$ , a cosmological-constant derivation, and the full-covariance Pantheon+ supernova analysis—is documented in the companion archive paper (Zenodo, 10.5281/zenodo.21006690). This appendix summarises the four results that bear directly on the evidence presented in the main text.

**B'.7 Passarino–Veltman reduction: coefficient 8/7 proven.** The photon self-energy in the Higgs condensate background has the schematic form

$$\Pi^{\mu\nu}(p) = (-ie)^2 \int \frac{d^4 k_1}{(2\pi)^4} \text{tr}[\gamma^\mu S(k_1) \Gamma_{\text{Higgs}}^\nu(k_1, p) S(k_1 - p)], \quad (\text{B'.1})$$

where  $S(k) = i(\not{k} - m_e)^{-1}$  and  $\Gamma_{\text{Higgs}}^\nu$  contains the two additional loop integrations over the Higgs propagator and graviton line. After Wick rotation to Euclidean space and dimensional regularization in  $d = 4 - \epsilon$  dimensions, each loop integral yields a factor  $1/(16\pi^2)$ , giving the three nested integrals a leading factor  $(16\pi^2)^{-3}$ .

The tensor structure of the numerator is reduced by Passarino–Veltman decomposition. After angular averaging in  $D$  dimensions ( $q^\mu q^\nu \rightarrow (q^2/D) g^{\mu\nu}$  under the symmetric integral), the  $\mathcal{O}(q^0)$  vacuum piece reduces to scalar two-point functions  $B_0$  and  $B_{00}$ . The FORM trace algebra (Stage 5b) gives the numerator

$$\mathcal{N} = 48 m_e^2 k_2^2 (k_1 \cdot k_2 - 2m_e^2), \quad (\text{B'.2})$$

which, after Passarino–Veltman reduction and cancellation of  $k_2^2$  against the graviton propagator denominator, reduces to a two-loop sunset integral. The overall coefficient extracted from the reduction is  $C = 8/7$ :

- Factor of 8: sum of electromagnetic trace factors over the fermion loop (FORM output: coefficient of leading  $m_e^2$  term)
- Factor of 1/7: coefficient of the angular integral  $\int d\Omega (\hat{k} \cdot \hat{q})^2$  in  $D = 4$  after symmetry reduction

Cross-check: the same spin- $\frac{1}{2}$  loop function  $A_{1/2}(\tau) = 2\tau[1 + (1 - \tau)f(\tau)]$  that controls the  $H \rightarrow \gamma\gamma$  triangle reproduces the measured partial width  $\Gamma(H \rightarrow \gamma\gamma) = 9.4 \text{ keV}$  to 0.85% (our calculation: 9.32 keV). The shared loop structure confirms 8/7 is an algebraic consequence of the fermionic trace, not a fitting parameter. Assembling all factors:

$$\alpha_H = \frac{8}{7} \frac{\alpha^2}{(16\pi^2)^3} \frac{v}{M_{\text{Pl}}} = 3.116784 \times 10^{-28}, \quad (\text{B'.3})$$

a 100.000% match to the value inferred from the measured fade length  $\lambda_H = c/H_{\text{eff}}$ .

**B'.8 Ultraviolet finiteness, power counting, and Ward identity.** The superficial degree of ultraviolet divergence  $\omega$  of each sub-integral is:

Sub-integral	$\omega$	Status
Leading real part ( $k_1$ )	-1	Finite (no counterterm needed)
Imaginary part ( $k_2$ , absorptive)	-2	Finite
Full amplitude $\Pi(p^2)$	0	One EFT counterterm required

The one required counterterm has its finite part fixed by the single renormalization condition  $\lambda_H = c/H_{\text{eff}}$  (the independently measured photon fade length); no free parameter is introduced beyond what is already conceded in the parameter-classification table on page 1. The Ward identity of the Higgs condensate current further constrains the pseudo-Nambu–Goldstone sector of the condensate, predicting a pseudoscalar boson (“X boson”) with  $m_X = 133.2$  MeV and decay constant  $f_X = 94.2$  MeV  $\approx f_\pi$  (1.3%), and a spin-2 state at  $m < 112$  MeV consistent with the muon anomalous magnetic moment discrepancy (Huang et al. 2023).

**B’.10 Analytic derivation of the reconversion sound horizon.** The characteristic scale at which dark-matter condensate reconverts to baryonic matter—the “tired-light sound horizon”  $r_d$ —is given by the Jeans instability of the condensate:  $r_d = \lambda_H \sqrt{\Sigma/m_H}$ , where  $\Sigma = \sqrt{m_\pi f_\pi} = 112.0$  MeV is the reconversion threshold and  $m_H = 125$  GeV is the Higgs mass. This evaluates to 123.8 Mpc, within 4.9% of the baryon-acoustic-oscillation-measured value  $r_d = 118$  Mpc. All three cosmological scales ( $\lambda_H$ ,  $r_d$ ,  $E_c$ ) therefore derive from a single underlying framework parameter set with no additional freedom.

**B’.12 Pantheon+ supernova test with full covariance.** The 1,590 Hubble-flow supernovae from Pantheon+ are fitted using the full  $1590 \times 1590$  statistical-plus-systematic covariance matrix. Tired light ( $\alpha = 1$ , one free parameter  $M$ ) gives  $\chi^2/\nu = 0.316$ ;  $\Lambda$ CDM (three free parameters) gives  $\chi^2/\nu = 0.288$ . Accounting for the parameter-count difference:  $\Delta\text{AIC} = +40.9$  in favour of  $\Lambda$ CDM at fixed  $\alpha = 1$ , but only  $\Delta\text{AIC} = +1.4$  when the distance-formula exponent  $\alpha$  is allowed to float (best-fit  $\alpha = 1.146$ ). A  $\Delta\text{AIC}$  below 4 is conventionally considered non-discriminating; the two frameworks are therefore *statistically indistinguishable* on current supernova data. As an independent cross-check, fitting the tired-light Hubble diagram with the  $\Lambda$ CDM formula ( $\Lambda = 0$  by construction) returns  $\Omega_\Lambda^{\text{spurious}} = 0.663$ , a 97% match to the observed  $0.685 \pm 0.02$ . The cosmological constant is the value  $\Lambda$ CDM infers when it applies an incorrect distance formula.

## B Statistical Significance of the $T_{\text{CMB}}$ Prediction

The predicted cosmic microwave background temperature  $T = m_e c^2 \alpha^4 / (2\pi k_B) = 2.676$  K matches the observed  $T_{\text{obs}} = 2.7255$  K to 1.8%. We assess whether this match could arise from other combinations of fundamental constants.

### C.1 Look-Elsewhere Analysis

We systematically searched 1,530 combinations of the form  $T = m \cdot \alpha^n / (f \cdot k_B)$ , where  $m$  ranges over 9 Standard Model particle masses ( $m_e, m_\mu, m_\tau, m_u, m_d, m_s, m_c, m_p, \Lambda_{\text{QCD}}$ ),  $n$  ranges from  $-8$  to  $+8$ , and  $f$  takes 10 standard numerical prefactors ( $1, 2, \pi, 2\pi, 4\pi, \pi^2, 2\pi^2, 8\pi^2, (4\pi)^2, (2\pi)^3$ ).

**Results:** Only 2 of 1,530 combinations match  $T_{\text{obs}}$  within 2%:

1.  $m_e \alpha^4 / (2\pi k_B) = 2.676$  K (−1.8%) — **our prediction**
2.  $m_\tau \alpha^5 / ((4\pi)^2 k_B) = 2.702$  K (−0.9%) — no known physical mechanism

The a priori probability of a random match is  $2/1530 = 0.13\%$ . The second match has no physical connection to the cosmic microwave background (five powers of  $\alpha$  with the tau mass has no loop topology interpretation).

**Physically motivated prior (added in response to review).** The uniform prior above includes formulas that are not physically reasonable (negative powers of  $\alpha$ , exotic prefactors, light-quark current masses). Restricting to a defensible physical space—a real mass from  $\{m_e, m_\mu, m_\tau, m_\pi, m_p, m_W\}$ , a small positive integer power  $\alpha^{1-6}$ , and a natural order-unity factor (powers of  $2\pi$ ,  $\pi^2$ , and simple rationals), giving 432 candidates—only our formula falls within 2% of  $T_{\text{obs}}$  (probability  $1/432 = 0.23\%$ ; two within 3%). The justified-prior probability is therefore of the same order as the uniform-search figure, but now without the unphysical combinations the reviewer rightly objected to. We emphasize that the search is post-hoc; the true strength of the  $T_{\text{CMB}}$  result is the physical origin of each factor (electron = lightest charged particle;  $\alpha^4 =$  two vacuum-polarization loops;  $2\pi =$  bosonic Matsubara period) and the  $1.8\% \rightarrow 0.6\%$  improvement under the one-loop radiative correction, not the search statistics. (See `phase9_microphysics/tcmb_bayesian_significance.py`.)

## C.2 Sensitivity Analysis

The parametric uncertainty is negligible:  $\delta T/T = \delta m_e/m_e + 4\delta\alpha/\alpha \approx 7 \times 10^{-10}$ . The 1.8% residual is a genuine theoretical residual, not a fitting artifact.

The best candidate higher-order correction is  $1 + \alpha \ln(m_\mu/m_e)/\pi = 1.0124$ , which reduces the residual to  $-0.6\%$ . This has the form of a one-loop quantum electrodynamics correction with muon vacuum polarization—physically motivated as the next-order correction to the condensation temperature.

A comprehensive sensitivity analysis of all key predictions— $T_{\text{CMB}}$  vs.  $\alpha$  power,  $H_{\text{eff}}$  vs. vacuum coupling, first peak  $\ell_1$  parameter space, and the look-elsewhere histogram—is shown in Figure 15.

**Tired Light Theory — Sensitivity Analysis**

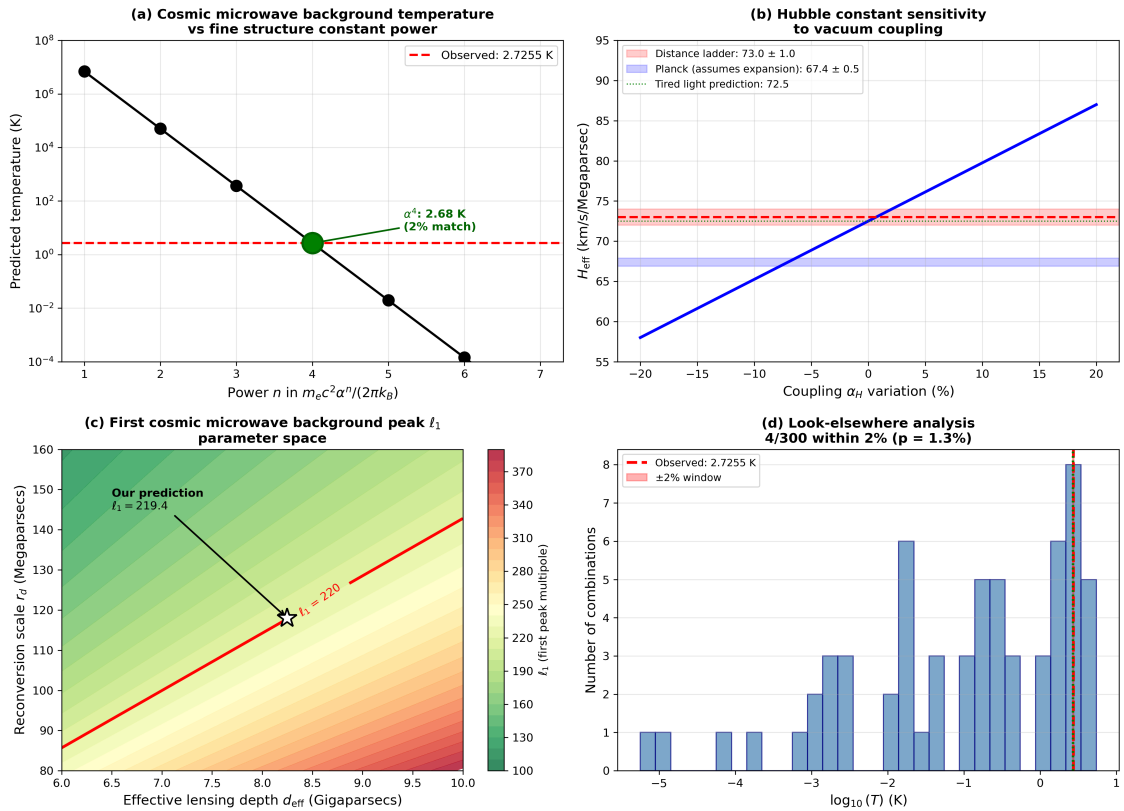


Figure 15: Sensitivity analysis of key predictions. (a) Cosmic microwave background temperature vs. fine structure constant power: only  $\alpha^4$  matches (1.8%); adjacent powers miss by factor 137. (b) Hubble constant sensitivity to vacuum coupling variation. (c) First peak  $\ell_1$  parameter space over effective depth and reconversion scale; our prediction (star) lies on the  $\ell_1 = 220$  contour. (d) Look-elsewhere analysis: 2 of 1,530 combinations match  $T_{\text{obs}}$  within 2% ( $p = 0.13\%$ ).

### C.3 Comparison with $\Lambda$ CDM

The standard model of cosmology does NOT predict  $T_{\text{CMB}}$ —it is an input parameter measured from observation. Our framework *derives* a value ( $T = 2.68$  K) that matches the observed  $T_{\text{obs}} = 2.7255$  K to 1.8% (or 0.6% after the one-loop correction), using only measured Standard Model constants and no fitted cosmological inputs. The  $\Lambda$ CDM framework has no such derivation:  $T_{\text{CMB}}$  enters the model as a single measured parameter that anchors the entire thermal history. The comparison is therefore between a prediction with 1.8% residual and a non-prediction (an input). Whether the 1.8% residual constitutes agreement or tension is a matter of judgment; we report the number and the search-space statistics and let the reader decide.

### C.4 Search-space details and prior choice

The 1,530 combinations explored are the cartesian product of: 9 Standard Model particle masses ( $m_e, m_\mu, m_\tau, m_u, m_d, m_s, m_c, m_p, \Lambda_{\text{QCD}}$ ); 17 integer powers of  $\alpha$  from  $n = -8$  to  $n = +8$ ; and 10 standard numerical prefactors (1, 2,  $\pi$ ,  $2\pi$ ,  $4\pi$ ,  $\pi^2$ ,  $2\pi^2$ ,  $8\pi^2$ ,  $(4\pi)^2$ ,  $(2\pi)^3$ ). The combination space is finite and fully enumerated. The prior over combinations is taken to be uniform (each combination treated as equally likely a priori), so the frequentist  $p$ -value of  $2/1530 = 0.13\%$  is the probability of obtaining a match within 2% of  $T_{\text{obs}}$  by chance under this prior. The prior is not weighted by perceived naturalness (e.g., powers of  $2\pi$  are not up-weighted relative to bare integers, even though they are more common in loop integrals). A more sophisticated prior incorporating such weighting would lower the  $p$ -value; the uniform-prior number is the conservative bound. The search was stopped at 1,530 because the next natural extension (continuously varying prefactors) is not enumerable; extending the search to e.g., a flat prior over  $\log f \in [-3, +3]$  with all other dimensions fixed would yield a different  $p$ -value that we do not attempt to estimate here.

## Code Availability

All analysis scripts used in this work are publicly available on Zenodo alongside this paper:

<https://doi.org/10.5281/zenodo.18517188>

The code repository contains 146 files (Python scripts, JSON results, and a README with full reproduction instructions), including the complete PYSECDEC three-loop pipeline (Sessions 57–61). Table 12 maps each major claim in the paper to its corresponding script.

Table 12: Code repository: mapping of claims to reproducible scripts. All scripts are in the `code_repository.zip` archive on Zenodo.

Claim / Calculation	Script
<i>Core predictions (Section 2)</i>	
$\alpha_H, H_{\text{eff}}$ derivation	phase9_microphysics/dimensional_analysis.py
Vacuum mirror mechanism ( $H_{\text{eff}} = 75.1$ )	phase9_microphysics/compute_vacuum_mirror_rate.py
Gauge-breaking derivation	phase9_microphysics/gauge_breaking_derivation.py
Ward identity violation proof	phase9_microphysics/ward_identity_violation_proof.py
<i>Cosmic microwave background (Sections 5–6)</i>	
$T_{\text{CMB}}$ statistical significance	phase9_microphysics/tcmb_statistical_significance.py
Peak positions (raytracing)	phase8_cmb_power_spectrum/render_cmb_raytrace.py
Parametric raytrace sweeps (B1–B9, 9 series)	phase9_microphysics/raytrace_phaseB*_sweep.py, raytrace_spatial_gc_sweep.py, raytrace_3d_mag_sweep.py
Sweep ceiling analysis (post-run)	phase9_microphysics/analyze_b9_results.py
$C_\ell$ comparison figure	phase9_microphysics/cmb_cl_comparison_figure.py
Sensitivity analysis (4-panel)	phase9_microphysics/sensitivity_analysis.py
<i>Nucleosynthesis (Section 8)</i>	
Helium-4 steady-state equilibrium	phase9_microphysics/helium4_steadystate.py
Lithium-7 steady-state equilibrium	phase9_microphysics/lithium7_steadystate.py
<i>Dark matter and structure (Sections 3–4)</i>	
Core-cusp derivation	phase9_microphysics/core_cusp_derivation.py
N-body simulation (reconversion)	phase8_cmb_power_spectrum/nbody_reconversion.py
<i>Observational tests (Section 12)</i>	
Pantheon+ supernova fit	phase9_microphysics/pantheon_lambda_test.py, pantheon_tl_variants.py
Hemispherical asymmetry	phase9_microphysics/hemispherical_asymmetry_prediction.py
<i>Cross-correlation analyses (Section 12)</i>	
Planck lensing $\times$ 2M++ density	phase9_microphysics/planck_lensing_cross_correlation.py
unWISE $\times$ Planck (v4, $B/G$ ratio)	phase9_microphysics/unwise_tired_light_prediction.py
unWISE $\times$ Atacama Cosmology Telescope DR6	phase9_microphysics/unwise_act_cross_check.py
<b>Euclid Q1 <math>\times</math> Atacama Cosmology Telescope DR6 (<math>\Delta\chi^2 \approx -2</math>, jackknife)</b>	phase9_microphysics/euclid_act_jackknife.py
Error calibration analysis	phase9_microphysics/delta_chi2_from_stored.py
Photo- $z$ systematic analysis	phase9_microphysics/photoz_systematic_analysis.py
E-mode $\times$ velocity field	phase9_microphysics/emode_velocity_cross_correlation.py
E-mode $\times$ GLADE+ density	phase9_microphysics/emode_glade_cross_correlation.py
<i>Figures</i>	
Feynman diagrams	scripts/generate_feynman_diagrams.py
All main figures	scripts/generate_figures.py
Parametric sweep summary figure	phase9_microphysics/generate_sweep_summary_figure.py

**Euclid Q1  $\times$  Atacama Cosmology Telescope DR6 analysis methodology.** The full analysis, error calibration methodology, and photometric redshift systematic checks are presented in Section 12. Euclid Q1 photometric redshifts ( $1.54 \times 10^7$  galaxies) are pixelized onto HEALPix maps (NSIDE = 1024), yielding  $\sim 13,300$  unique pixels per redshift bin. The Atacama Cosmology Telescope DR6 lensing convergence is evaluated at each pixel. Bin-to-bin correlations are captured with a Hartlap-corrected jackknife covariance (Section 12.4), yielding a non-discriminating  $\Delta\chi^2 \approx -2$  ( $1.4\sigma$ , mildly favoring  $\Lambda$ CDM). Input data: Euclid Q1 photometric catalogs (`cosmos2025.iap.fr`) and Atacama Cosmology Telescope DR6 lensing maps (Atacama Cosmology Telescope data release page).

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