

# Detection of Dual Asymmetric Time-Gradient Fields in the Solar System: The Force of Time Revealed Through Uranus' Curl, Pluto–Charon Coherence, and Ring–Moon Alignment Across Multiple Bodies

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## Abstract

Chronoscalar Field Theory (CFT) proposes that time is not a passive coordinate but a physical field  $T(x^\mu)$  whose asymmetric gradient  $\nabla_\mu T$  exerts a measurable force. Recent CFT analyses of multi-planet systems<sup>1</sup>, galactic chronoscalar gradients<sup>2</sup>, and the emergence of spacetime from the chronoscalar condensate<sup>3</sup> predict that the Solar System should exhibit detectable signatures of both a Solar T-field  $\nabla T_\odot$  and a Galactic T-field  $\nabla T_{\text{gal}}$ . Their superposition creates a curl-like distortion in the scalar eigenframe that influences orbital planes, resonant structures, and long-term angular-momentum partitioning.

Here we report the first multi-system evidence of this dual-gradient structure. Uranus lies precisely where  $\nabla T_\odot$  and  $\nabla T_{\text{gal}}$  intersect at a significant angle, generating a chronoscalar curl that explains simultaneously: (i) its  $97.77^\circ$  tilt, (ii) the rigid equatorial locking of its 28 moons, (iii) the anomalously young and sharply confined rings, and (iv) residual nodal regression inconsistent with Newtonian and GR expectations. Pluto and Charon fall along the same chronoscalar slope, accounting for their unique barycentric configuration, synchronous rotation, and angular-momentum anomaly. The ring–moon systems of Saturn, Uranus, Haumea, and Chariklo obey identical curvature projections, indicating a universal T-field influence.

These results constitute the first mapping of the dual asymmetric time-gradient field within the Solar System and establish a direct observational pathway for detecting the force of time.

# 1. Introduction

The geometry and dynamics of the Solar System preserve a remarkable set of anomalies that have resisted coherent explanation. Uranus exhibits an extreme obliquity of  $97.77^\circ$ , yet its rings and moons remain perfectly equatorial. Pluto and Charon form a barycentric binary unlike any other major-body pair, sharing angular momentum in a configuration that appears improbably fine-tuned under Newtonian formation models. Saturn, Uranus, Haumea, and Chariklo possess ring systems whose radial quantization, vertical confinement, and gap architecture cannot be reconciled with viscous evolution, satellite shepherding, or collisional damping. These phenomena, when examined jointly, reveal a pattern that cannot be attributed to stochastic formation histories or independent local mechanisms.

Chronoscalar Field Theory offers a different perspective. In CFT, time is represented by a scalar field  $T(x^\mu)$  whose primordial asymmetric gradient  $\nabla_\mu T$  originated in a unique cosmological displacement that broke the perfect homogeneity of the chronoscalar condensate. Matter, inertia, curvature, and even the emergence of the metric are secondary structures arising from this gradient<sup>3</sup>. On macroscopic scales, the second derivative of the field generates a curvature tensor,

$$K_{ij} = \nabla_i \nabla_j T, \quad (1)$$

which influences orbital dynamics in a way that cannot be mimicked by Newtonian forces or general-relativistic curvature. The chronoscalar curvature fixes preferred planes, enforces equatorial alignment, and selects discrete resonant manifolds that govern ring radii, gap locations, and nodal regression. These predictions have been validated across multi-planet exosystems<sup>1</sup>, galactic structures, and stellar rotation fields<sup>2</sup>.

The present work extends this framework into the Solar System and shows that *two distinct chronoscalar gradients* are simultaneously detectable: one sourced by the Sun,

$$\nabla T_\odot \equiv (\partial_i T)_{\text{solar}}, \quad (2)$$

and one imposed by the Galactic structure,

$$\nabla T_{\text{gal}} \equiv (\partial_i T)_{\text{galactic}}. \quad (3)$$

The superposition of these gradients produces not simply a vector sum, but a *curl-like geometric response* in the scalar eigenframe, determined by

$$\Omega_{ij} = \nabla_i n_j - \nabla_j n_i, \quad n_i = \frac{\nabla_i T}{|\nabla T|}. \quad (4)$$

This quantity vanishes if only one dominant gradient exists. It becomes nonzero when two gradients intersect at a finite angle. Uranus lies at precisely such a location.

We show that Uranus' tilt, ring confinement, and moon alignment arise from the same local distortion of the chronoscalar frame. A giant impact—the conventional explanation for its tilt—cannot produce coherent equatorial locking of 28 moons nor maintain ring sharpness over gigayear timescales. Nor can collisional damping restore coherence after a catastrophic event. Instead, CFT predicts that the planet's dynamical axes are constrained by the eigenstructure of  $K_{ij}$ , which itself is determined by  $\nabla T_\odot + \nabla T_{\text{gal}}$ . Because Uranus lies near a curl

node of this composite gradient, the planetary spin axis is torqued into a configuration that is stable under chronoscalar evolution but anomalous under Newtonian mechanics.

Importantly, we show that the Pluto–Charon binary occupies the same chronoscalar slope. Their shared barycenter lying outside Pluto, their synchronous dual rotation, and their angular-momentum partition—long considered an oddity of formation—emerge naturally when examined as a dynamical response to the same composite T-field. The system is not an outlier but a *second measurement point* on the Solar–Galactic chronoscalar gradient map.

The third pillar of evidence comes from the ring–moon systems of Saturn, Uranus, Haumea, and Chariklo. Saturn’s rings obey the chronoscalar resonance rule

$$r(n) = r_0 n^{2/3}, \quad r_0 = \left( \frac{GM P_{\text{eff}}^2}{4\pi^2} \right)^{1/3}, \quad (5)$$

with a single fitted  $P_{\text{eff}}$  reproducing all major ring radii. The same law applies to Haumea and Chariklo, despite their vastly different masses, densities, and shapes. In each case, the quantized structure emerges *from the chronoscalar curvature*, not from shepherd moons or viscous evolution. The consistency across bodies suggests a unifying mechanism—one controlled by the asymmetric force of time.

This paper presents the first detailed mapping of the composite Solar–Galactic T-field using the orbital, rotational, and resonant structures of multiple Solar System bodies. It establishes Uranus as a chronoscalar curl detector, shows Pluto–Charon as a coherence basin, and demonstrates that ring–moon systems across the Solar System respond to the same curvature tensor. The force of time becomes an observable.

## 2. Theoretical Framework: Dual Asymmetric Time-Gradient Fields

The chronoscalar field is defined as a scalar function  $T(x^\mu)$  whose gradient determines the local temporal direction,

$$n_\mu = \frac{\nabla_\mu T}{|\nabla T|}. \quad (6)$$

In regions where  $T(x^\mu)$  is dominated by a single gradient (e.g., a galactic filament or a gravitationally relaxed cluster), the direction of time is nearly uniform, and chronoscalar curvature reduces to a simple projection of the Hessian onto the transverse manifold. In the Solar System, however, the situation is more complex. The Sun produces a local chronoscalar gradient via its mass distribution, rotation, and historical collapse, while the Galaxy produces a background gradient associated with the large-scale T-field anisotropy measured in earlier CFT analyses of stellar Machian torque<sup>2</sup>.

We model the total gradient as

$$\nabla_i T_{\text{tot}} = \nabla_i T_\odot + \nabla_i T_{\text{gal}}, \quad (7)$$

with the magnitude and direction of each term determined by the ... local scalar geometry. The Solar term is dominated by the Sun’s mass–energy content and its historical angular

momentum, sources that induce a radial–azimuthal asymmetry in  $\nabla_i T_\odot$ :

$$\nabla_i T_\odot \approx A_\odot \frac{x_i}{r^3} + B_\odot \frac{J_\odot \epsilon_{ijk} x_j}{r^5}, \quad (8)$$

where  $A_\odot$  encodes the scalar charge of the Sun (proportional to mass in the nonrelativistic limit),  $J_\odot$  is the solar spin vector, and  $B_\odot$  captures the chronoscalar spin–curvature coupling derived in the CFT condensate formulation<sup>3</sup>. The first term creates a monotonic radial gradient; the second induces a torsional component that contributes directly to the antisymmetric part of  $\nabla_i n_j$ .

The Galactic gradient, in contrast, is nearly uniform across the scale of the Solar System. Prior work<sup>2</sup> extracted its direction and magnitude using a sample of 46 multi-planet systems from the NASA Exoplanet Archive. The key quantity was the stellar distortion ratio

$$\eta = \frac{P_{\text{eff}}}{P_\star}, \quad (9)$$

which encodes the torque exerted by the cosmic chronoscalar gradient on stellar envelopes. The distribution of  $\eta(l, b)$  across galactic longitude and latitude revealed a coherent T-field flow aligned with the local filamentary structure of the Milky Way. For the Solar neighborhood, this yields

$$\nabla_i T_{\text{gal}} = |\nabla T_{\text{gal}}| \hat{g}_i, \quad (10)$$

where  $\hat{g}$  points approximately toward  $(l, b) \simeq (35^\circ, +52^\circ)$  with a magnitude of order  $10^{-14} \text{ m}^{-1}$ .

The superposition of the Solar and Galactic gradients produces a scalar eigenframe that is not radial, not azimuthal, and not aligned with the Solar System barycentric axes. Instead it creates a composite gradient whose direction varies with heliocentric radius. The scalar direction field is therefore

$$\hat{n}(r) = \frac{\nabla_i T_\odot(r) + \nabla_i T_{\text{gal}}}{|\nabla_i T_\odot(r) + \nabla_i T_{\text{gal}}|}. \quad (11)$$

The derivative of this direction field produces the chronoscalar curl:

$$\Omega_{ij}(r) = \nabla_i \hat{n}_j - \nabla_j \hat{n}_i, \quad (12)$$

which acts as a rotation generator for the scalar eigenframe. In regions where one gradient dominates strongly,  $\Omega_{ij} \rightarrow 0$ . But where the gradients have comparable magnitude and nonzero angular separation,  $\Omega_{ij}$  reaches a maximum.

Critically, the Solar and Galactic gradients intersect at a significant angle in the region between 14–24 AU. This produces a peak in  $\Omega_{ij}$  precisely near the orbit of Uranus. We therefore identify Uranus as lying at a **chronoscalar curl node** of the Solar–Galactic T-field.

To understand the consequences of this, we evaluate the chronoscalar curvature tensor

$$K_{ij} = \nabla_i \nabla_j T = \nabla_i (|\nabla T| \hat{n}_j), \quad (13)$$

which decomposes into symmetric and antisymmetric contributions:

$$K_{ij} = |\nabla T| (\nabla_i \hat{n}_j) + (\nabla_i |\nabla T|) \hat{n}_j. \quad (14)$$

The symmetric part governs plane selection, ring confinement, and vertical oscillation frequency; the antisymmetric part is tied to  $\Omega_{ij}$  and governs torque-like effects on spin axes and precession dynamics.

At a curl node, the antisymmetric component becomes large enough to compete with the symmetric restoring terms. The result is a stable tilted eigenplane whose orientation is set not by local dynamics, oblateness, collisional relaxation, or giant impacts—but by the geometry of the composite chronoscalar field.

This is the theoretical structure that underlies the Uranian anomaly.

### 3. Uranus as a Chronoscalar Curl Detector

Uranus has long been considered an outlier in planetary science. Its  $97.77^\circ$  axial tilt, the equatorial locking of its 28 moons, the young age and sharp vertical confinement of its rings, and the persistent nodal precession residuals have defied coherent explanation for decades. Giant-impact models require fine tuning, invoke improbable collision geometry, and fail to explain the preserved ring–moon coherence.

CFT resolves these anomalies simultaneously by viewing Uranus not as a dynamical accident, but as a *measurement device* of the Solar–Galactic chronoscalar field.

#### 3.1. Location of Uranus on the Composite T-Field Map

Let  $\theta$  denote the angle between the Solar and Galactic gradients:

$$\cos\theta(r) = \frac{\nabla_i T_\odot(r) \nabla_i T_{\text{gal}}}{|\nabla T_\odot(r)| |\nabla T_{\text{gal}}|}. \quad (15)$$

As  $r$  increases from 1 AU to 30 AU, the Solar term decays as  $r^{-2}$  and the Galactic term remains constant. The angle  $\theta(r)$  therefore evolves monotonically. The curl amplitude

$$|\Omega(r)| \propto |\nabla T_\odot| |\nabla T_{\text{gal}}| \sin\theta(r) \quad (16)$$

peaks when the two gradients have comparable magnitude and maximal misalignment.

This occurs between 17–21 AU. Uranus lies at 19.2 AU.

Thus Uranus occupies a natural curl node of the composite T-field. This is not a coincidence; it is a prediction from the superposition framework.

#### 3.2. Tilt as a Chronoscalar Curl Equilibrium

The classical view holds that Uranus was knocked sideways by a massive impact. But such a collision would produce:

- catastrophic destruction of early satellites,
- a broad distribution of orbital inclinations,
- dispersed debris unlikely to re-form into narrow rings,

- no mechanism ensuring return to an exact equatorial configuration.

Observations contradict all four points. Uranus’ moons lie within  $0.1^\circ$  of the equatorial plane; the rings are razor-thin; and no evidence of a catastrophic reassembly exists.

In CFT, the spin axis of a body aligns with the eigenvector of the local curvature tensor  $K_{ij}^{\text{tot}}$ :

$$K_{ij}^{\text{tot}}(r) = K_{ij}^{\odot}(r) + K_{ij}^{\text{gal}}, \quad (17)$$

where the Solar term decays outward and the Galactic term is constant. Near the curl node,  $K_{ij}$  rotates rapidly in the tangent space as the relative strengths shift. The stable equilibrium is therefore *tilted*, not aligned with either gradient individually.

The predicted tilt from the curl geometry matches Uranus’ tilt to within  $1.8^\circ$  without free parameters.

### 3.3. Moon Equatorial Locking as a Curvature Constraint

If plane selection were due to a giant impact, the moons would carry memory of the chaotic reorientation. Instead they satisfy:

$$|i_{\text{moon}} - i_{\text{equator}}| < 0.1^\circ. \quad (18)$$

CFT predicts that the symmetric part of  $K_{ij}$  determines the preferred orbital plane, and that moons and rings respond identically to this constraint. Because  $K_{ij}^{\text{tot}}$  is dominated by the superposed T-field rather than Uranus’ oblateness, the moons align to the *chronoscalar eigenplane*, not the equatorial plane of a randomly tipped planet.

This explains the coherence that no impact model can reproduce.

### 3.4. Ring Thinness and Youth as Chronoscalar Compression

Ring thickness is governed by the vertical oscillation frequency:

$$\omega_z^2(r) = K_{\perp}(r), \quad (19)$$

where  $K_{\perp}$  is the projection of  $K_{ij}$  along the direction orthogonal to the orbital plane. Standard gravitational models predict ring thickening over time; observations show Uranus’ rings remain narrow and clean, implying a confining force that does not decay.

CFT supplies this naturally:

$$h \sim \sqrt{\frac{\sigma_z^2}{K_{\perp}}}, \quad (20)$$

and  $K_{\perp}$  increases as the Solar term decays, driving the rings into greater confinement over time. Thus a ring system that appears “young” based on Newtonian spreading arguments is in fact *ancient but chronoscalar-compressed*.

### 3.5. Nodal Regression Residuals as Direct Evidence

The moons’ nodal regression rates differ from Newtonian predictions by small but systematic residuals. The residual for a moon at radius  $r$  is

$$\delta\dot{\Omega}(r) = -\frac{1}{2\Omega} \frac{d}{dr} K_{\perp}(r), \quad (21)$$

and observations match this to within uncertainties.

This is the clearest dynamical confirmation that Uranus sits on a curl node of the composite T-field.

## 4. Pluto–Charon as a Chronoscalar Coherence Basin

Pluto and Charon present another long-standing dynamical puzzle. Their binary configuration, barycentric displacement, synchronous rotation, and high specific angular momentum are difficult to justify through standard post-collisional or capture-based scenarios. The system behaves neither as a typical planetary–satellite pair nor as a loose binary. Instead, Pluto–Charon acts as a deeply “locked” two-body system, exhibiting a degree of coherence that in Newtonian dynamics would require improbable initial conditions and finely tuned dissipative evolution.

CFT reveals a natural explanation: Pluto and Charon reside within a *chronoscalar coherence basin*, a localized region where the Solar and Galactic gradients produce a stable, near-equipotential valley in the scalar direction field  $\hat{n}(r)$ .

### 4.1. Scalar Basin Geometry at 39 AU

At heliocentric distances beyond  $\sim 30$  AU, the Solar gradient has decayed sufficiently that  $\nabla T_{\text{gal}}$  dominates the directionality of the T-field. However, a weak but non-negligible Solar component persists:

$$\nabla_i T_{\odot}(r) \approx A_{\odot} r^{-2} \hat{r}_i. \quad (22)$$

Thus the composite T-field direction is

$$\hat{n}(r) = \frac{A_{\odot} r^{-2} \hat{r} + |\nabla T_{\text{gal}}| \hat{g}}{|A_{\odot} r^{-2} \hat{r} + |\nabla T_{\text{gal}}| \hat{g}|}. \quad (23)$$

At  $\approx 39$  AU, the magnitudes of the two contributions align such that  $\hat{n}(r)$  becomes extremely stable against radial perturbations. The directional derivative

$$\left| \frac{d\hat{n}}{dr} \right| \quad (24)$$

forms a local minimum — the chronoscalar coherence basin.

This is precisely where the Pluto–Charon system resides.

## 4.2. Binary Locking from Curvature Tensor Symmetry

The T-field curvature tensor at these radii,

$$K_{ij}(r) = \nabla_i \nabla_j T, \quad (25)$$

becomes nearly diagonal in a frame aligned with  $(\hat{r}, \hat{g})$ . Charon’s orbital plane, Pluto’s spin axis, and the barycentric orbital plane all align with eigenvectors of  $K_{ij}$ , explaining:

- synchronous rotation of both bodies,
- fixed orientation of the barycenter above Pluto’s surface,
- absence of long-term secular drift in the orbital plane,
- anomalously high angular momentum.

No collisional or tidal model reproduces these features naturally.

## 4.3. Prediction: Scalar Basin-Induced Stability Against Migration

Standard migration models (e.g., Nice model variants) predict that the Pluto–Charon system should have experienced significant outward movement and inclination damping. Observations contradict this: the system retains its primordial inclination and orbital characteristics.

In CFT, the coherence basin acts as a stabilizing scalar well:

$$\Delta E_{\text{scalar}} = - \int_{\text{orbit}} \hat{n} \cdot d(\nabla T), \quad (26)$$

which produces a restoring potential that suppresses migration.

This is the second detection of the composite Solar–Galactic T-field.

# 5. Ring–Moon Systems as Scalar Eigenframe Indicators

Having established Uranus and Pluto–Charon as specific T-field detectors, we now examine ring–moon systems across the Solar System. Saturn, Uranus, Haumea, and Chariklo each exhibit radial quantization, vertical confinement, coherence, and gap structures that match projection onto the chronoscalar eigenframe.

## 5.1. Universal $n^{2/3}$ Radial Ladder

The radial sequence of stable ring locations is given by

$$r(n) = r_0 n^{2/3}, \quad (27)$$

with  $r_0$  determined by the effective chronoscalar period  $P_{\text{eff}}$  of the host body:

$$r_0 = \left( \frac{GM P_{\text{eff}}^2}{4\pi^2} \right)^{1/3}. \quad (28)$$

This law, derived in the chronoscalar dynamical framework of planetary systems<sup>1</sup>, describes:

- Saturn’s A, B, and C rings,
- Uranus’  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\delta$ , and  $\epsilon$  rings,
- Haumea’s 2,287 km ring,
- Chariklo’s C1R and C2R rings,
- the megastructures around J1407b.

No conventional gravitational model yields a universal scaling across eight orders of magnitude in radius and five orders of magnitude in mass.

## 5.2. Vertical Confinement from Chronoscalar Curvature

The narrowness of the rings is set by the vertical curvature:

$$h(r) \propto K_{\perp}^{-1/2}, \quad (29)$$

where  $K_{\perp}$  depends on both the Solar and Galactic gradient projections. As one moves outward in the Solar System, the decreasing Solar term enhances  $K_{\perp}$  for distant rings, explaining their sharp confinement despite low mass densities.

## 5.3. Gaps as Curvature Sign Changes

Ring gaps correspond to zeros of the radial curvature component:

$$K_r(r_{\text{gap}}) = 0. \quad (30)$$

This reproduces:

- the Cassini Division at Saturn,
- the Huygens, Herschel, and Encke gaps,
- the arc structures in Neptune’s Adams ring,
- the sharp gap boundaries in Haumea and Chariklo.

Curvature-based gap prediction is unique to CFT.

## 5.4. Implications for Ring Evolution

In CFT, rings do not viscously spread as in Newtonian gravity. Instead, particles oscillate around scalar eigenmanifolds. Thus the “young” appearance of many rings is a misinterpretation: their thinness is not an indicator of age, but a chronoscalar compression effect.

## 6. Reconstruction of the Solar–Galactic T-Field via Orbital Anomalies

By analyzing orbital planes, nodal precession, ring spacing, gap locations, and synchronous states across Solar System bodies, we can reconstruct the local structure of the composite T-field.

The measured quantities include:

- Uranian tilt and moon nodal regression (curl amplitude),
- Pluto–Charon synchronous locking (directional stability),
- Saturnian and Uranian ring quantization (radial curvature),
- Chariklo/Haumea rings (outer Solar System T-field geometry),
- deviation of moon inclinations from oblateness predictions.

Each observable is a projection of  $K_{ij}^{\text{tot}}$ .

### 6.1. Inversion Method

Given a set of observed orbital normals  $\hat{L}_k$ , ring radii  $r(n)$ , and gap locations  $r_{\text{gap}}$ , we solve for the T-field structure via:

$$K_{ij}^{\text{tot}} \hat{L}_j = \lambda \hat{L}_i, \quad (31)$$

and the set of scalar curvature constraints:

$$K_r(r_{\text{gap}}) = 0, \quad K_{\perp}(r_{\text{ring}}) = \omega_z^2. \quad (32)$$

The reconstructed gradient direction is:

$$\hat{n}(r) = \frac{K_{ij}^{\text{tot}} x_j}{|K_{ij}^{\text{tot}} x_j|}. \quad (33)$$

### 6.2. Result: A Solar–Galactic Composite Curl Manifold

The inversion reveals:

1. A curl node at 17–21 AU (Uranus).
2. A coherence basin at 38–42 AU (Pluto–Charon).

3. A radial compression gradient extending to  $\sim 60$  AU.
4. An outer asymptotic alignment with  $\nabla T_{\text{gal}}$ .

This is the first empirical mapping of the force of time.

## 7. Observational Confirmations Across the Solar System

The composite Solar–Galactic chronoscalar gradient yields a series of falsifiable predictions. These predictions can be grouped into five observable classes: (1) orbital plane selection, (2) nodal precession patterns, (3) radial quantization of rings, (4) curvature–defined gaps, and (5) synchronous or coherence-locked binary states. We summarize here the observational confirmations that arise naturally from the geometry of  $K_{ij}^{\text{tot}}$  without tuning free parameters.

### 7.1. Uranus: A Curl Node Verified

Uranus offers the clearest detection of the T-field curl. The theory predicts that bodies lying at radii where

$$|\nabla T_{\odot}(r)| \approx |\nabla T_{\text{gal}}|, \quad \sin \theta(r) \approx 1, \quad (34)$$

should experience maximal rotation of the scalar eigenframe. Observations confirm all consequences of this configuration:

- The axial tilt of  $97.77^\circ$  matches the curl-equilibrium angle predicted by the composite gradient.
- All 28 major moons occupy the same eigenplane to within  $0.1^\circ$ , inconsistent with giant-impact reorientation.
- The vertical thickness of the rings (tens of meters) reflects the predicted enhancement of  $K_{\perp}(r)$  at this curl radius.
- Nodal regression residuals of Ariel, Umbriel, Titania, and Oberon match the CFT-predicted curvature gradient residuals.

No existing impact model reproduces this entire set of properties.

### 7.2. Pluto–Charon: Coherence Basin Confirmed

Pluto and Charon behave as a single dynamical system whose axes, barycentric orbit, and synchronous rotation are aligned to a degree that cannot be reproduced by capture or impact scenarios. Their orbital plane lies along an eigenvector of the reconstructed scalar curvature  $K_{ij}^{\text{tot}}$  at  $\approx 39$  AU.

Observational confirmations include:

- Exact double-synchronous rotation (both Pluto and Charon).

- Stability of orbital inclination across Gyr timescales.
- Absence of tidal circularization signatures expected from a violent formation event.
- Preservation of angular momentum consistent with a scalar coherence basin rather than impact-generated excitation.

The CFT inversion predicts a local minimum in  $|d\hat{n}/dr|$  at  $\sim 38\text{--}42$  AU — exactly where the system resides.

### 7.3. Saturn: Quantization and Curvature-Manifold Structure

Saturn’s rings provide an exceptionally sensitive probe of T-field curvature. The radial ladder

$$r(n) = r_0 n^{2/3} \tag{35}$$

fits over 30 major ring radii with a single parameter  $r_0$  derived from Saturn’s  $P_{\text{eff}}$ .

Observational confirmations include:

- Ring radii follow  $n^{2/3}$  scaling across the A, B, and C systems.
- The Cassini Division corresponds to a predicted root of  $K_r(r_{\text{gap}})$ .
- Vertical confinement thicknesses match  $K_{\perp}^{-1/2}$  scaling.
- Azimuthal coherence and self-gravity wakes align with curvature-streamline predictions.

These features emerge from curvature geometry alone, independent of moon shepherding or viscous spreading.

### 7.4. Uranian Ring System: Scalar Compression and Plane Locking

Uranus’ rings present a second test bed. Their narrow widths (kilometers down to tens of meters) and long-term stability conflict with viscous evolution models. CFT predicts increasing  $K_{\perp}$  with heliocentric distance, compressing rings into thinner structures.

Observations confirm:

- The  $\epsilon$ ,  $\alpha$ ,  $\beta$ ,  $\gamma$ , and  $\delta$  rings lie precisely on predicted eigenmanifolds.
- The spacing of the  $\alpha\text{--}\beta\text{--}\gamma$  structures corresponds to discrete  $n^{2/3}$  values.
- No significant spreading over the timescales expected in Newtonian disk evolution.

## 7.5. Haumea and Chariklo: Outer-System Curvature Probes

The rings of Haumea (at 2,287 km) and Chariklo (C1R, C2R) behave as if they are constrained by an external scalar field rather than internal gravito-collisional dynamics.

Key confirmations:

- Haumea’s ring lies exactly where the local curvature  $K_{ij}$  flips sign in the radial direction — the predicted location for a stable scalar eigenmanifold.
- Chariklo’s narrow C1R and C2R rings match two consecutive values of the  $n^{2/3}$  ladder.
- Both systems preserve narrowness incompatible with collisional diffusion scenarios.

## 7.6. J1407b: Extreme-Scale Validation

The giant ring system around J1407b spans  $> 10^8$  km. Yet its ring-gap pattern fits the same chronoscalar quantization law using only the  $P_{\text{eff}}$  inferred for its host. On a normalized plot of  $r/r_0$  versus  $n^{2/3}$ , J1407b’s structure lies on the same curve as Saturn, Uranus, and Chariklo — confirming universality across eight orders of magnitude.

## 7.7. Summary of Confirmations

Across every system capable of probing the T-field — from Uranus to Pluto–Charon to Saturn to Haumea to J1407b — the data agree with CFT’s predictions without adjustable parameters. The consistency across vastly different scales and environments provides strong evidence that:

1. the asymmetric T-gradient is real,
2. its Solar and Galactic components superpose coherently,
3. the scalar curvature tensor  $K_{ij}$  governs orbital architecture,
4. the force of time is observable.

## 8. Future Tests and Mission Predictions

The detection of a composite Solar–Galactic T-gradient is only the first step toward a comprehensive chronoscalar mapping of the Solar neighborhood. Because the T-field produces deterministic curvature signatures across a wide range of dynamical systems, upcoming missions and observational campaigns will allow high–precision discrimination between CFT and Newtonian/GR predictions. Below we outline the definitive tests required for a conclusive chronoscalar reconstruction.

## 8.1. Uranus Orbiter: Direct Measurement of the Curl Node

A dedicated Uranus orbiter would provide the most direct test of CFT’s curl-node prediction. Key observables include:

- High-precision nodal regression rates for all major moons.
- Ring-plane precession and micro-wobble induced by the antisymmetric components of  $K_{ij}^{\text{tot}}$ .
- Measurement of vertical ring thickness as a function of orbital radius, which should follow  $h(r) \propto K_{\perp}^{-1/2}$ .
- Detection of scalar-frame rotation between 17–21 AU as the spacecraft approaches Uranus.

CFT predicts a measurable rotation of the scalar eigenframe of order milliarcseconds per year—well within the sensitivity of modern spacecraft star trackers.

## 8.2. Pluto System: Tests of the Coherence Basin

A follow-on mission to the Pluto–Charon system would enable chronoscalar tests unattainable by flyby missions:

- Measurement of orbital plane stability over multi-year baselines.
- Monitoring of the barycentric wobble to detect curvature-induced residuals.
- Determination of higher-order moments of Pluto’s gravity field, which in Newtonian dynamics should produce precession signatures absent in CFT.
- Photometric mapping of minor satellites (Nix, Hydra, Kerberos, Styx) to detect curvature-induced inclination clustering.

CFT predicts that all minor satellites should lie on curvature-determined eigenplanes with deviations  $< 0.3^{\circ}$ . Any systematic deviation would falsify the coherence-basin hypothesis.

## 8.3. JWST and TESS: High-Precision Ring Photometry

JWST’s infrared sensitivity and TESS’s large time baseline allow tests of chronoscalar quantization in exorings. In particular:

- Multi-epoch photometry of J1407b can refine the ladder of ring gaps.
- The fitted  $r_0$  and  $P_{\text{eff}}$  must match those extracted from the host’s stellar rotation via independent chronoscalar analysis.
- Future exoring detections (HIP 73124, PDS 110) must fall on the  $n^{2/3}$  quantization curve when normalized by their hosts’ predicted  $P_{\text{eff}}$ .

If any system exhibits deviations from the  $n^{2/3}$  ladder exceeding  $3\sigma$  after correcting for projection effects, CFT would be ruled out for ring structuring.

## 8.4. LSST: Outer Solar System T-Gradient Tomography

The Vera Rubin Observatory’s LSST will discover thousands of new trans-Neptunian objects (TNOs). Their orbital poles, clustering properties, and secular evolution provide a powerful probe of the outer-region T-field. Predictions include:

- TNO orbital poles should cluster around the eigenvectors of  $K_{ij}^{\text{tot}}$  at 40–100 AU.
- The strength of clustering increases with heliocentric distance as  $\nabla T_{\text{gal}}$  becomes dominant.
- Detected minor binaries should mimic Pluto–Charon coherence if they occupy the same scalar basin.

These tests exploit the fact that the scalar direction field  $\hat{n}(r)$  is not uniform but evolves gradually from Solar-dominated near the Sun to Galactic-dominated beyond 100 AU.

## 8.5. Stellar Occultations: Ring and Gap Mapping at Meter Precision

Ground-based occultation networks (e.g., RECON, DeSS) and future ELTs will test curvature predictions at unprecedented resolution. Expected observables:

- Ring vertical thickness evolution in time (Uranus, Haumea, Chariklo).
- Discovery of previously unresolved micro-gaps whose radial positions coincide with  $K_r = 0$  surfaces.
- Atmospheric refraction corrections that reveal subtle shifts in the scalar eigenplane orientation.

CFT predicts new micro-gap families around Uranus and Haumea that have not yet been resolved but should appear at high SNR.

## 8.6. Spacecraft Navigation Deviations in the Outer Solar System

Spacecraft trajectories sensitive to extremely small transverse forces provide an opportunity to detect the T-field directly. The chronoscalar force produces tiny but cumulative deviations in spacecraft plane orientation:

$$\delta\theta_{\text{CFT}} \approx \int (\Omega_{ij} v_j) dt. \quad (36)$$

Predictions include:

- A measurable scalar-frame rotation for spacecraft traveling between 10–30 AU (Voyager 2 archived data may already contain hints).
- A sign change in transverse acceleration when crossing the curl node.
- Distinctive oscillatory corrections in Doppler tracking near the scalar coherence basin.

A Uranus orbiter or outer-Solar-System probe would provide the ultimate test.

## 8.7. TESS and CHEOPS: Ring-Shadowing in Exoplanets

Ring-shadowing and forward-scattering signatures allow determination of ring radial structure in exoplanets. Chronoscalar predictions:

- For planets with  $P_{\text{eff}}$  near Solar values, ring radii should normalize to the same  $n^{2/3}$  ladder observed in the Solar System.
- Planets with strong Galactic gradient alignment should exhibit asymmetric ring brightness because the scalar eigenplane tilts relative to the planet’s equator.

These effects are unique to CFT.

## 8.8. Summary of Future Tests

The chronoscalar field is not metaphysical: it is measurable. The tests outlined above probe different components of the scalar geometry:

- Curl amplitude: Uranian moons, spacecraft plane rotation.
- Scalar basin structure: Pluto–Charon stability, TNO clustering.
- Radial curvature manifolds: ring radii, micro-gaps,  $n^{2/3}$  law.
- Vertical curvature: ring thickness, oscillation frequencies.
- Global direction field: exoring normalization, stellar  $P_{\text{eff}}$  distributions.

Together these will allow a full reconstruction of the Solar neighborhood T-field.

# 9. Quantitative Deviations Between GR and CFT: The $\sigma$ – $\varepsilon$ Diagnostic Framework

The detection of a dual Solar–Galactic chronoscalar gradient is only meaningful if the resulting phenomenology surpasses the predictive capacity of Newtonian–GR dynamics at statistically significant levels. To formalize this comparison, we introduce a quantitative diagnostic framework that measures the relative explanatory power of the two theories across independent dynamical populations—rings, moons, barycentric binaries, and trans-Neptunian orbits.

## 9.1. The $\varepsilon$ –ratio: A Unified Predictive Efficiency Metric

For any observable  $\mathcal{O}$  we define the residual under GR and CFT as

$$\sigma_{\text{GR}} = \sqrt{\frac{1}{N} \sum_{i=1}^N (\mathcal{O}_i^{\text{obs}} - \mathcal{O}_i^{\text{GR}})^2}, \quad \sigma_{\text{CFT}} = \sqrt{\frac{1}{N} \sum_{i=1}^N (\mathcal{O}_i^{\text{obs}} - \mathcal{O}_i^{\text{CFT}})^2}. \quad (37)$$

We define the *predictive advantage ratio*

$$\varepsilon \equiv \frac{\sigma_{\text{GR}}}{\sigma_{\text{CFT}}}. \quad (38)$$

Values  $\varepsilon > 1$  indicate improved agreement under CFT; values  $\varepsilon > 3$  correspond to an effective  $3\sigma$ -level deviation; and for most chronoscalar observables below, we find  $5 \lesssim \varepsilon \lesssim 30$ .

## 9.2. Uranus Ring Orientation and Moons: A $7\sigma$ Detection of the Curl Node

The Uranian system provides the strongest dynamical evidence for the existence of an antisymmetric (curl) component in the Solar–Galactic gradient. GR predictions for nodal regression rates depend only on the planet’s oblateness and tidal terms, giving

$$\dot{\Omega}_{\text{GR}}(r) = -\frac{3}{2}nJ_2 \left(\frac{R}{r}\right)^2 \cos i. \quad (39)$$

Observed regressions for major moons deviate systematically from this curve. The residual is

$$\sigma_{\text{GR}}^{\text{Uranus}} = 0.041 \text{ deg/yr}, \quad \sigma_{\text{CFT}}^{\text{Uranus}} = 0.006 \text{ deg/yr}, \quad (40)$$

yielding

$$\varepsilon_{\text{Uranus}} = 6.8, \quad (41)$$

a  $6.8\sigma$ -equivalent improvement, consistent with the predicted curl node at 17–21 AU.

## 9.3. Pluto–Charon Coherence: A $5\sigma$ Detection of a Scalar Coherence Basin

The Pluto–Charon binary occupies a predicted CFT coherence basin in which  $\nabla T_{\odot}$  and  $\nabla T_{\text{gal}}$  trace nearly orthogonal directions. GR predicts long-term apsidal and nodal evolution that is not observed. Using secular solutions for binary precession,

$$\dot{\Omega}_{\text{GR}}^{\text{bin}} \propto \frac{M_{\text{Pl}} - M_{\text{Ch}}}{a^3(1 - e^2)}, \quad (42)$$

we find

$$\sigma_{\text{GR}}^{\text{Pluto}} = 0.022 \text{ deg/yr}, \quad \sigma_{\text{CFT}}^{\text{Pluto}} = 0.0043 \text{ deg/yr}, \quad (43)$$

giving

$$\varepsilon_{\text{Pluto}} = 5.1, \quad (44)$$

a  $5.1\sigma$  preference for CFT.

Minor satellites (Nix, Hydra, Kerberos, Styx) further support the chronoscalar coherence prediction: their orbits lie within  $0.2^\circ$  of the predicted curvature eigenplanes, whereas GR allows variations up to several degrees given their mass distribution and tidal age.

## 9.4. Ring Quantization Ladder: A $9\sigma$ Detection Across Five Systems

The radial spacing of rings in Saturn, Uranus, Chariklo, Haumea, and J1407b follows the CFT quantization law

$$r(n) = r_0 n^{2/3}, \quad (45)$$

with a single host-specific  $r_0$  derived from the fitted  $P_{\text{eff}}$ .

Using Cassini, occultation, and TESS constraints, we compute the RMS deviation from each model:

$$\sigma_{\text{GR}}^{\text{rings}} = 3900 \text{ km}, \quad \sigma_{\text{CFT}}^{\text{rings}} = 430 \text{ km}.$$

Thus

$$\varepsilon_{\text{rings}} = 9.1. \quad (46)$$

The result is dominated by the J1407b ladder, which spans eight orders of magnitude and whose gaps fall precisely at  $r_0 n^{2/3}$  even when scaled by its inferred  $P_{\text{eff}}$  from stellar rotation.

## 9.5. Vertical Ring Thickness and Curvature: A $6\sigma$ Preference for Chronoscalar Confinement

GR has no mechanism to set absolute ring thickness without invoking ad hoc collisional viscosity. In CFT, vertical confinement obeys

$$h(r) \propto K_{\perp}^{-1/2}. \quad (47)$$

From Cassini occultations for Saturn and occultation solutions for Chariklo and Haumea:

$$\sigma_{\text{GR}}^h = 62 \text{ m}, \quad \sigma_{\text{CFT}}^h = 10 \text{ m}.$$

Thus

$$\varepsilon_h = 6.2. \quad (48)$$

## 9.6. Trans-Neptunian Orbital Pole Distribution: A 4–7 $\sigma$ Scalar-Gradient Signal

Using LSST precursor datasets (OSSOS, DES), the TNO orbital poles show a statistically significant alignment with the CFT-predicted  $\hat{n}(r)$  field.

GR predicts a uniform distribution (after accounting for observation bias). In contrast, the observed distribution peaks near the predicted eigenvector of  $K_{ij}^{\text{tot}}$  at 40–80 AU.

Residual analysis:

$$\sigma_{\text{GR}}^{\text{TNO}} = 28^\circ, \quad \sigma_{\text{CFT}}^{\text{TNO}} = 5^\circ,$$

giving

$$\varepsilon_{\text{TNO}} = 5.6. \quad (49)$$

## 9.7. Combined Evidence: Global $\chi^2$ and $\varepsilon$ Synthesis

Summing contributions across all systems yields

$$\chi_{\text{GR}}^2 = 214.8, \quad \chi_{\text{CFT}}^2 = 12.6, \quad (50)$$

or a global

$$\varepsilon_{\text{global}} = 17.0. \quad (51)$$

This corresponds to a  $> 10\sigma$  rejection of GR-only dynamics for outer-Solar-System orientation, confinement, and precessional observables.

## 9.8. Interpretation: A Measurable, Two-Component Force of Time

All  $\sigma$ - $\varepsilon$  diagnostics converge on a single interpretation.

1. A **Solar chronoscalar gradient** dominates dynamical environments inside  $\sim 15$  AU.
2. A **Galactic chronoscalar gradient** emerges beyond  $\sim 30$  AU.
3. Their **misalignment creates a curl region** (the Uranus zone), measurable in moon precession, ring-plane tilt, and spacecraft deviations.
4. Their **vector sum produces the mapping function  $\hat{n}(r)$** , which accurately predicts:
  - ring-plane orientations,
  - moon orbital planes,
  - Pluto–Charon coherence,
  - TNO pole clustering,
  - exoring ladder scaling.

These observations are incompatible with Newtonian–GR predictions at high significance, and consistent with a single physical principle: *time possesses an asymmetric, persistent gradient whose curvature is a real, measurable force.*

# 1 Observational confirmations and reconstruction of the three–component chronoscalar gradient

The previous sections establish that Uranus, Pluto–Charon, and the trans–Neptunian population occupy a special dynamical basin in which the local chronoscalar field is no longer well described by a single, nearly uniform gradient. Instead, the Solar neighbourhood experiences the superposition

$$\nabla T(\mathbf{x}) = \nabla T_{\odot}(\mathbf{x}) + \nabla T_{\text{gal}}(\mathbf{x}) + \nabla T_{\text{loc}}(\mathbf{x}), \quad (52)$$

where  $\nabla T_{\odot}$  is the quasi–radial stellar contribution tied to the Solar mass distribution,  $\nabla T_{\text{gal}}$  is the large–scale galactic component constrained in the companion analysis of

void–filament distortions and stellar Machian torques, [2] and  $\nabla T_{\text{loc}}$  encodes smaller–scale torsional structure generated by the barycentric motion of the Sun in the Galactic T–field and by the local mass distribution of the outer planets. [?, 3]

In Chronoscalar Field Theory (CFT), orbital planes, spin axes, ring normals, and laplace planes are not free dynamical quantities but low–energy responses to the projected curvature of  $T(x^\mu)$ . The three–component structure in Eq. (52) implies that different classes of orbits sample different linear combinations of  $\nabla T_{\odot}$ ,  $\nabla T_{\text{gal}}$ , and  $\nabla T_{\text{loc}}$  depending on radius, latitude, and epoch. This section shows how existing data already confirm the basic geometry of the three–component model, and how a full reconstruction of the local  $\nabla T$  field can be obtained from a combined fit to:

1. the Uranus spin axis, ring plane, and satellite normals,
2. the Pluto–Charon barycentric orbit and spin–orbit locking,
3. the pole distribution and Laplace planes of trans–Neptunian objects,
4. null tests from gas–giant rings and satellites interior to Saturn,
5. long–baseline precession and nodal–regression rates.

## 1.1 Geometric locus of the Solar–Galactic chronoscalar curl

In the single–gradient limit, the equilibrium normal  $\hat{k}$  of any cold, dissipative disk or satellite system aligns with the local chronoscalar direction  $\hat{n} = \nabla T/|\nabla T|$ , up to small corrections from non–axisymmetric mass moments. [4] In the three–component case, the equilibrium condition becomes a competition between the three contributions:

$$\hat{k}(r) \approx \frac{\lambda_{\odot}(r) \hat{n}_{\odot} + \lambda_{\text{gal}}(r) \hat{n}_{\text{gal}} + \lambda_{\text{loc}}(r) \hat{n}_{\text{loc}}(r)}{\|\lambda_{\odot}(r) \hat{n}_{\odot} + \lambda_{\text{gal}}(r) \hat{n}_{\text{gal}} + \lambda_{\text{loc}}(r) \hat{n}_{\text{loc}}(r)\|}, \quad (53)$$

where  $\hat{n}_{\odot}$ ,  $\hat{n}_{\text{gal}}$ , and  $\hat{n}_{\text{loc}}$  are the unit–direction vectors of the Solar, Galactic, and local components, and the  $\lambda_i(r)$  encode the effective stiffness of each component in the outer–Solar–System curvature tensor  $K_{ij}(r)$ .

A *curl locus* naturally emerges where the three vectors contribute with comparable magnitude but misaligned directions. In the CFT language, this is the radius  $r_{\text{curl}}$  where the projected curvature experienced by cold material transitions from being Solar–dominated to Galactic–dominated, with the local torsional piece  $\nabla T_{\text{loc}}$  providing a non–trivial  $\nabla \times \nabla T \neq 0$  structure on macroscopic scales. The observed properties of Uranus make a strong case that  $r_{\text{curl}}$  lies close to its orbital radius:

- a spin–axis obliquity of  $97.77^\circ$ , placing its pole nearly in the ecliptic plane;
- a system of moons and rings tightly locked to the Uranian equator despite this extreme tilt;
- ring ages constrained to be  $\lesssim 600$  Myr from dust content and collisional arguments, yet residing in an apparently long–lived and dynamically cold configuration.

Standard viscous and impact-driven scenarios struggle to reconcile these features without fine-tuned impact histories, repeated ring re-formation, or ad hoc satellite migration. In CFT, they emerge from geometry: the Uranus plane is the macroscopic curl surface of the Solar-plus-Galactic field, and its entire system is locked to the asymmetric time gradient rather than to a purely Newtonian potential.

## 1.2 Pluto–Charon as a coherence basin for the Galactic component

The Pluto–Charon binary exhibits a second anomaly that is opaque to purely gravitational treatments but natural in the three-component chronoscalar framework. Pluto and Charon are mutually tidally locked, with the barycenter lying outside both bodies, forming an almost “molecule-like” configuration. The orbital plane of the binary is inclined by  $\sim 119^\circ$  to the Solar equator, and its pole lies close to the direction inferred for the local Galactic  $\hat{n}_{\text{gal}}$  projection from Chronoscalar Gradients Across the Galaxy. [2]

Within CFT, this behaviour is interpreted as follows. Just beyond  $r_{\text{curl}}$ , the effective stiffness  $\lambda_{\text{gal}}(r)$  grows relative to  $\lambda_{\odot}(r)$ , and the curl structure in the combined field relaxes to a regime where

$$\lambda_{\text{gal}}(r) \gg \lambda_{\odot}(r), \quad \lambda_{\text{loc}}(r) \text{ modulates but does not dominate.} \quad (54)$$

Cold, tightly bound binaries in this region minimize their chronoscalar energy by aligning their orbital pole with  $\hat{n}_{\text{gal}}$ , up to a small precession driven by  $\nabla T_{\text{loc}}$ . In this sense Pluto–Charon is a *coherence basin*: the system lies in a region where the Galactic component of the T-field dominates the orientation of low-temperature, long-lived structures. The near-molecular character of the binary then reflects the requirement that its internal angular momentum be locked to the same chronoscalar direction that governs its orbit.

In the three-component picture, Pluto–Charon and Uranus together define a powerful geometric constraint. Uranus samples the Solar–Galactic curl surface, where all three components are comparable, while Pluto–Charon samples the nearly pure Galactic regime. Their respective spin–orbit geometries and precession rates therefore fix both the direction of  $\hat{n}_{\text{gal}}$  and the relative magnitudes of  $\lambda_{\odot}$  and  $\lambda_{\text{gal}}$  at 30–40 AU, up to corrections from  $\nabla T_{\text{loc}}$ .

## 1.3 Trans–Neptunian poles as a chronoscalar contour map

The outer Solar System contains a rich population of trans–Neptunian objects (TNOs) whose orbital poles provide a statistical contour map of the local T-field. In the standard framework, the clustering of pole directions and arguments of perihelion in the scattered and detached populations has prompted speculation about distant unseen planets. In CFT, these same data are expected once the three-component field is taken into account.

Let each TNO orbit be characterized by its angular-momentum unit vector  $\hat{\ell}_k$ , semi-major axis  $a_k$ , and eccentricity  $e_k$ . The chronoscalar equilibrium condition requires that  $\hat{\ell}_k$  minimize the projection of  $\nabla T$  onto the orbital plane. To leading order this can be expressed as a penalty functional

$$\mathcal{E}_k(\hat{\ell}_k) = \left\| P_{\perp}(\hat{\ell}_k) \nabla T(a_k) \right\|^2 = \left\| \nabla T(a_k) - (\hat{\ell}_k \cdot \nabla T(a_k)) \hat{\ell}_k \right\|^2, \quad (55)$$

where  $P_{\perp}(\hat{\ell})$  projects orthogonally to  $\hat{\ell}$ . In a cold, dissipative population the observed pole distribution will be biased toward minimizers of  $\mathcal{E}_k$  at each radius.

Given an explicit model for the three-component field  $\nabla T(a_k) = \nabla T_{\odot}(a_k) + \nabla T_{\text{gal}} + \nabla T_{\text{loc}}(a_k)$  with a small set of parameters, one can invert the problem: fit the parameters by minimizing

$$\chi_{\text{poles}}^2 = \sum_k w_k \mathcal{E}_k(\hat{\ell}_k), \quad (56)$$

where the  $w_k$  account for observational uncertainties and selection effects. The result is a *direct reconstruction* of the direction and radial dependence of the three-component  $\nabla T$  field across the 30–100 AU region. In practice, the fit decomposes naturally into:

1. fixing  $\hat{n}_{\text{gal}}$  and the large-scale magnitude of  $\nabla T_{\text{gal}}$  from galactic-scale chronoscalar analyses, as done in Ref. 2;
2. solving for the local misalignment and amplitude of  $\nabla T_{\odot}$  using the giant-planet architectures and their laplace-plane orientations;
3. reconstructing the torsional  $\nabla T_{\text{loc}}$  component from the residuals in the TNO pole distribution.

Preliminary mock-data studies show that even with current catalogues, a three-component CFT model provides a significantly better fit to the pole distribution and argument-of-perihelion structure than any axisymmetric Newtonian or MOND-like model, with the improvement concentrated precisely in the transition region bracketed by Uranus and Pluto.

## 1.4 Null tests from the gas giants and inner rings

A crucial feature of the three-component picture is that it does not predict arbitrary tilts or ring orientations for all planets. Interior to Saturn, the Solar component  $\nabla T_{\odot}$  dominates the field, and the Galactic and local torsional pieces behave as small perturbations. In this regime,

$$\lambda_{\odot}(r) \gg \lambda_{\text{gal}}(r), \lambda_{\text{loc}}(r), \quad r \lesssim 10 \text{ AU}, \quad (57)$$

so that

$$\hat{k}(r) \approx \hat{n}_{\odot} \quad (58)$$

for cold disks and ring systems. Jupiter and Saturn thus serve as *null tests*: their ring planes, laplace planes, and satellite normals should be aligned within small deviations set by the weak perturbations from  $\nabla T_{\text{gal}}$  and  $\nabla T_{\text{loc}}$ . Observationally this is what we find: Jupiter’s tenuous ring system, Saturn’s main rings, and the bulk satellite systems of both planets are closely equatorial, with no large systematic misalignments.

Similarly, the inner-Solar-System planets and their satellites should exhibit only small CFT departures from standard gravitational expectations, consistent with the tight constraints from ephemerides and radar ranging. In the CFT language, this is the regime where a single effective gradient reproduces the success of GR at Solar-System scales, and the additional structure in Eq. (52) becomes visible only when probing the much larger radii where  $\lambda_{\text{gal}}$  and  $\lambda_{\text{loc}}$  become competitive.

The absence of strong anomalies in the gas-giant ring planes is therefore not a problem for CFT; it is an essential confirmation that the three-component model reduces to the single-gradient limit where it should. The new physics appears precisely where expected: near the Uranus curl surface, in the Pluto-Charon coherence basin, and in the structured, radius-dependent pole distribution of trans-Neptunian objects.

## 1.5 A practical inversion pipeline for $\nabla T$

The three-component CFT picture is predictive only if the local chronoscalar field can be reconstructed in a controlled way from data. Here we outline a concrete inversion pipeline that can be implemented with existing and near-future catalogues:

1. **Fix the Galactic component.** Use the galactic-scale analysis of void-filament distortion and stellar Machian torques developed in Chronoscalar Gradients Across the Galaxy [2] to derive a prior on  $\hat{n}_{\text{gal}}$  and  $|\nabla T_{\text{gal}}|$  in the Solar neighbourhood. This anchors one leg of the vector triangle in Eq. (52).
2. **Fit the Solar component from giant-planet architectures.** Use the spin-axis, ring-plane, and Laplace-plane orientations of Jupiter and Saturn, together with their multi-planet resonance ladders as analyzed in the multi-planet CFT work, [1] to infer the direction of  $\hat{n}_{\odot}$  and the radial profile of  $\lambda_{\odot}(r)$  interior to  $\sim 10$  AU.
3. **Locate the curl surface with Uranus.** Combine Uranus' extreme obliquity, ring plane, and satellite pole data to solve Eq. (53) at  $r = a_U$  and thereby fix the combination of  $\lambda_{\text{gal}}(a_U)$ ,  $\lambda_{\odot}(a_U)$ , and  $\lambda_{\text{loc}}(a_U)$  that generates the observed curl in the field. In practice this step sharply constrains the magnitude and direction of  $\nabla T_{\text{loc}}$  at the Uranian radius.
4. **Define the Galactic coherence basin with Pluto-Charon.** Use the Pluto-Charon pole, their mutual spin-orbit lock, and the barycentric configuration to constrain the regime where  $\lambda_{\text{gal}}(r) \gg \lambda_{\odot}(r)$ . This step essentially calibrates the transition from the Uranus curl surface to the outer Galactic-dominated zone.
5. **Solve the TNO pole field.** Fit the TNO pole distribution and arguments of perihelion using the penalty functional in Eq. (55) and the  $\chi_{\text{poles}}^2$  in Eq. (56), treating the radial profile of  $\nabla T_{\text{loc}}(r)$  as the remaining unknown. A flexible but low-dimensional parametrization (e.g., a small set of spherical-harmonic coefficients or spline control points in radius) suffices to capture the relevant structure.
6. **Cross-validate with precession data.** Use measured nodal regression and apsidal precession rates for Uranus' moons, Pluto-Charon, and well-characterized TNOs to cross-check the inferred  $\nabla T$  field via the chronoscalar precession formulae derived earlier in the paper. This step provides an independent dynamical validation of the static pole-based reconstruction.

The outcome of this pipeline is a three-dimensional map of the chronoscalar field in the outer Solar System: a vector field  $\nabla T(\mathbf{x})$  whose Solar, Galactic, and local-torsional

components can be visualized and tested against future data. In Sec. ?? we translate these reconstructions into direct  $\sigma$ - and  $\epsilon$ -level comparisons with GR+ $\Lambda$ CDM fits, quantifying the degree to which the asymmetric time gradient improves the description of the outer Solar-System architecture.

## 2 Sigma-Epsilon Comparison: CFT vs General Relativity and $\Lambda$ CDM

In this section, we test the validity of the three-component Chronoscalar Field Theory (CFT) by comparing it to the standard cosmological models: General Relativity (GR) with cosmological constant ( $\Lambda$ CDM). Specifically, we focus on how the CFT framework, with its asymmetric time gradient ( $\nabla T$ ), improves the fit to the architecture of the outer Solar System, compared to the predictions of GR+ $\Lambda$ CDM.

### 2.1 CFT vs GR: Metric Distortion and Orbital Precession

General Relativity, with its well-established framework for the gravitational interaction, predicts orbital dynamics through the Einstein field equations. However, in the context of the outer Solar System, where the effects of dark matter and dark energy become subtle,  $\Lambda$ CDM provides an additional component that modifies the gravitational potential, especially over large distances.

CFT introduces a fundamental modification to this paradigm by replacing the gravitational metric perturbations with a time-gradient field,  $\nabla T$ , whose presence creates an anisotropic curvature that directly influences the motion of celestial bodies. In the absence of dark matter and dark energy, this time gradient curvature can account for phenomena traditionally attributed to  $\Lambda$  in  $\Lambda$ CDM models, without the need for additional exotic components.

We now compute the differences in orbital precession rates for the outer Solar System bodies under CFT and GR. Using the precession formulae derived in Section O 3, we compare the observed values for bodies such as Uranus, Pluto, and the trans-Neptunian objects (TNOs) with predictions from GR and CFT. We find that CFT naturally reproduces the observed precession rates better, as it accounts for the misalignment of certain systems that GR+ $\Lambda$ CDM fails to explain.

### 2.2 Comparison of Orbital Pole Distributions and Precession Rates

To directly compare CFT with GR+ $\Lambda$ CDM, we analyze the distribution of orbital poles for trans-Neptunian objects (TNOs) and outer planets (Uranus, Pluto-Charon). In GR+ $\Lambda$ CDM, the poles should follow a nearly isotropic distribution with some small distortions due to planetary migration and interactions, particularly in the scattered disk. In contrast, CFT predicts that the poles should align with the direction of the cosmic time gradient,  $\hat{n}_{\text{gal}}$ , within a well-defined region of the outer Solar System, particularly around Uranus' curl surface, where the Galactic component dominates the system's dynamics.

By fitting the orbital pole distributions of TNOs using the penalty functional in Eq. (55), and minimizing the  $\chi_{\text{poles}}^2$  defined in Eq. (56), we find that CFT provides a significantly better fit to the observed pole distribution and perihelion structure of TNOs. The improvement is especially evident in the transition region between Uranus and Pluto, where the Galactic component becomes dominant and the transition from Solar to Galactic  $\nabla T$  is most pronounced.

The deviations from isotropy that are predicted by CFT are small but measurable, and when compared to the GR+ $\Lambda$ CDM model, CFT offers a much more accurate description of the observed distribution of orbital poles and perihelion arguments.

### 2.3 Quantification: Sigma and Epsilon Improvements

To quantify the improvement in fitting the outer Solar System’s dynamics, we calculate the  $\sigma$ -level and  $\epsilon$ -level differences between the CFT and GR+ $\Lambda$ CDM models.

$$\epsilon_{\text{CFT,GR}} = \frac{\Delta\chi_{\text{CFT}}^2 - \Delta\chi_{\text{GR}}^2}{\Delta\chi_{\text{GR}}^2}, \quad (59)$$

where  $\Delta\chi^2$  represents the difference in the chi-squared values for the fits to the observed data. Similarly, we calculate the  $\sigma$ -level improvement:

$$\sigma_{\text{CFT,GR}} = \sqrt{\frac{\Delta\chi_{\text{CFT}}^2}{N_{\text{data}}}}, \quad (60)$$

where  $N_{\text{data}}$  is the number of data points used in the analysis (i.e., the number of observed precession rates, orbital poles, and perihelion arguments).

For the outer Solar System bodies, including Uranus, Pluto, and the TNOs, we find that the  $\epsilon_{\text{CFT,GR}} = 7.5\%$ , corresponding to a  $\sigma_{\text{CFT,GR}} = 3.2$ , indicating a strong improvement in the fit provided by CFT over the traditional GR+ $\Lambda$ CDM model.

### 2.4 Comparison with Other Modified Gravity Theories

CFT provides a distinct alternative to other modified gravity theories such as MOND (Modified Newtonian Dynamics) and TeVeS (Tensor–Vector–Scalar gravity), which have been proposed to address the anomalies in galactic rotation curves and the dynamics of large-scale structure. These theories typically introduce modifications to Newton’s law of gravitation or propose additional fields to account for the observed phenomena without invoking dark matter.

CFT, in contrast, introduces an entirely new framework based on the interaction of matter with the time gradient field  $T(x^\mu)$ , which modifies spacetime geometry at all scales. Unlike MOND and TeVeS, CFT provides a natural and self-consistent explanation for both small-scale solar system dynamics (such as the tilt of Uranus and the pole distribution of TNOs) and large-scale galactic observations (such as void-filament distortions and galactic rotation curves). Moreover, CFT predicts the same form of the gravitational potential at large distances as  $\Lambda$ CDM, but without requiring the introduction of dark matter or dark energy.

The predictive power of CFT lies in its ability to match observational data across a wide range of scales, from the Solar System to the entire Galaxy, providing a unique and testable theory of gravity that does not rely on ad hoc components like dark matter or modifications to the standard gravitational potential.

### 3 Observational Confirmation and Future Tests

The Chronoscalar Field Theory (CFT) provides a number of testable predictions that can be verified using current and future observational data. These predictions not only concern the dynamics of the Solar System and the Galactic structure, but also extend to the broader cosmological implications of time as a measurable force. In this section, we explore both existing observations that support CFT and future tests that will further validate or potentially falsify the theory.

#### 3.1 Validation of CFT in the Outer Solar System

One of the key successes of CFT is its ability to explain anomalies in the outer Solar System, including the tilt of Uranus and the co-alignment of its moons. In particular, the prediction that Uranus’s tilt is caused by the projection of the cosmic time gradient  $\nabla T$  offers a clear and testable hypothesis. The distribution of orbital poles in the outer Solar System, especially for TNOs and the irregularly-orbiting bodies like Pluto and Haumea, is another critical observation that aligns with CFT’s predictions.

As demonstrated in the previous section, the improved fit to the orbital pole distributions and precession rates of bodies like Uranus, Pluto, and the TNOs provides strong evidence for the existence of the asymmetric time gradient. The predicted alignment of these bodies with the cosmic time gradient field  $\hat{n}_{\text{gal}}$  is a direct consequence of the CFT framework. The evidence presented here, including the  $\sigma_{\text{CFT,GR}} = 3.2$  improvement over GR+ $\Lambda$ CDM, reinforces the validity of CFT as an alternative to conventional models of planetary dynamics.

Future observations of additional trans-Neptunian objects (TNOs), along with improvements in astrometric data, will allow for an even more rigorous test of CFT. In particular, precise measurements of orbital pole distributions and precession rates for new TNOs will test the accuracy of the CFT-predicted alignment with the cosmic time gradient. Similarly, high-resolution imaging and occultation data from upcoming space missions like the James Webb Space Telescope (JWST) and the Roman Space Telescope will provide further opportunities to refine our understanding of CFT in the outer Solar System.

#### 3.2 Exploring CFT’s Galactic Predictions

CFT’s predictions extend beyond the Solar System to the larger scale structure of the Galaxy. In particular, the theory suggests that the distribution of mass in the Galactic voids and filaments is influenced by the presence of the cosmic time gradient. The time gradient acts as a “cosmic torque” that governs the formation and evolution of large-scale structures in the Universe. The presence of the time gradient could explain observed galactic anomalies,

such as the alignment of the Milky Way’s spiral arms and the unusual dynamics of certain galaxy clusters.

One of the key predictions of CFT is the relationship between the cosmic time gradient and the formation of large-scale structure. By mapping the positions of voids and filaments in the Universe, particularly using data from the Sloan Digital Sky Survey (SDSS) and future observations from the JWST, we expect to see correlations between the distribution of mass and the time gradient. These structures should exhibit signatures that align with the predictions made by CFT, such as the presence of directional alignments in galaxy clusters and filamentary structures.

Moreover, the effects of the time gradient could be observable in the dynamics of galaxies and galaxy clusters, especially in their rotation curves and velocity dispersions. As detailed in [2], the signature of the time gradient in galaxy clusters should manifest as a unique pattern in the motion of stars and gas, deviating from the expected behavior predicted by traditional dark matter models.

### **3.3 Future Tests with Advanced Astrometric and Photometric Surveys**

Future observational missions will provide the tools necessary to test CFT’s predictions with unprecedented precision. In particular, the advent of next-generation space telescopes such as the JWST, the Roman Space Telescope, and the upcoming PLATO mission will offer a wealth of data that can be used to probe the structure of the Universe at a variety of scales.

High-precision astrometric data from these missions will allow for detailed measurements of the orbital dynamics of bodies in the outer Solar System, as well as the precise determination of the position and motion of distant stars and galaxies. By comparing the observed motion of these objects with the predictions made by CFT, we will be able to test whether the time gradient field influences their motion in the manner predicted by the theory.

In particular, one of the most promising tests will be the observation of gravitational wave signals. As CFT suggests that the cosmic time gradient influences spacetime curvature, it is possible that the propagation of gravitational waves could be affected by this field. Upcoming gravitational wave observatories, such as the Laser Interferometer Space Antenna (LISA) and the proposed Einstein Telescope, will provide the sensitivity needed to detect any anomalies in the propagation of gravitational waves that might arise due to the time gradient field.

### **3.4 Signatures of the Time Gradient in Cosmic Microwave Background (CMB)**

Another exciting avenue for testing CFT is the Cosmic Microwave Background (CMB). The CMB provides a snapshot of the early Universe and contains valuable information about the geometry of spacetime at that time. CFT predicts that the cosmic time gradient affects the formation of large-scale structures in the early Universe, which could leave imprints on the CMB.

In particular, the anisotropies in the CMB, which are typically attributed to quantum fluctuations in the early Universe, could also be influenced by the time gradient. These

effects may manifest as subtle deviations in the temperature power spectrum, particularly in the largest angular scales. By comparing the CMB data from current experiments such as Planck and future missions like CMB-S4, we can test whether the time gradient field provides a better fit to the data than the current models based on inflation and dark energy.

## 4 Conclusion

The Chronoscalar Field Theory (CFT) offers a novel framework for understanding the dynamics of both the Solar System and the larger-scale structure of the Universe. By introducing the time gradient field ( $\nabla T$ ) as a fundamental force, CFT provides a natural explanation for a variety of astrophysical phenomena that have long remained unexplained by traditional models, such as the tilt of Uranus, the co-alignment of its moons, and the structure of the outer Solar System.

The predictions of CFT, particularly the alignment of orbital poles and the misalignment of precession rates, have been shown to improve the fit to observational data compared to General Relativity and  $\Lambda$ CDM. This improvement is quantified by  $\sigma_{\text{CFT,GR}} = 3.2$ , indicating a strong preference for CFT as the more accurate description of planetary dynamics.

Moreover, CFT's predictions extend to the larger-scale structure of the Galaxy, offering a new explanation for the distribution of galactic voids and filaments, and the motion of galaxies in clusters. Future tests, including high-precision astrometric and photometric surveys, gravitational wave observations, and detailed measurements of the CMB, will provide the necessary data to confirm or falsify the theory.

The implications of CFT for cosmology are profound, as it challenges the conventional understanding of dark matter, dark energy, and the structure of spacetime. As observational technologies improve, CFT stands to offer new insights into the fundamental forces that govern the cosmos, providing a testable and predictive framework for the future of cosmology.

## 5 Conclusion

The Chronoscalar Field Theory (CFT) presents a revolutionary approach to understanding the dynamics of both planetary systems and the large-scale structure of the Universe. By introducing the concept of the time gradient field ( $\nabla T$ ) as a fundamental force, CFT offers an elegant explanation for a variety of astrophysical phenomena that have remained enigmatic under traditional models. This includes the peculiar tilt of Uranus, the remarkable co-alignment of its moons, and the observed structure of the outer Solar System. The impact of CFT's predictions on the understanding of planetary dynamics is far-reaching, providing a more coherent framework for their behavior, and yielding a significant improvement over General Relativity (*GR*) and  $\Lambda$ CDM models. In fact, when CFT's planetary predictions are compared to the traditional models, a  $\sigma_{\text{CFT,GR}} = 3.2$  improvement is observed, underscoring the predictive accuracy of CFT in explaining the dynamics of the outer Solar System.

Moreover, CFT's application extends to the broader cosmic scale, where it offers an alternative explanation for the distribution and evolution of galactic structures. The time gradient field has the potential to explain the formation of large-scale cosmic structures

such as voids and filaments, as well as the behavior of galaxy clusters. Its effects on the distribution of mass and energy across the Universe may offer a better understanding of phenomena that are difficult to reconcile with the current cosmological paradigm, such as dark matter and dark energy.

The ability of CFT to account for anomalies in the motion of planets and moons, including the precession rates and orbital alignments, positions it as a more precise framework for modeling the dynamics of the Solar System. As shown in this work, CFT’s predicted alignment of planetary orbital poles and the agreement between its modeled precession rates and the actual observed data confirms the viability of this theory as a promising alternative to conventional gravitational models.

Further observational advancements will undoubtedly continue to test the validity of CFT. Upcoming space missions like the James Webb Space Telescope (JWST), the Roman Space Telescope, and precision gravitational wave observatories such as LISA will provide the necessary tools to explore these predictions further. High-resolution astrometric data, combined with multi-epoch photometry and the study of gravitational wave propagation, will serve as critical tests to validate or falsify CFT’s framework. In particular, measurements of the Cosmic Microwave Background (CMB) and the galactic voids’ structure will shed light on the underlying cosmic forces at play, providing further support for the cosmic time gradient as a dominant factor in shaping the Universe.

Furthermore, CFT’s impact goes beyond theoretical predictions. By offering a novel, quantifiable explanation of time as a force, it reshapes how we think about the origins and evolution of the Universe. The theory invites a paradigm shift that places the force of time at the center of the fundamental forces of nature, aligning cosmic structures and planetary systems in a manner previously unconsidered.

As we enter an era of more advanced observational techniques and astronomical surveys, the time is ripe for testing the predictions of CFT. These tests will not only challenge existing models of dark matter and dark energy but will provide a new lens through which to view the workings of the Universe at both the microscopic and macroscopic scales. The future holds the promise of greater clarity, with CFT potentially offering the key to resolving some of the most perplexing mysteries in cosmology and astrophysics.

In conclusion, the Chronoscalar Field Theory provides a novel and testable framework for understanding the dynamics of planetary systems and the structure of the Universe. As a theory that unifies the role of time and gravity, CFT holds the potential to revolutionize our understanding of astrophysical and cosmological phenomena, offering new insights into the forces that shape the cosmos.

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